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# **Extraction of pseudo-PDFs**

## ***or why Ioffe-Time Distributions are our friend***

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# Outline

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- LaMET, quasi-PDFs, pseudo-PDFs and Good Lattice Cross Sections
- Pion as a theatre for PDFs - pPDFs and GLCS
- pPDFs in Nucleon
- Summary

# Introduction

- **First Challenge:**

- Euclidean lattice precludes calculation of light-cone/time-separated correlation functions

$$q(x, \mu) = \int \frac{d\xi^-}{4\pi} e^{-ix\xi^- P^+} \langle P | \bar{\psi}(\xi^-) \gamma^+ e^{-ig \int_0^{\xi^-} d\eta^- A^+(\eta^-)} \psi(0) | P \rangle$$

So.... ...Use *Operator-Product-Expansion to formulate in terms of Mellin Moments* with respect to Bjorken x.

$$\longrightarrow \langle P | \bar{\psi} \gamma_{\mu_1} (\gamma_5) D_{\mu_2} \dots D_{\mu_n} \psi | P \rangle \rightarrow P_{\mu_1} \dots P_{\mu_n} a^{(n)}$$

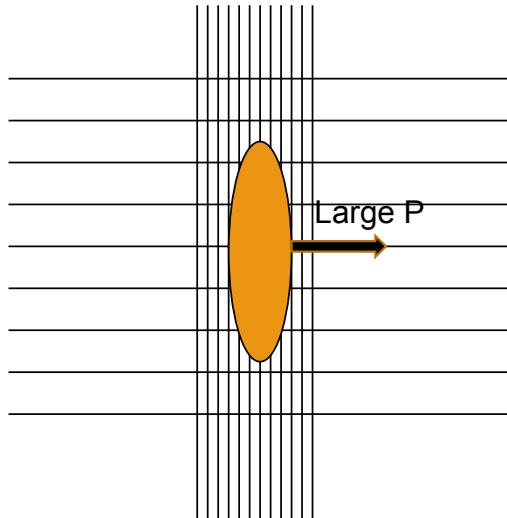
- **Second Challenge:**

- Discretised lattice: power-divergent mixing for higher moments

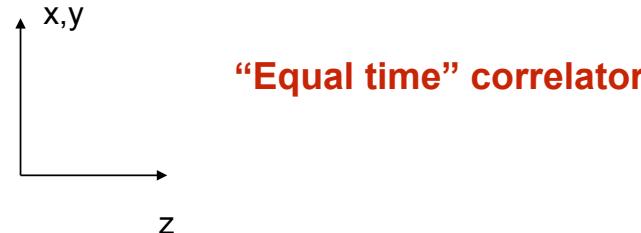
## Moment Methods

- Extended operators: Z.Davoudi and M. Savage, PRD 86,054505 (2012)
  - Valence heavy quark: W.Detmold and W.Lin, PRD73, 014501 (2006)

# Solution....



Large-Momentum Effective Theory (LaMET)



X. Ji, Phys. Rev. Lett. 110, 262002 (2013).

X. Ji, J. Zhang, and Y. Zhao, Phys. Rev. Lett. 111, 112002 (2013).

J. W. Qiu and Y. Q. Ma, arXiv:1404.686.

$$q(x, \mu^2, P^z) = \int \frac{dz}{4\pi} e^{izk^z} \langle P | \bar{\psi}(z) \gamma^z e^{-ig \int_0^z dz' A^z(z')} \psi(0) | P \rangle + \mathcal{O}((\Lambda^2/(P^z)^2), M^2/(P^z)^2)$$



$$q(x, \mu^2, P^z) = \int_x^1 \frac{dy}{y} Z\left(\frac{x}{y}, \frac{\mu}{P^z}\right) q(y, \mu^2) + \mathcal{O}(\Lambda^2/(P^z)^2, M^2/(P^z)^2)$$

# Pseudo-PDFs

- Pseudo-PDF (pPDF) recognizing generalization of PDFs in terms of *Ioffe Time*.  $\nu = p \cdot z$

A.Radyushkin, Phys. Rev. D 96, 034025 (2017)

B.Ioffe, PL39B, 123 (1969); V.Braun  
et al, PRD51, 6036 (1995)

$$M^\alpha(p, z) = \langle p | \bar{\psi} \gamma^\alpha U(z; 0) \psi(0) | p \rangle$$

$p = (p^+, m^2/2p^+, 0_T)$      $\downarrow$      $z = (0, z_-, 0_T)$                   Ioffe-Time Distribution

$$M^\alpha(z, p) = 2p^\alpha \mathcal{M}(\nu, z^2) + 2z^\alpha \mathcal{N}(\nu, z^2)$$

Ioffe-time pseudo-Distribution (**pseudo-ITD**) generalization to *space-like z*

Lattice “building blocks” that of quasi-PDF approach.

$$\mathcal{M}(\nu, z^2) = \int_{-1}^1 dx e^{i\nu x} \mathcal{P}(x, z^2)$$

$\Downarrow$  *Lorentz covariant*

$$f(x) = \mathcal{P}(x, 0) \underset{z_3^2 \rightarrow 0}{=} \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} \mathcal{M}(\nu, -z_3^2)$$

# pPDFs - II

To deal with UV divergences, introduce reduced distribution  $\mathfrak{M} = \frac{\mathcal{M}(\nu, z^2)}{\mathcal{M}(0, z^2)}$

$$\mathfrak{M}(\nu, z^2) = \int_0^1 du K(u, z^2 \mu^2, \alpha_s) Q(u\nu, \mu^2)$$

Computed on lattice

Perturbatively calculable

Ioffe-time Distribution

$$Q(\nu, \mu) = \mathfrak{M}(\nu, z^2) - \frac{\alpha_s C_F}{2\pi} \int_0^1 du \left[ \ln \left( z^2 \mu^2 \frac{e^{2\gamma_E+1}}{4} \right) B(u) + L(u) \right] \mathfrak{M}(u\nu, z^2).$$

K. Orginos et al.,  
PRD96 (2017),  
094503

Match data at different  $z$

Inverse problem

$$Q(\nu) = \int_{-1}^1 dx q(x) e^{i\nu x}$$

$$q(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} Q(\nu)$$

Need data for all  $\nu$ , or  
*additional physics input*

# Moments

J Karpie, K Orginos, S Zafeiropoulos, arXiv:1807.10933

- Can obviate need for inverse through computation of moments

$$\mathfrak{M}(\nu, z^2) = \sum_{n=0}^{\infty} i^n \frac{\nu^n}{n!} a_n(\mu^2) K_n(\mu^2 z^2) + \mathcal{O}(z^2)$$

