Acknowledgements: The PREX, CREX, E158 and MOLLER Collaboration The Accelerator Divisions at SLAC & Jefferson Laboratory V. Cirigliano, J. Erler, C. Horowitz, W. Marciano, M. Ramsey-Musolf, J. Piekerewicz

The Parity-Violation Program at Jefferson Lab

Krishna S. Kumar UMass, Amherst

EINN2019 Paphos, Cyprus November 2, 2019



Introduction to Parity-Violating Electron Scattering

- Relativistic electron scattering and nuclear size
- Neutral weak interactions
- 0
 - **PREX Experimental Overview**
 - First Result from 2010 Run
 - Summer 2019 PREX-II Run and upcoming CREX run

Scattering off electrons, protons and quarks

- - The past: E1 58, Qweak. The future: MOLLER, SOLID, P2
- Novel probes of partonic structure
- **Neutral Current Structure Functions at an EIC**
- **Summary and Outlook**

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The Parity Violation Program at Jefferson Lab

Outline

Historical introduction for graduate students

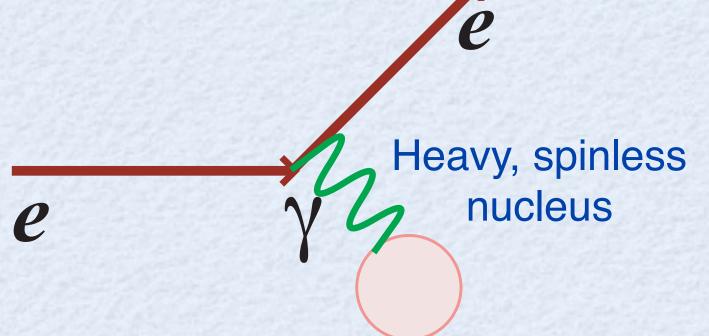
Elastic Scattering off a Heavy Doubly Magic Nucleus

PRL 108 (2012) 112502 PRC 85 (2012) 032501

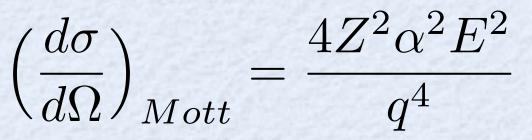
Program to Search for New Neutral Current Interactions Beyond the Standard Model



Relativistic Electron Scattering and nuclear size



Differential Cross Section

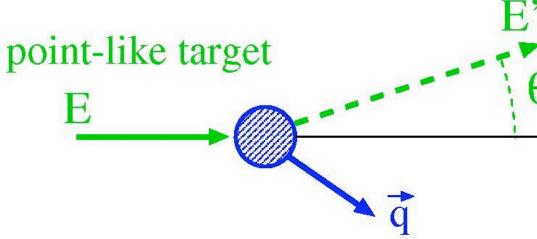


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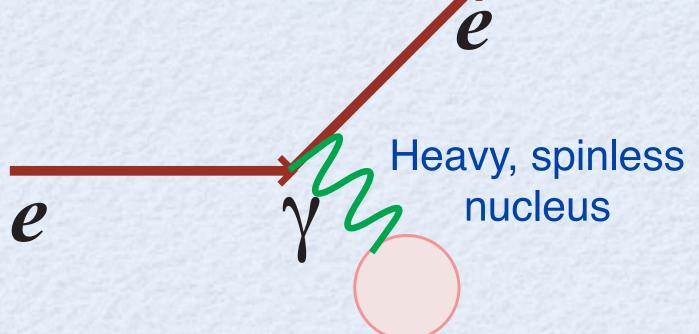
4-momentum transfer

 $q^{2} = -4 EE' \sin^{2} \frac{\theta}{2}$ $Q \approx \frac{hc}{\lambda}$ 4-momentum transf Q^{2}: -(4-momentum)^{2}
of the virtual photon

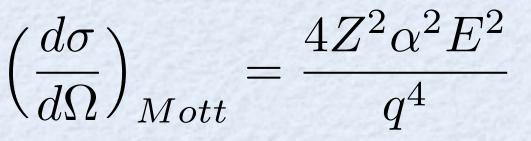




Relativistic Electron Scattering and nuclear size



Differential Cross Section



$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{\rm Mott} \left|F(q)\right|^2$$

As Q increases, nuclear size modifies formula

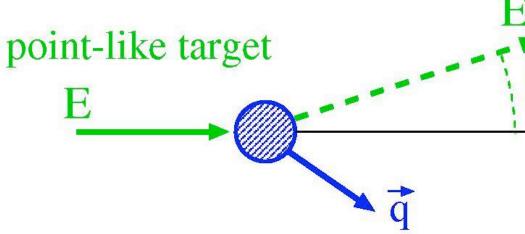
Neglecting recoil, form factor F(q) is the **Fourier transform of charge distribution**

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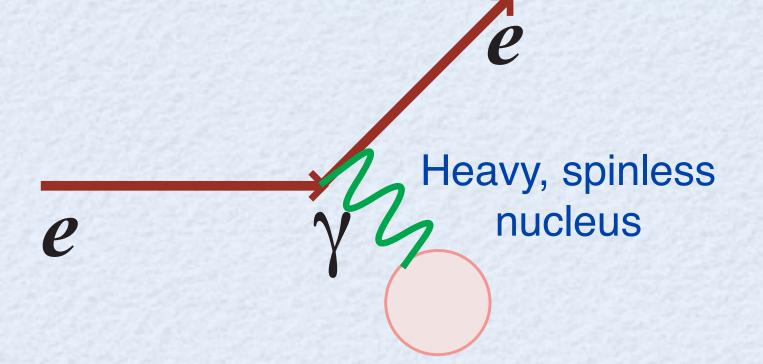
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4-momentum transfer