Mellin moments of PDF

Matching coefficient

# “Good Lattice Cross Sections”

$$\sigma_n(\nu, \xi^2, P^2) = \langle P | T\{\mathcal{O}_n(\xi)\} | P \rangle$$

Ma and Qiu, Phys. Rev. Lett. 120 022003

*Expressed in coordinate space*

where

$$\sigma_n(\nu, \xi^2, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(x\nu, \xi^2, x^2 P^2, \mu^2) + \mathcal{O}(\xi^2 \Lambda_{\text{QCD}}^2)$$

*Calculated in  
LQCD*

Parton Distribution  
function

Short distance scale

*Calculated in perturbation  
theory (“process dependent”)*

$$\mathcal{O}(\xi) = \bar{\psi}(0) \Gamma W(0, 0 + \xi) \psi(\xi)$$

← Encompasses qPDF/pPDF

$$\mathcal{O}_S(\xi) = \xi^4 Z_S^2 [\bar{\psi}_q \psi_q](\xi) [\bar{\psi}_q \psi](0)$$

Gauge-Invariant Currents

$$\mathcal{O}_{V'}(\xi) = \xi^2 Z_{V'}^2 [\bar{\psi}_q \xi \cdot \gamma \psi_{q'}](\xi) [\bar{\psi}_{q'} \xi \cdot \gamma \psi](0)$$

← Flavor-changing

+ analogous gluon operators

# Quasi- vs Pseudo- vs GLCS

- All integrals of Ioffe-Time Distribution Function
- Should yield same PDF after matching and systematic controls

## Quasi-PDF

$$Q(x, p_3^2) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} \mathcal{M}(\nu, -\nu^2/p_3^2)$$

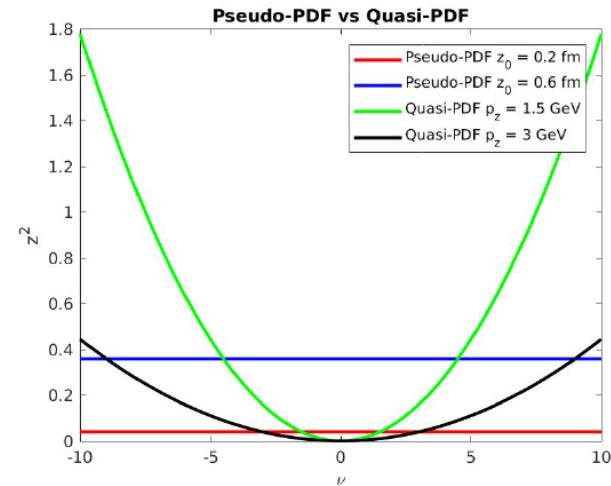
$$\mathcal{P}(x, -z_3^2) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\nu e^{-i\nu x} \mathcal{M}(\nu, -z_3^2)$$

## Pseudo-PDF and GLCS

For pPDF + GLCS,  $z$  sets short-distance scale.

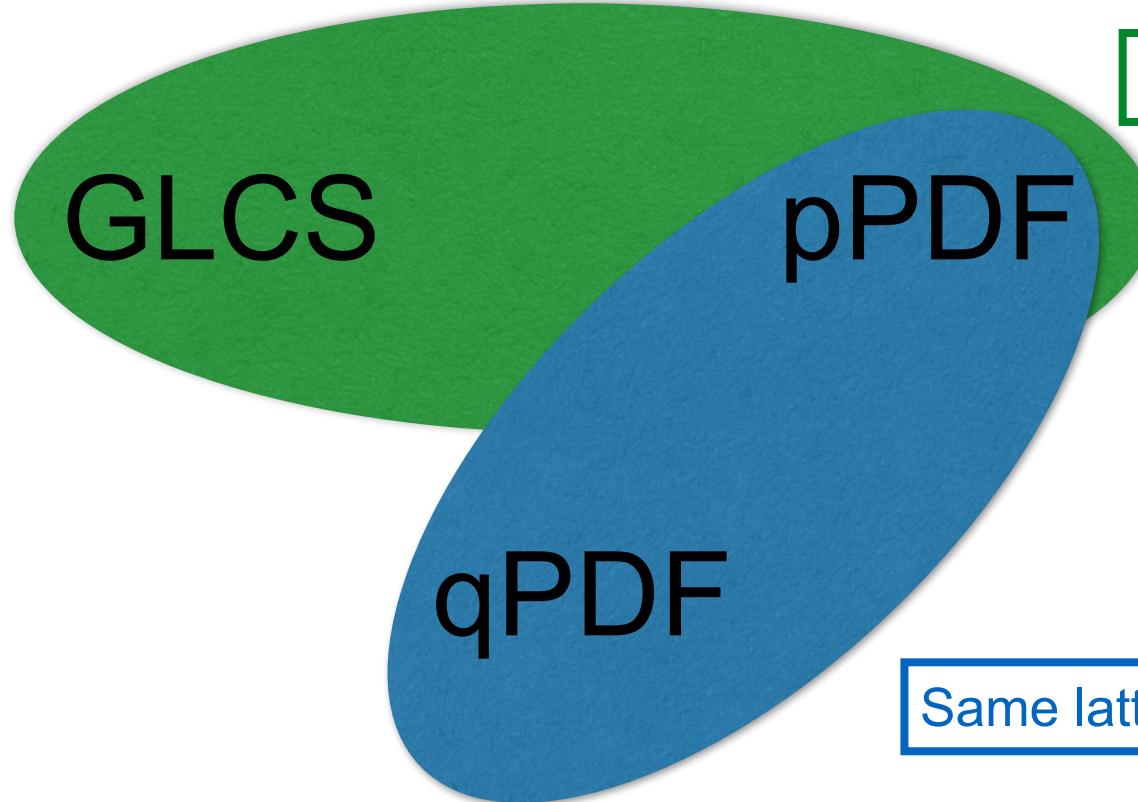
$$z \ll \frac{1}{\Lambda_{\text{QCD}}}$$

N.B. All approaches require large momentum - but for pPDF and GLCS to ensure range in Ioffe time to solve *inverse problem*.