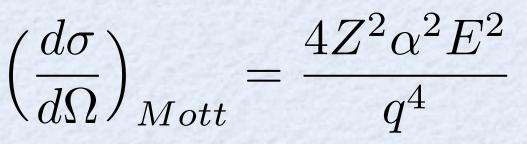
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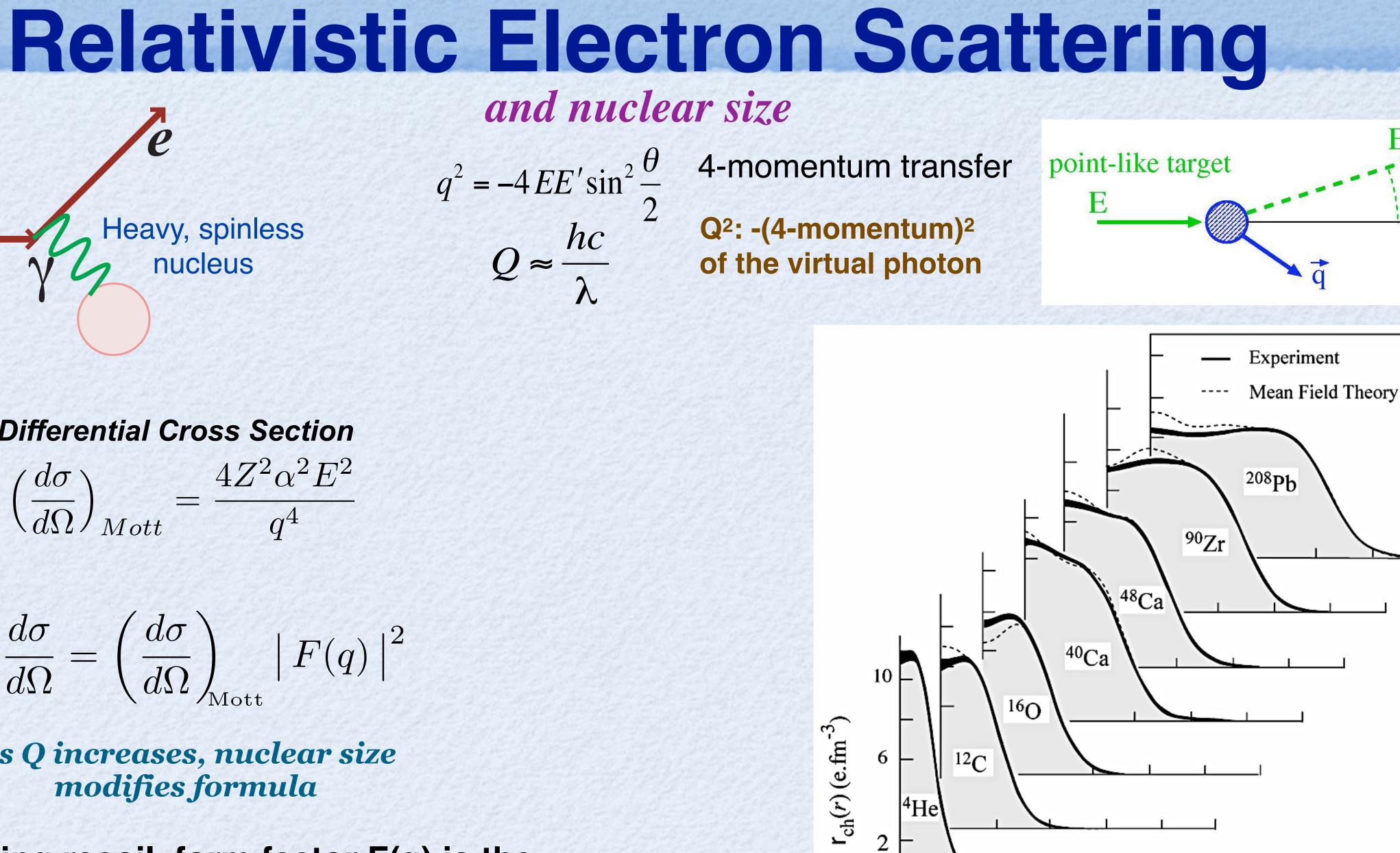
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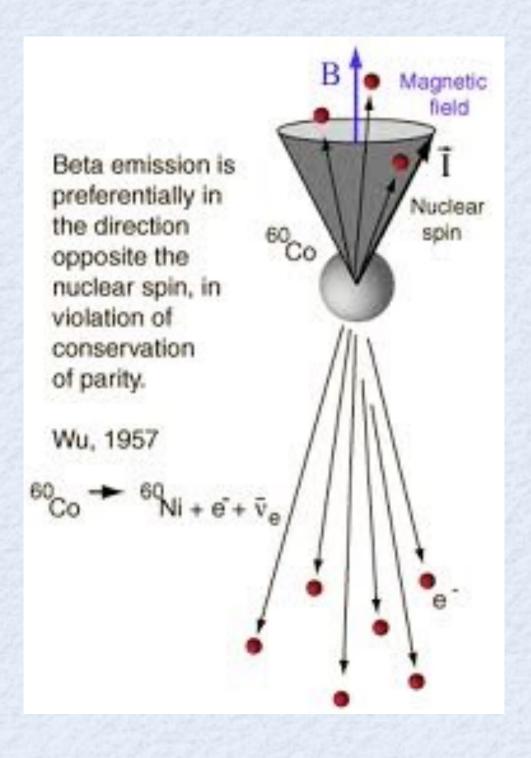
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 $r(\mathrm{fm})$

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1958: Zel'dovich speculation: Is Electron Scattering Parity-Violating? Nuclear B Decay