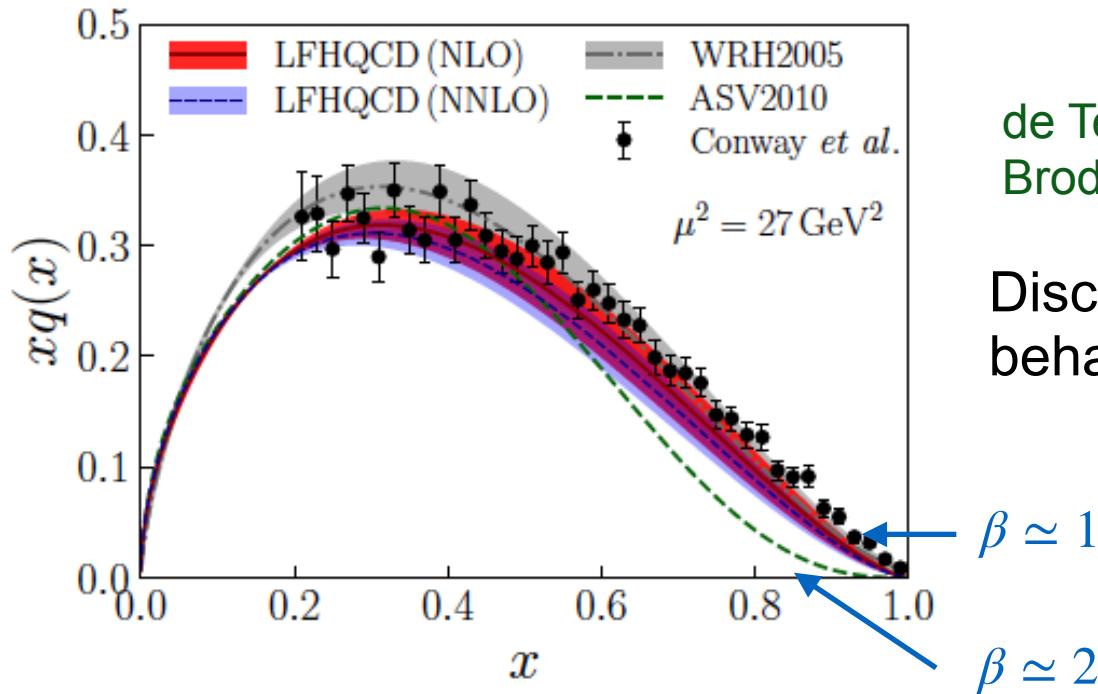
$$P \longrightarrow \sqrt{s} \quad \text{Collision energy}$$
$$z \longrightarrow \frac{1}{Q} \quad \text{Hard Probe}$$

Analogous matching to light-cone PDFs



# Pion Valence PDF

- u distribution of FNAL E615 to leading order
- C12-15-006 at Hall A will look at structure of pion
- C12-15-006A at Hall A will look at structure of Kaon
- No free pion target



de Teramond, liu, Sufian, Dosch,  
Brodsky, Deur, PRL (2018)

Discrepancy in large- $x$   
behavior of pion distribution

# Why the Pion?

- Pion less computationally demanding than nucleon?

- *Larger signal-to-noise ratio*

$$C(t, \vec{p}) \equiv \sum \langle 0 | \mathcal{O}(t, \vec{x}) \mathcal{O}^\dagger(0, 0) | 0 \rangle e^{-i\vec{p}\cdot\vec{x}} \rightarrow e^{-E(\vec{p})t}$$

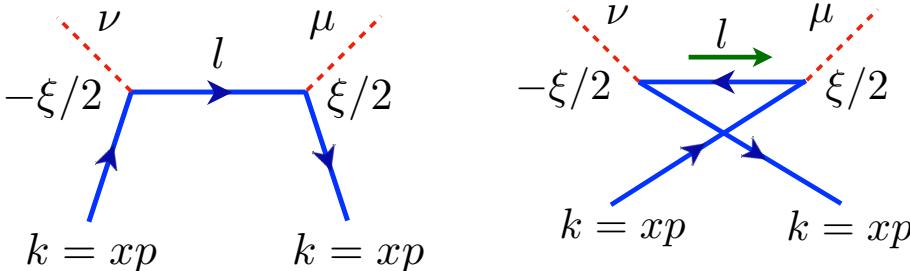
$$C_{\sqrt{\sigma^2}}(t, \vec{p}) \rightarrow \begin{cases} e^{-m_\pi t} & \text{Mesons} \\ e^{-(3m_\pi/2)t} & \text{Baryons} \end{cases}$$

- Important constraint on systematic uncertainty is understanding operator renormalization
  - *Operator renormalization “independent” of external states*
- Several different calculations using the different approaches
  - *Lattice cross-section approach straightforward for mesons, challenging for baryons*

# Good Lattice Cross Section

$$\begin{aligned}\sigma_{ij}^{\mu\nu}(\xi, p) &= \langle \pi(p) | \mathcal{O}_{ij}^{\mu\nu}(\xi) | \pi(p) \rangle \\ &= \xi^4 \langle \pi(p) | \mathcal{J}_i^\mu(\xi/2) \mathcal{J}^\nu + j(-\xi/2) | \pi(p) \rangle\end{aligned}$$

Calculate  $K$  at tree-level between quark states

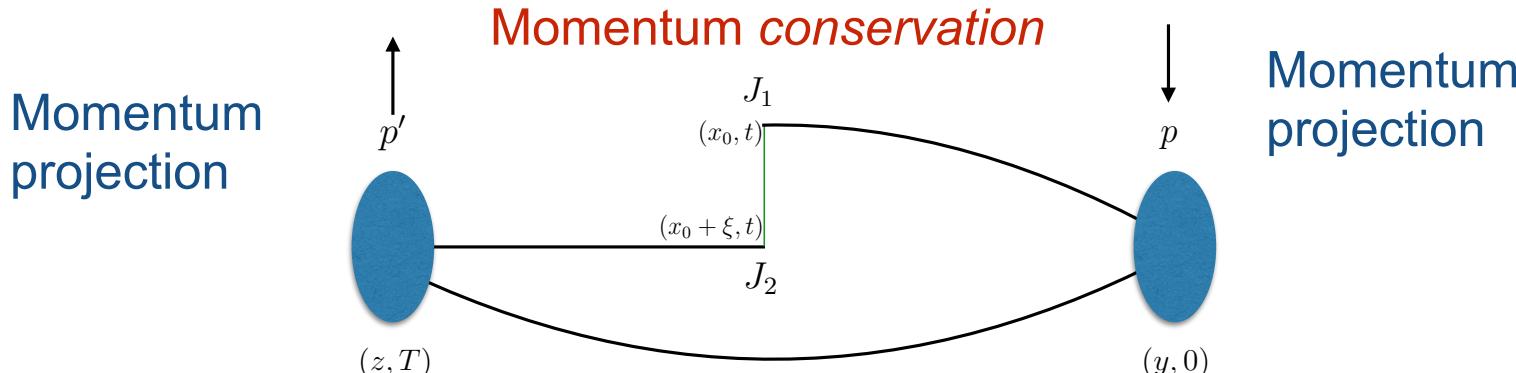


Process, i.e. current, dependent

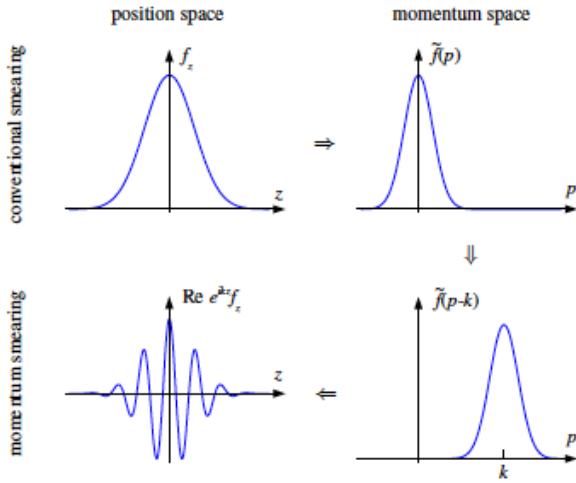
$$\frac{1}{2} [\sigma_{V,A}^{\mu\nu}(\xi, p) + \sigma_{A,V}^{\mu\nu}(\xi, p)]$$

$$\equiv \epsilon^{\mu\nu\alpha\beta} \xi_\alpha p_\beta T_1(\nu, \xi^2) + (p^\mu \xi^\nu - \xi^\mu p^\nu) T_2(\nu, \xi^2)$$

Sequential-Source Approach

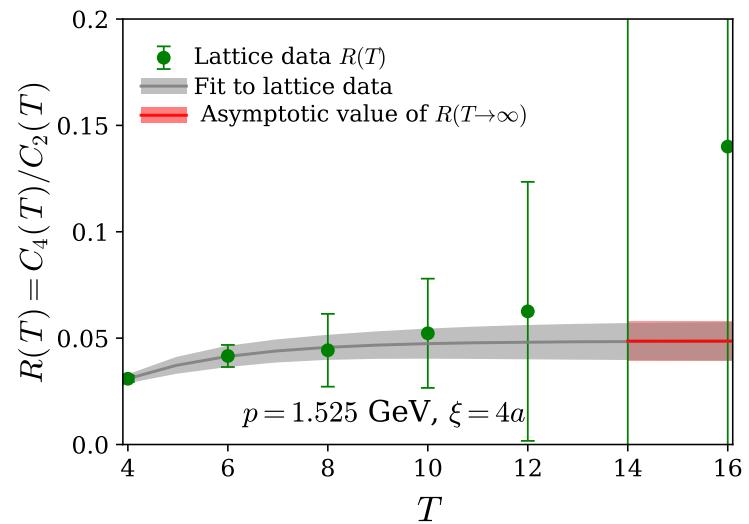
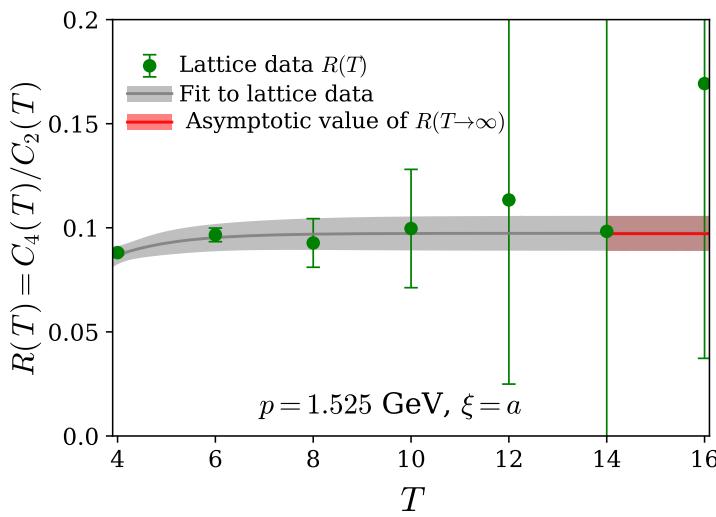


# Challenges of Higher Momenta

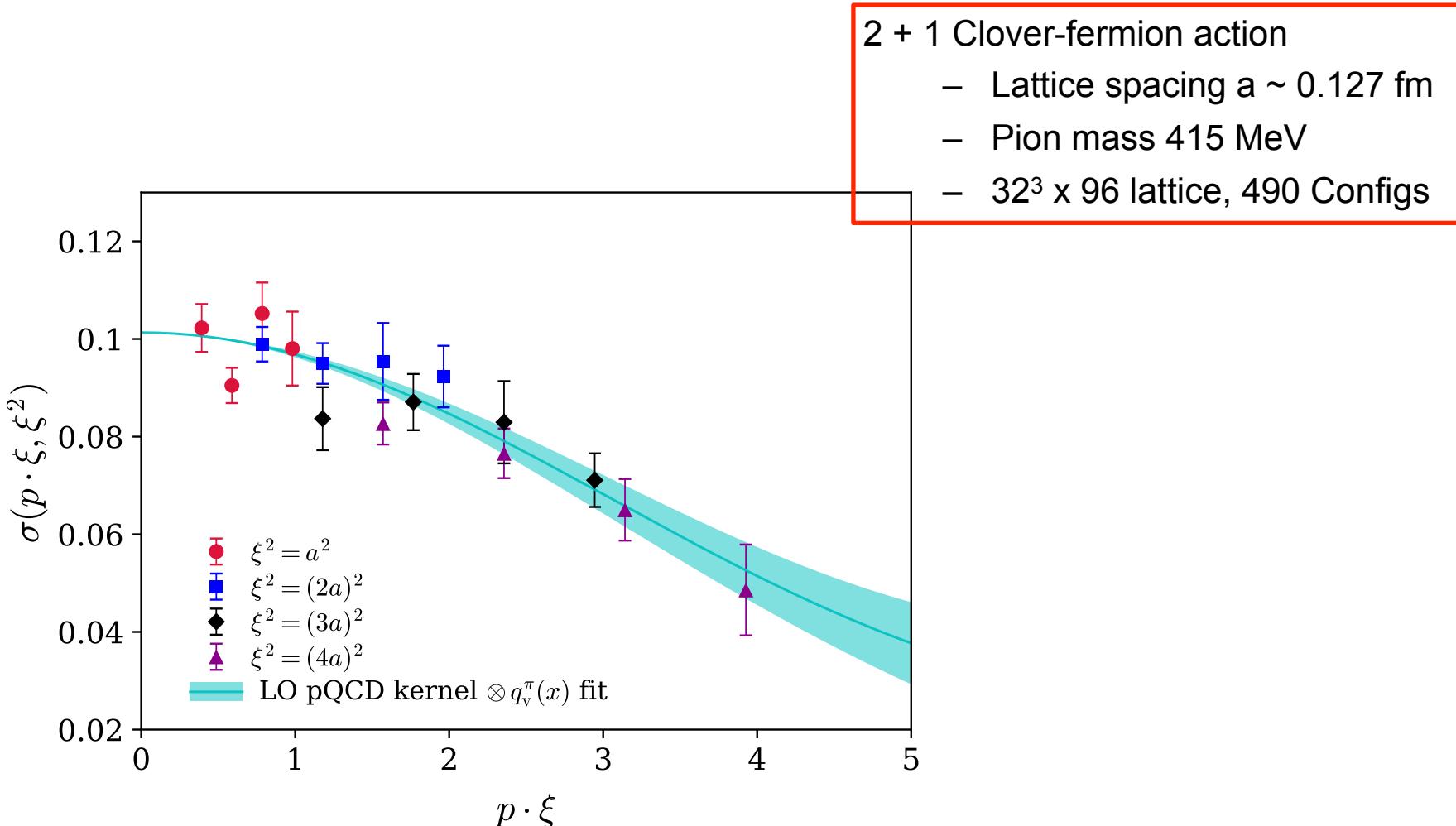


Boosted interpolating operators

Bali *et al.*, Phys. Rev. D 93, 094515 (2016)



# Good Lattice Cross Section



# Inverse problem: extract PDF

“Inverse Problem” - ill-posed inverse Fourier transform.