> V n

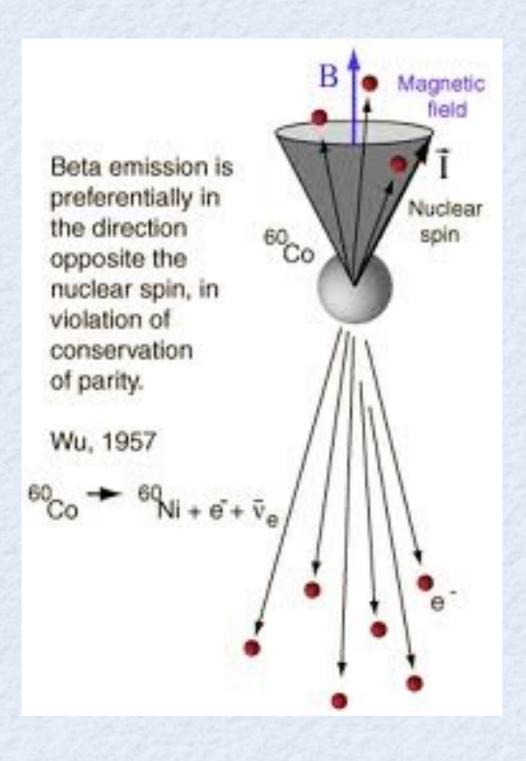
charge and flavor-changing

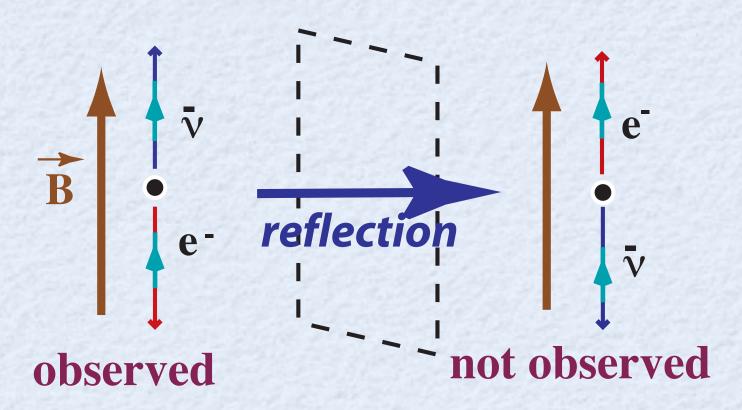
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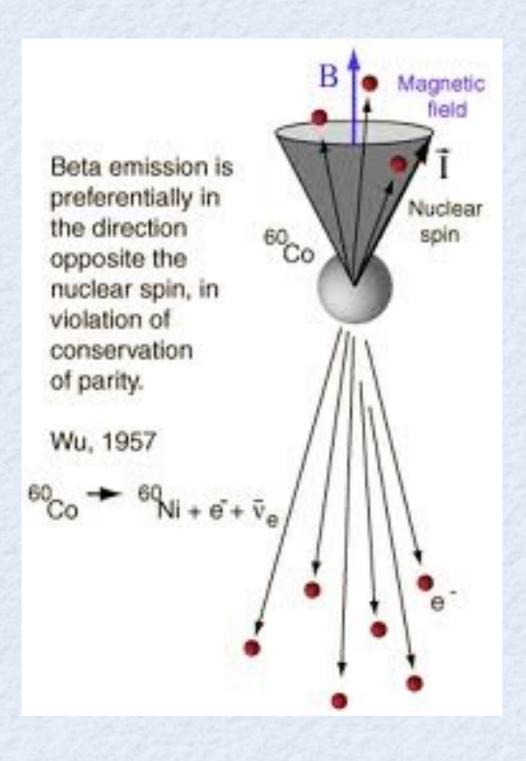


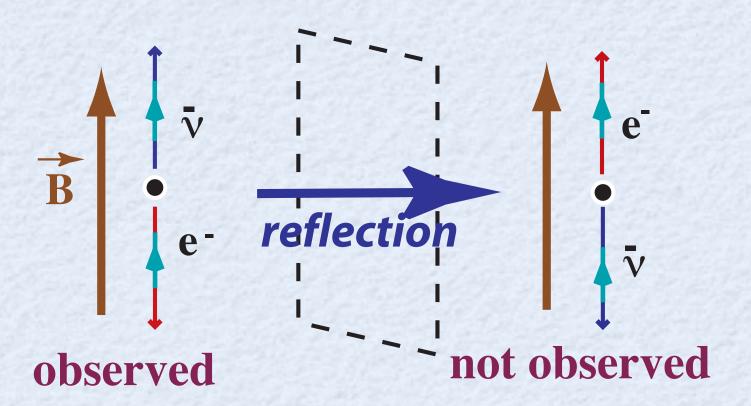
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1958: Zel'dovich speculation: Is Electron Scattering Parity-Violating?





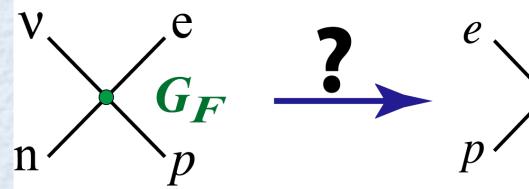


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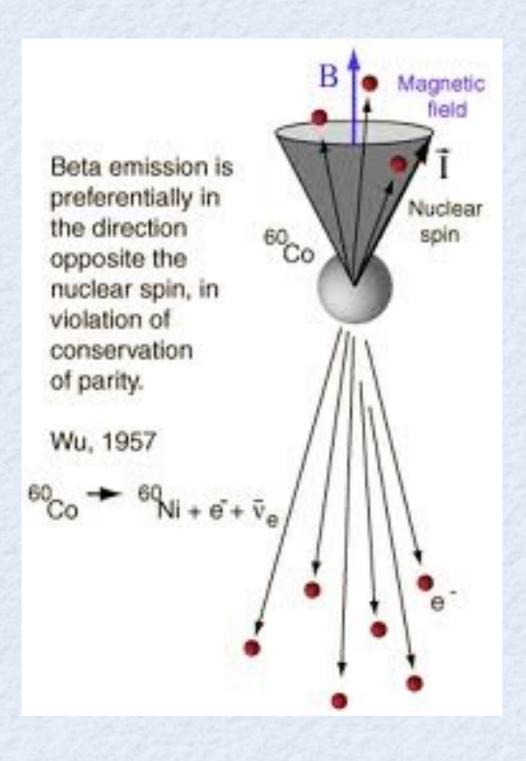
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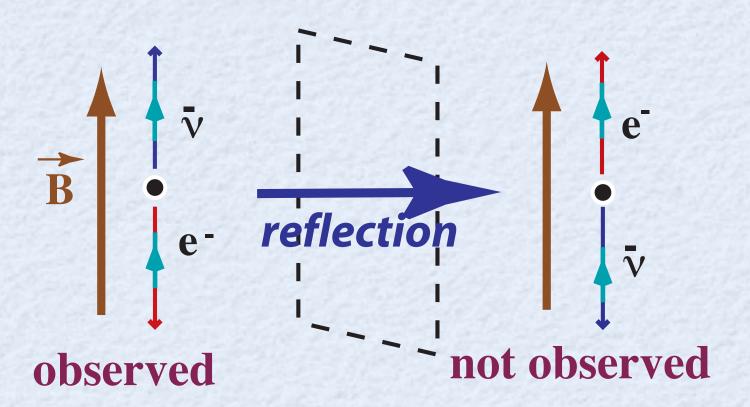
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Neutron β *Decay*







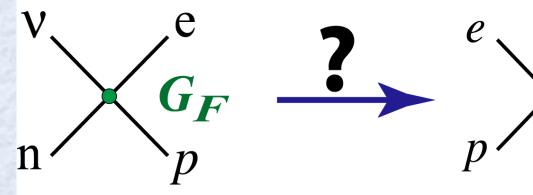


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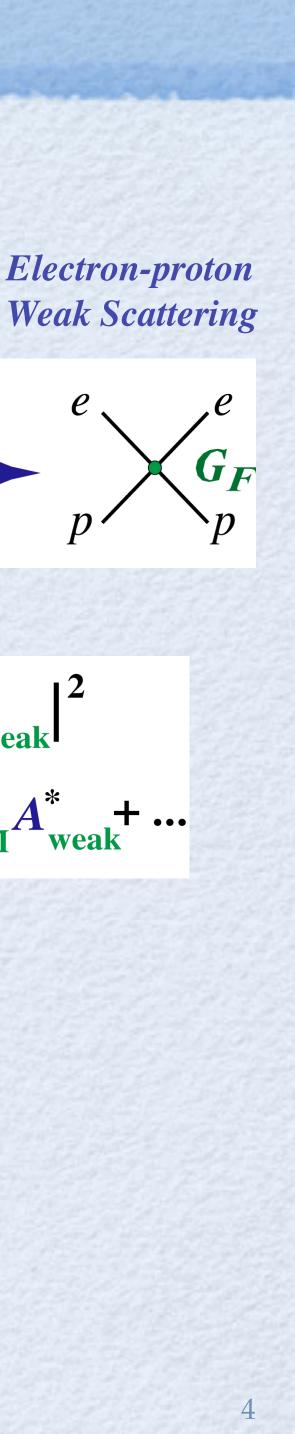
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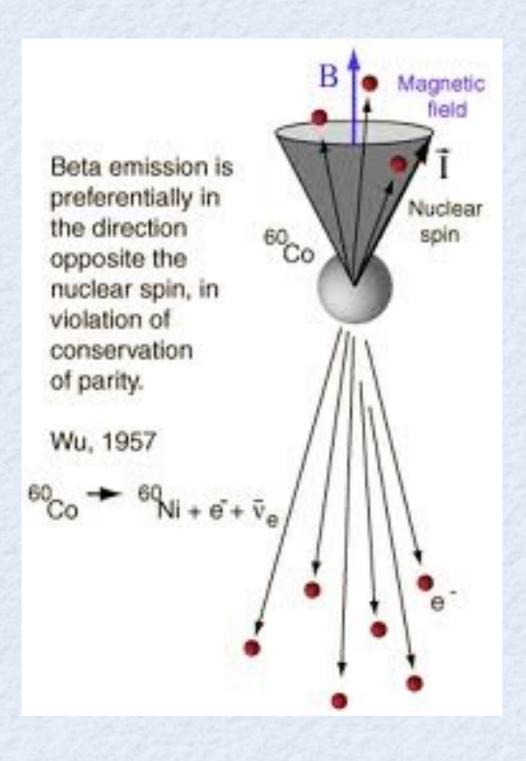
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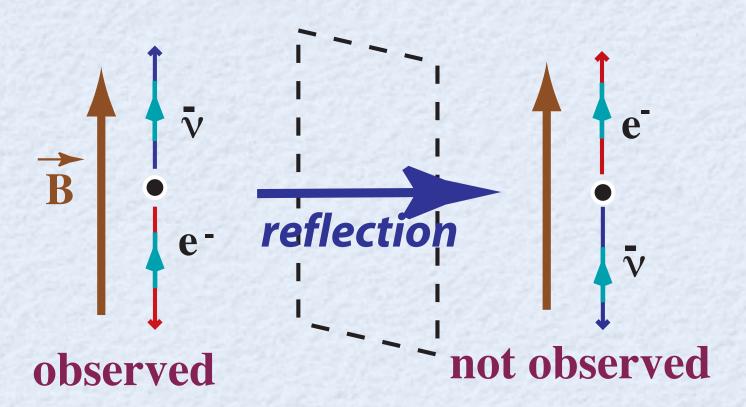
Neutron *β* Decay



 $\sigma \alpha |A_{EM} + A_{weak}|^2$ ~ $|A_{EM}|^2 + 2A_{EM}A_{weak}^* + \dots$







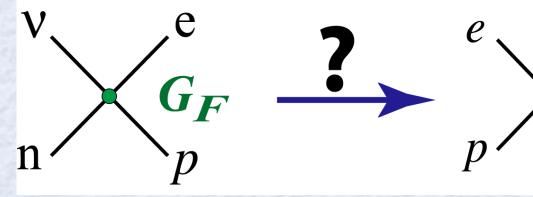
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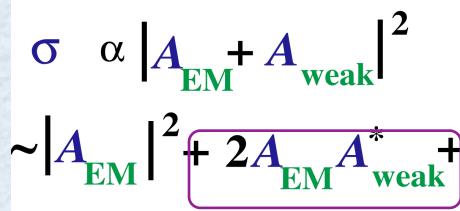
The Parity Violation Program At Jefferson Lab