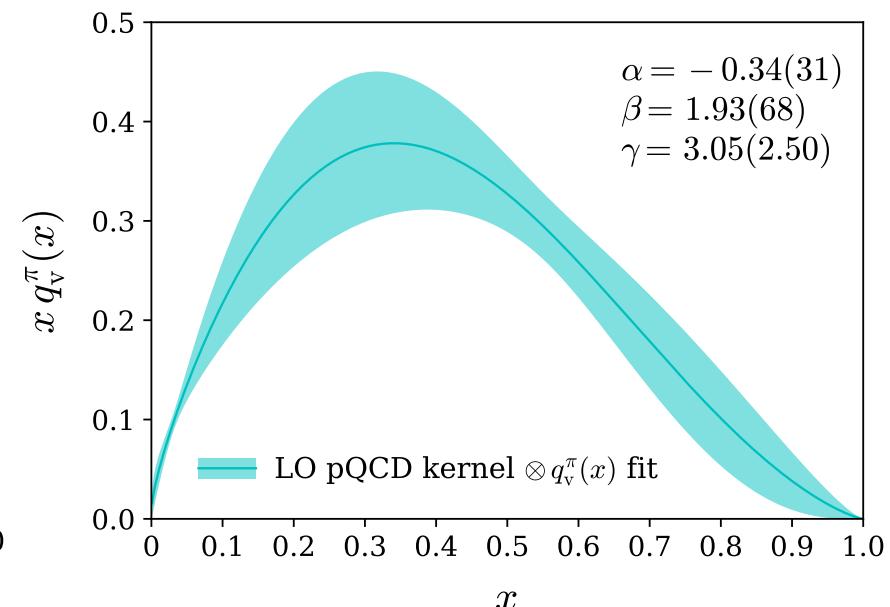
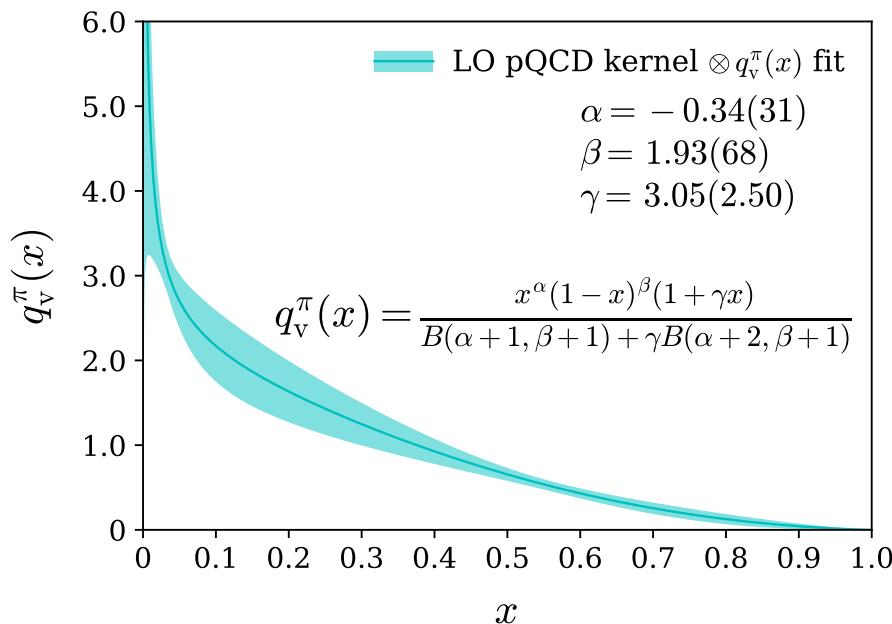
$$\sigma_n(\nu, \xi^2, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(x\nu, \xi^2, x^2 P^2, \mu^2) + \mathcal{O}(\xi^2 \Lambda_{\text{QCD}}^2)$$

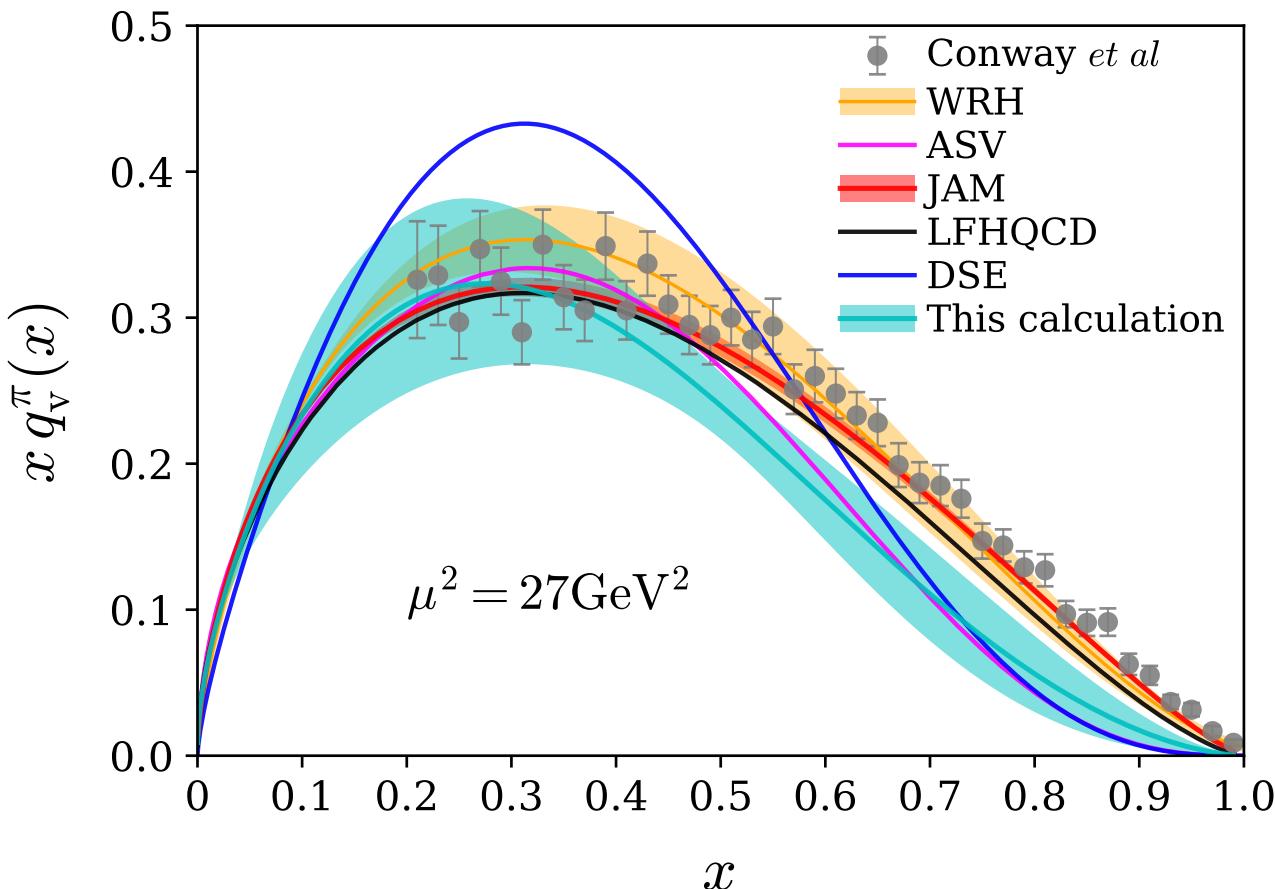
Calculate on Lattice

Calculate in PQCD

Extract PDF?

Similar challenge to global fitting community!





LO Calculation: “guess” scale

Systematic study at final lattice  
spacing NLO kernel: coming soon

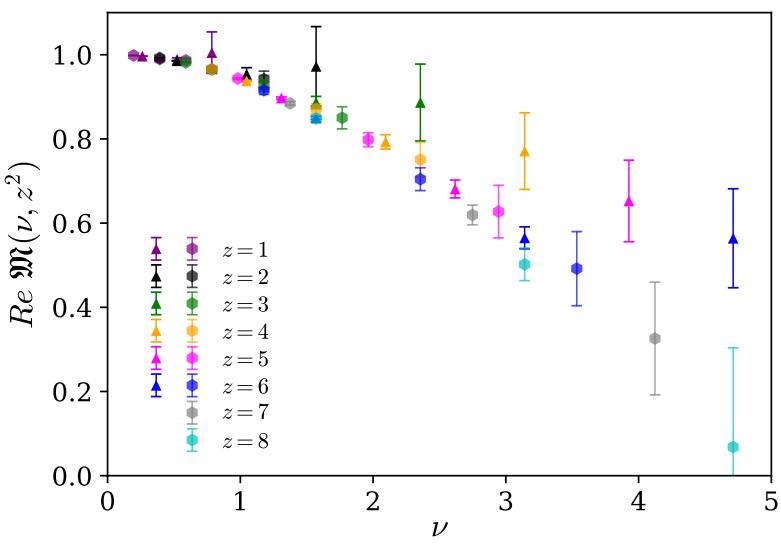
# Pseudo-PDF Approach

ID	$a$ (fm)	$m_\pi$ (MeV)	$\beta$	$am_t$	$am_s$	$L^3 \times T$	$N_{\text{cfg}}$
al27m415	0.127(2)	415(23)	6.1	-0.280	-0.245	$24^3 \times 64$	2147
al27m415L	0.127(2)	415(23)	6.1	-0.280	-0.245	$32^3 \times 96$	2560

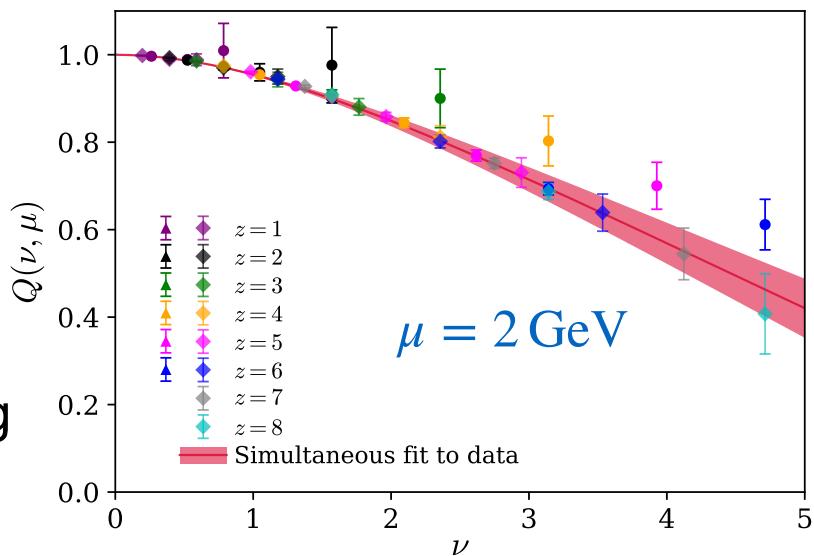
See Savvas Zeferopoulos, Thursday PM



Same ensemble as GLCS

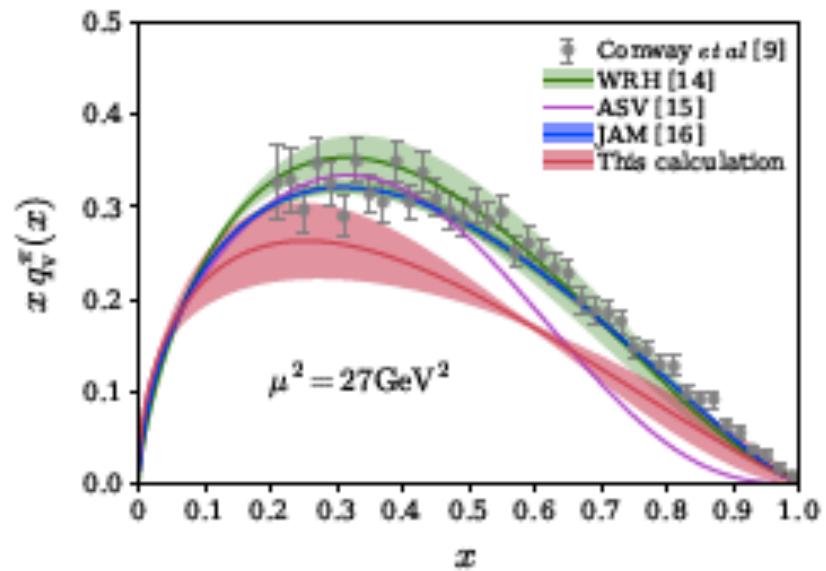
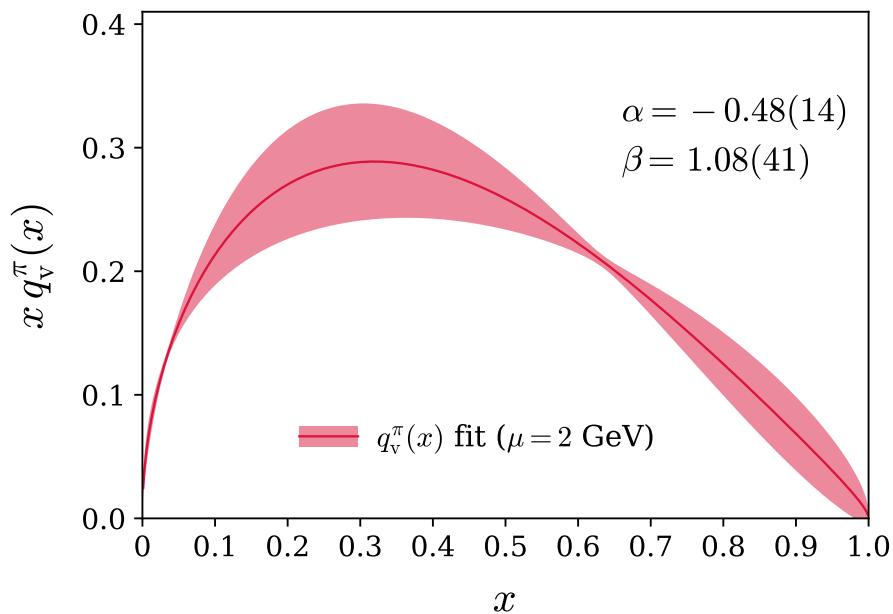