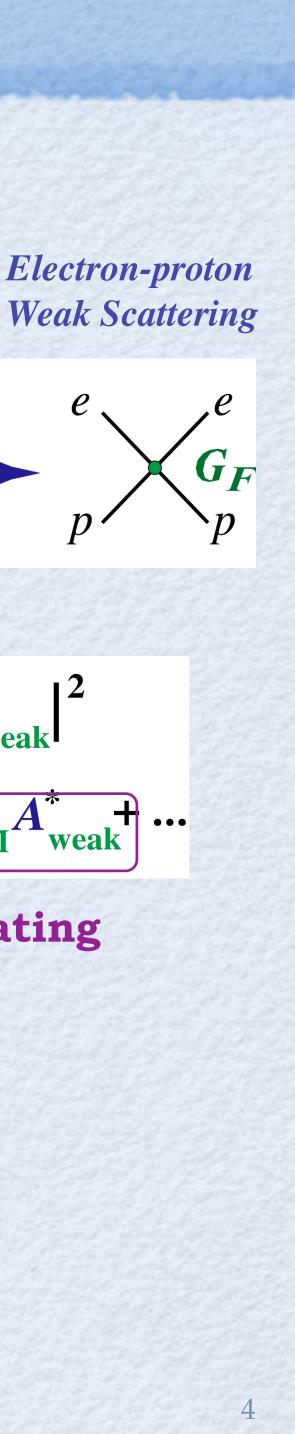
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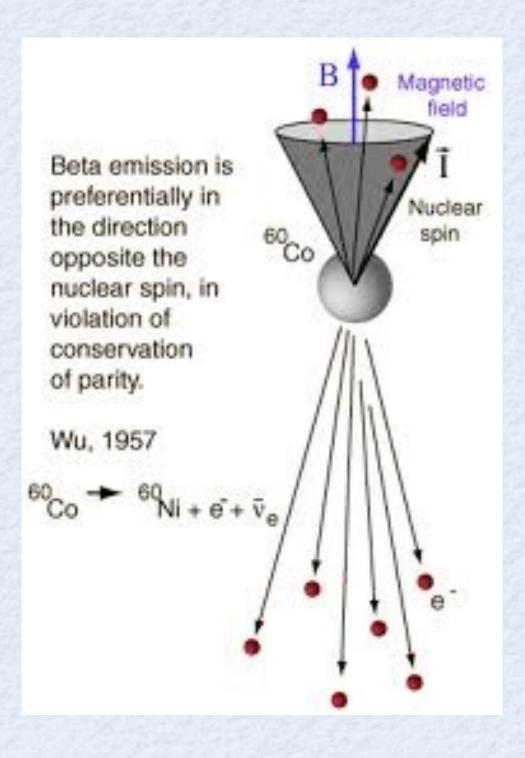
Neutron *β* Decay

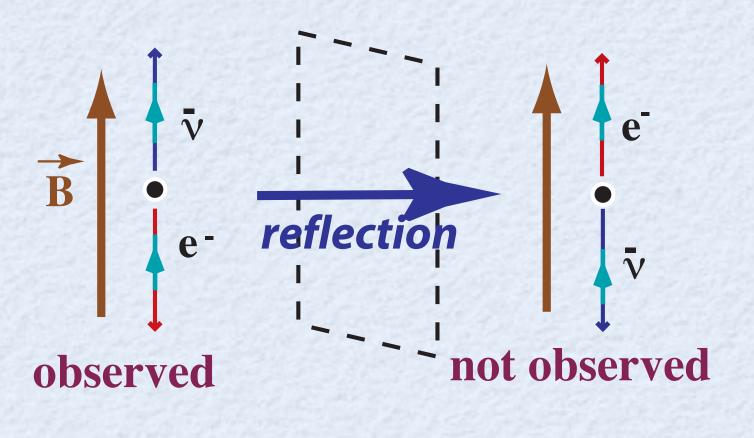
Weak Scattering











 $\frac{\sigma_{\bullet} - \sigma_{\bullet}}{\sigma_{\bullet} + \sigma_{\bullet}} = -A$

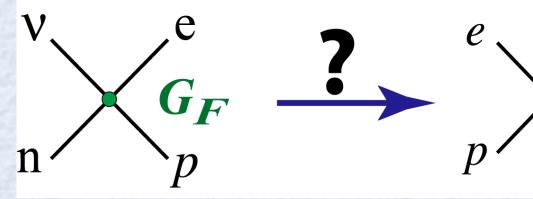
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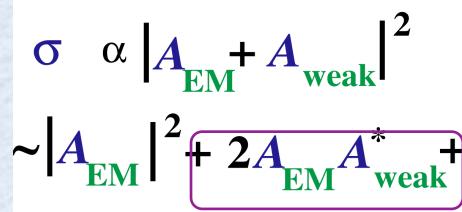
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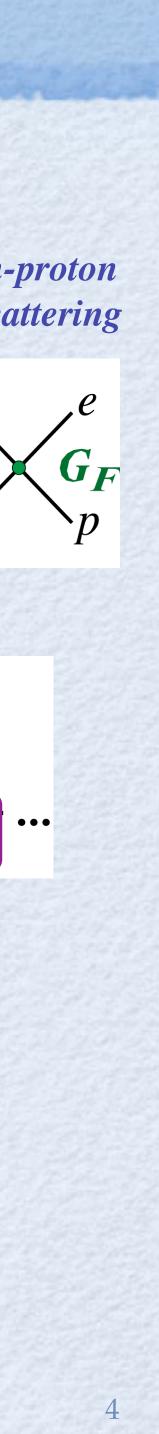
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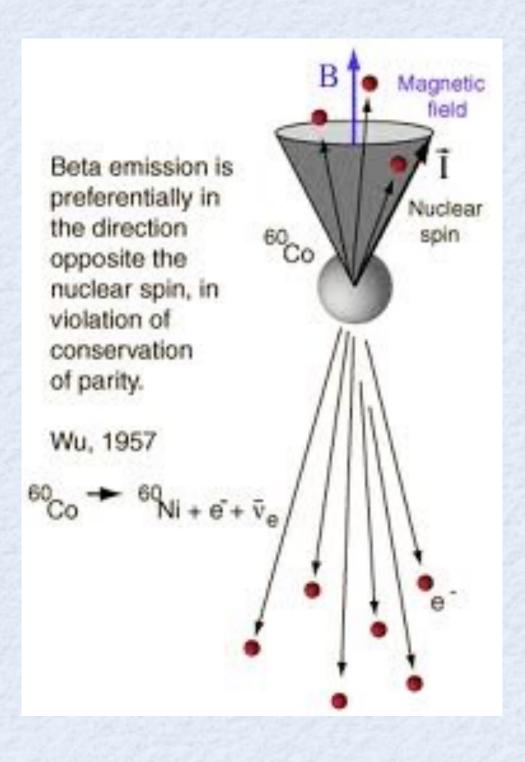
Neutron *B* Decay