After  
Matching

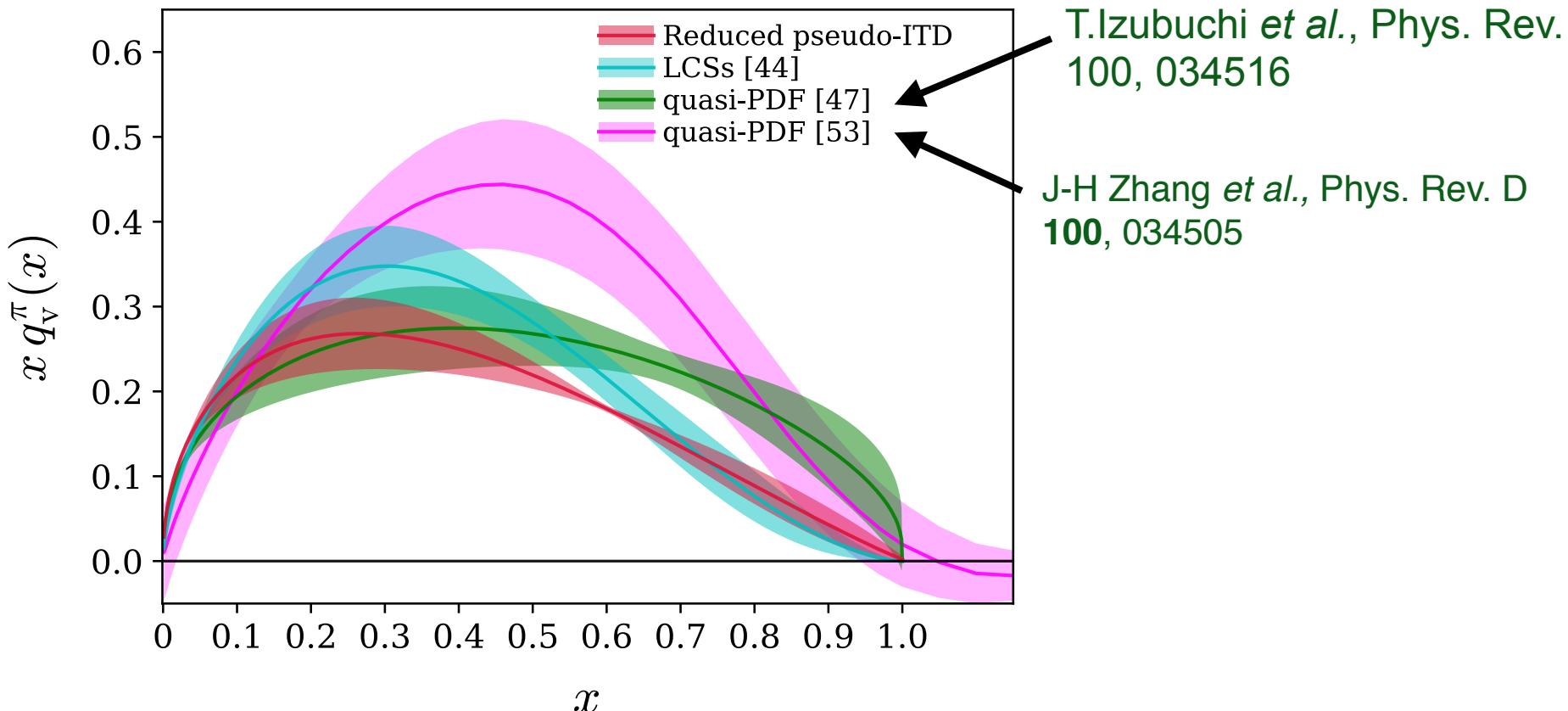


B.Joó, J.Karpie, K.Orginos, A.Radyushkin, D.Richards, R.Sufian, S.Zafeiropoulos, arXiv:1909.08517

# Pion pPDF - II



# Pion pPDF - III



# Nucleon pseudo-PDF

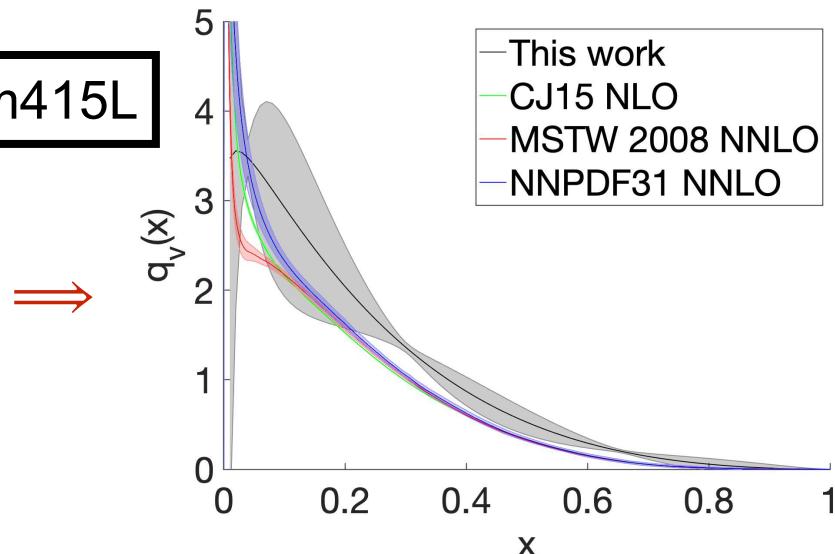
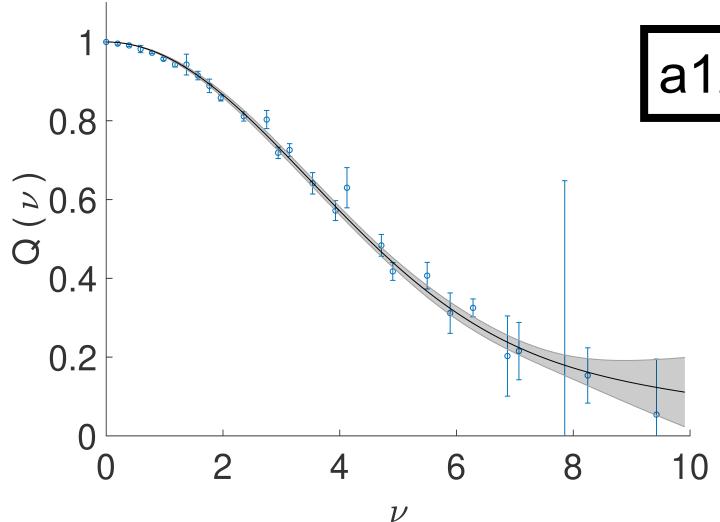
See Savvas Zeiferopoulos, Thursday PM