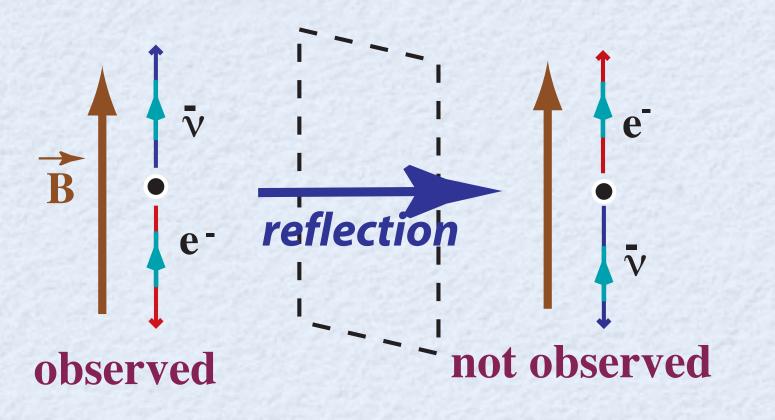
Electron-proton Weak Scattering











 $= \frac{\sigma_{\downarrow} - \sigma_{\downarrow}}{\sigma_{\downarrow} + \sigma_{\downarrow}} = -A \sim \frac{A_{\text{weak}}}{A_{\text{EM}}} \sim \frac{G_F Q^2}{4 \pi \alpha}$ $A_{PV} \sim 10^{-4} \cdot Q^2 \,(\text{GeV}^2)$

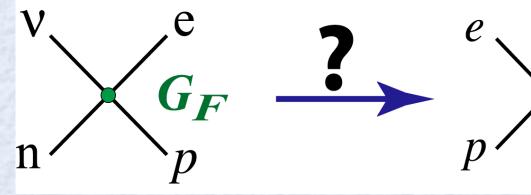
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1958: Zel'dovich speculation: Is Electron Scattering Parity-Violating?