Ground-breaking quenched calculation: K. Orginos et al., PRD96 (2017), 094503

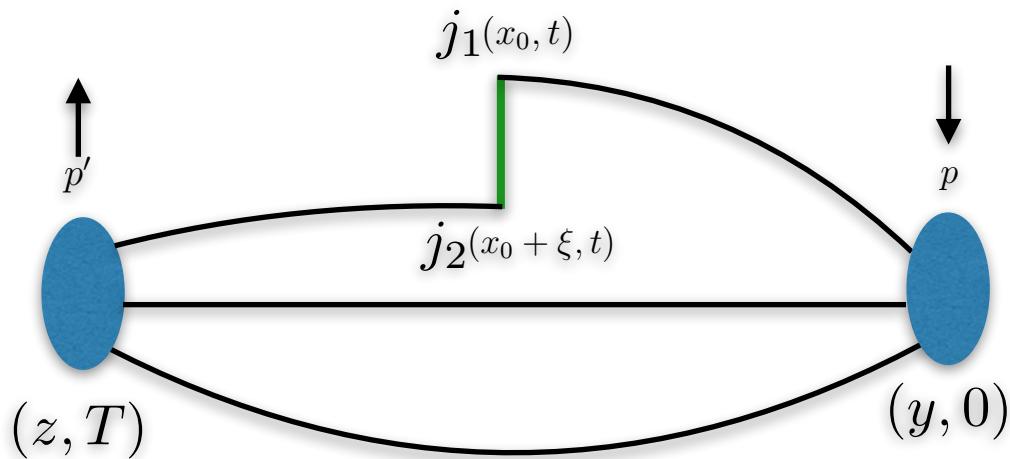
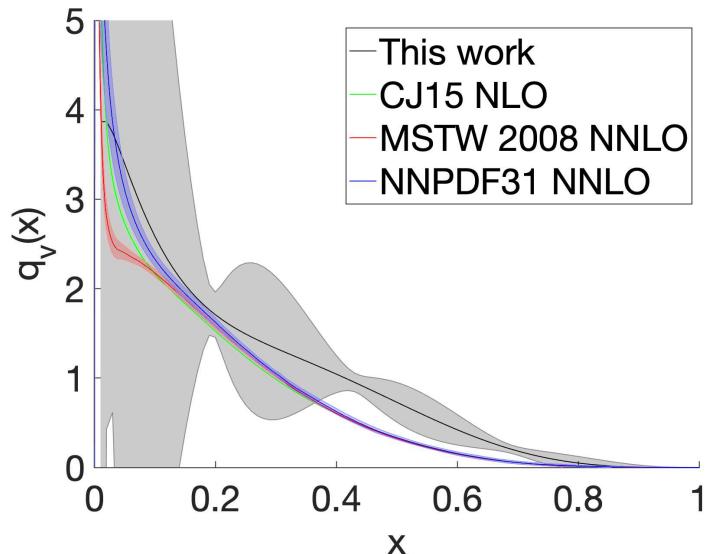
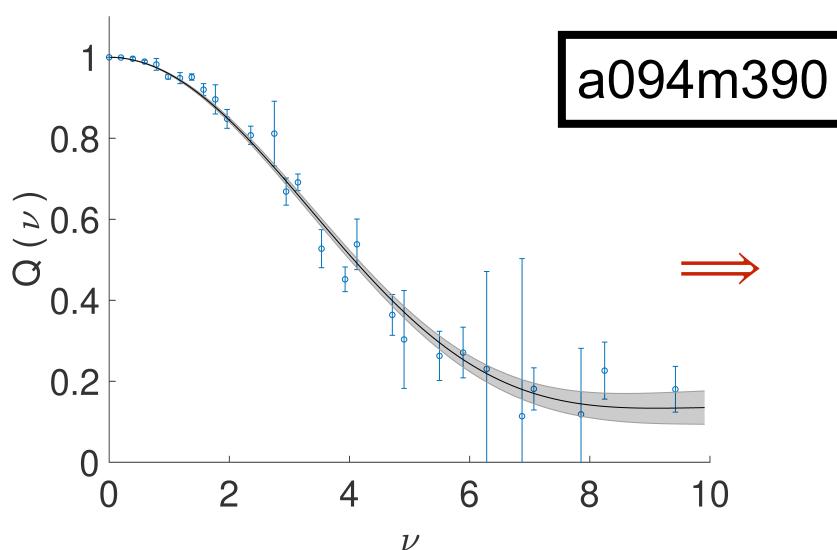
B.Joo et al., arXiv:1908.09771

ID	$a(\text{fm})$	$M_\pi(\text{MeV})$	$\beta$	$c_{\text{SW}}$	$am_\ell$	$am_s$	$L^3 \times T$	$N_{cfg}$
$a127m415$	0.127(2)	415(23)	6.1	1.24930971	-0.2800	-0.2450	$24^3 \times 64$	2147
$a127m415L$	0.127(2)	415(23)	6.1	1.24930971	-0.2800	-0.2450	$32^3 \times 96$	2560
$a094m390$	0.094(1)	390(71)	6.3	1.20536588	-0.2350	-0.2050	$32^3 \times 64$	417

← Finer lattice spacing



# Nucleon GLCS?



# Summary

- Revolution in the study of  $x$ -dependent measures of hadron structure
  - Impact global fitting community? Unclear for valence PDFs BUT
  - First-Principles calculation
- Pseudo-PDF/GLCS approach has a well-defined short-distance scale: factorize short-distance physics from perturbative scale.
- Solution of inverse problem: common to all attempts to extract PDFs. Appeal to global fitting community
- Systematic study of pion PDF using GLCS approach is in preparation; NLO perturbative kernel. **Q. Ma**
- To control systematics
  - fine lattices - to ensure in perturbative regime
  - large momenta - to provide range in Ioffe time

Extension to 3D imaging through GPDs and TMDs, see Michael Engelhardt - *opportunity to predict and constrain experiment*

# Moments of Pion PDF