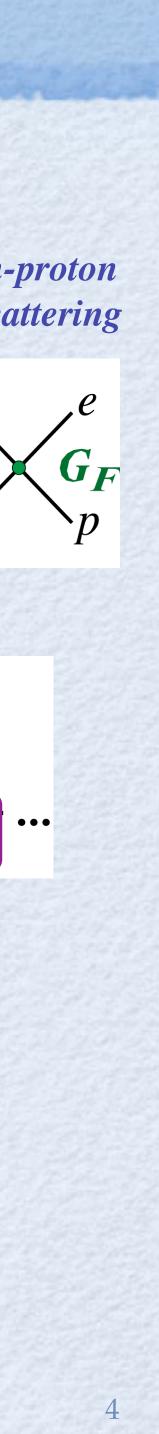
Neutron β *Decay*

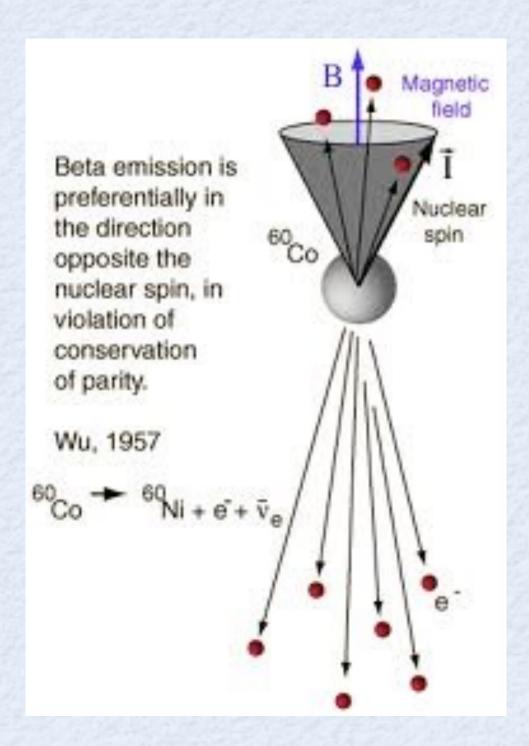
Electron-proton Weak Scattering

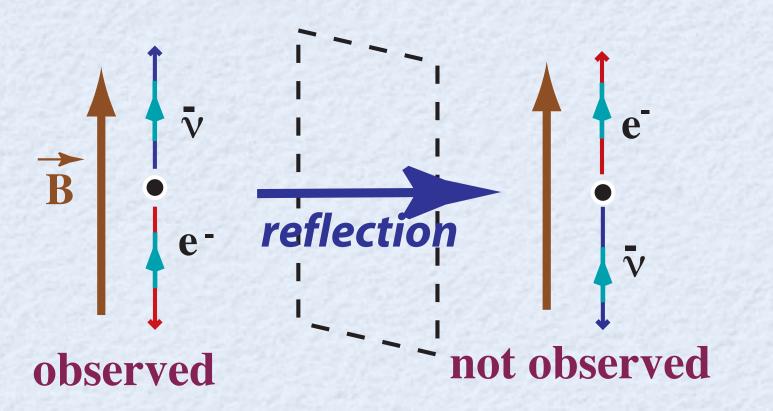


$$\sigma \alpha |A_{EM} + A_{weak}|^2$$

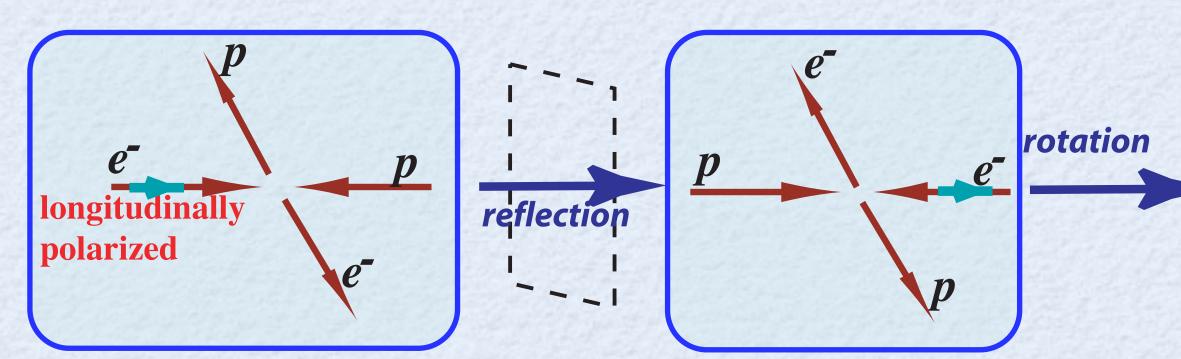
~ $|A_{EM}|^2 + 2A_{EM}A_{weak}^* +$







 $A_{PV} \sim$



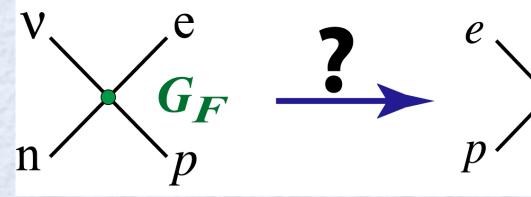
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1958: Zel'dovich speculation: Is Electron Scattering Parity-Violating?

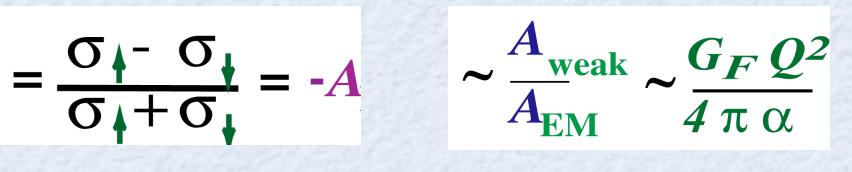
Neutron β *Decay*

Electron-proton Weak Scattering

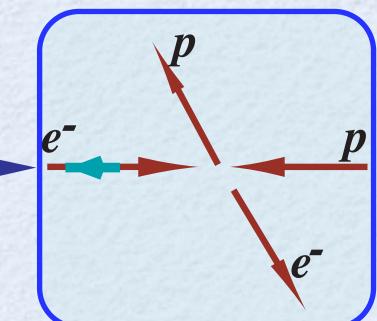


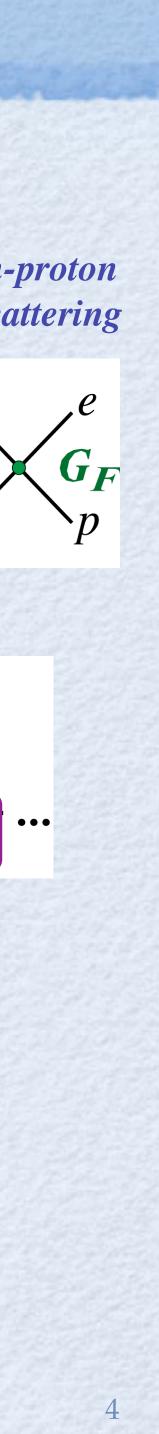
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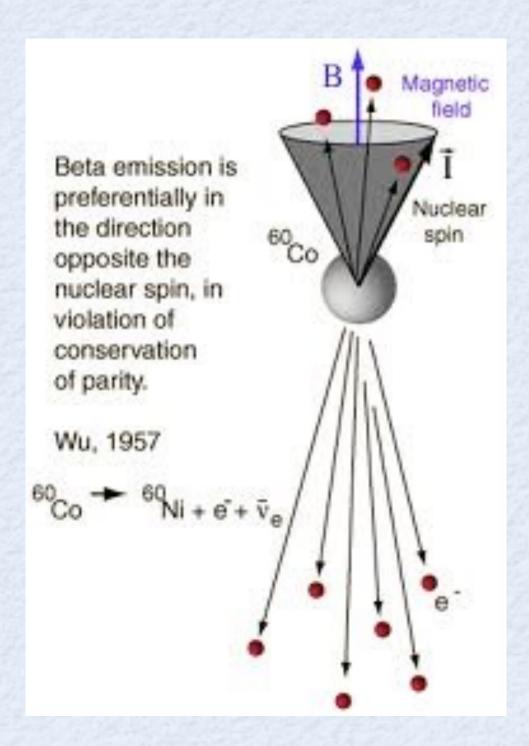
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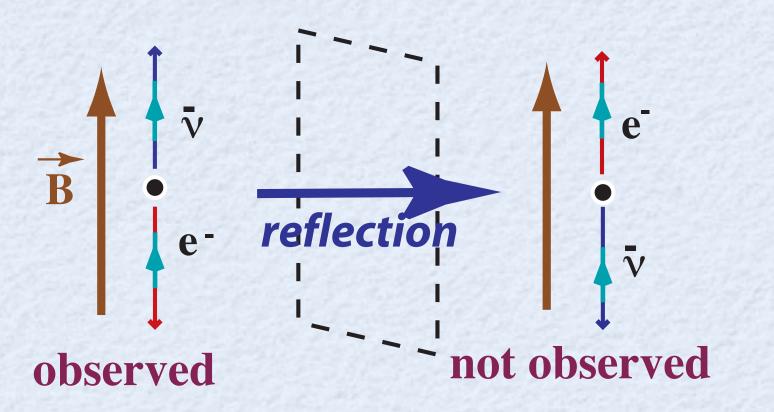


$$10^{-4} \cdot Q^2$$
 (GeV²)

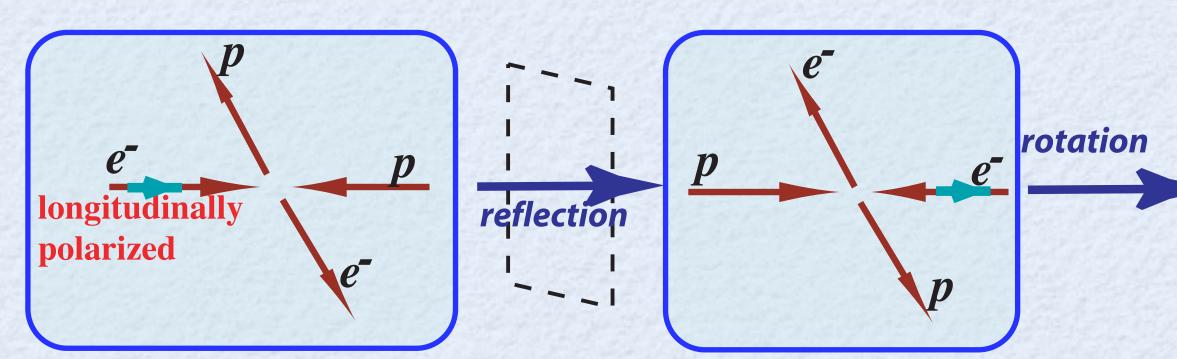








 $A_{PV} \sim$

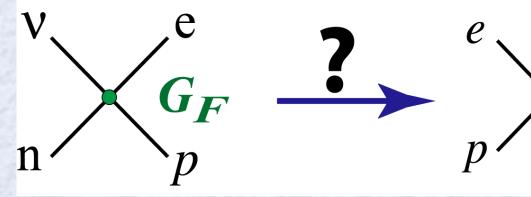


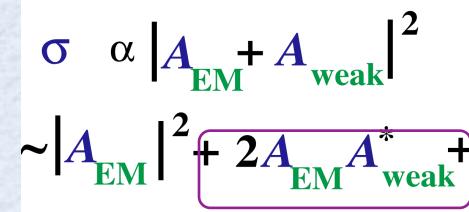
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1958: Zel'dovich speculation: Is Electron Scattering Parity-Violating?

Neutron β *Decay*



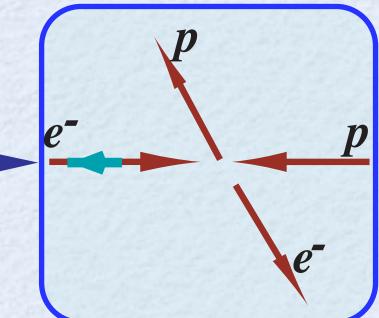


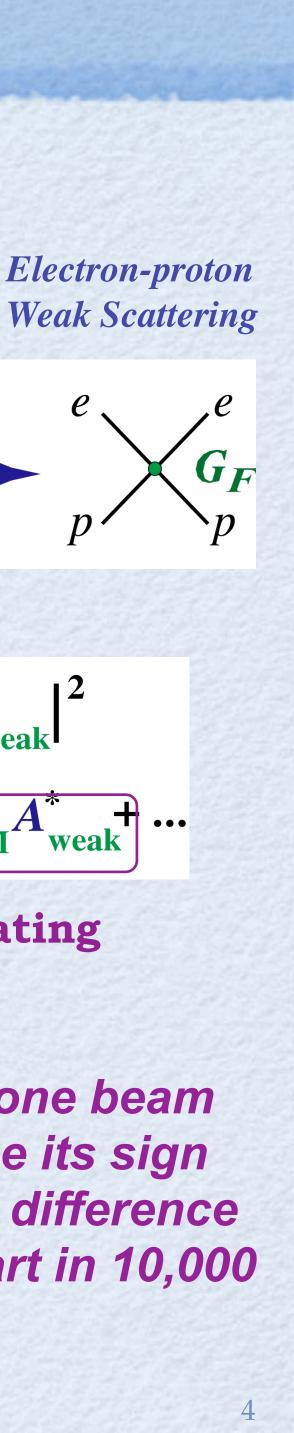
Parity-violating

 Iongitudinally polarize one beam with the ability to change its sign Measure fractional rate difference with a sensitivity of a part in 10,000

 $= \frac{\sigma_{\downarrow} - \sigma_{\downarrow}}{\sigma_{\downarrow} + \sigma_{\downarrow}} = -A \qquad \sim \frac{A_{\text{weak}}}{A_{\text{EM}}} \sim \frac{G_F Q^2}{4 \pi \alpha}$

$$10^{-4} \cdot Q^2$$
 (GeV²)





The Z boson incorporated

Left-	Right-	
$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	γ Charge
$T = \pm \frac{1}{2}$	zero	W Charg
$T - q\sin^2\theta_W$	$-q\sin^2\theta_W$	Z Charg

One free parameter: weak mixing angle θ_W

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Neutral Weak Interaction Theory

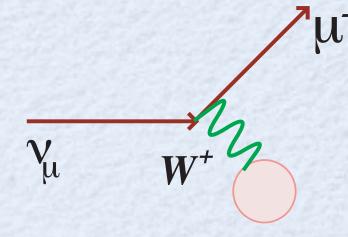
Glashow, Weinberg and Salam: $SU(2)_L X U(1)_Y$

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One free parameter: weak mixing angle θ_W

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Charged Current

The Parity Violation Program At Jefferson Lab

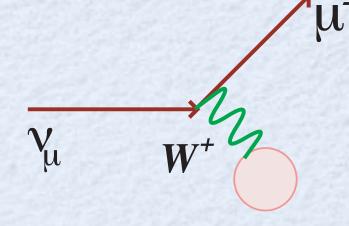
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	1.64.5
	1.47.4
1.19	11.56
1100	181.5.5
	1.6.5
6036	1.5
	1000
ge	100
50	1.1.1
-	1000
^	



Charged Current

 v_{μ} Z^0

Neutral Current

The Z boson incorporated

Left-	Right-	
$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	γ Charge
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Do lepton-nucleon neutral current interactions exhibit parity violation?

 $(e)_r$

or

Weinberg model Parity is violated

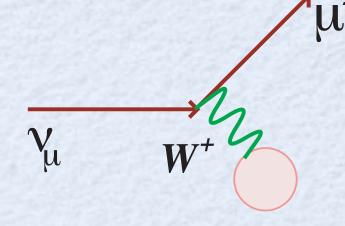
$$A_{PV} \sim 10^{-4}$$

Parity is conserved

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One free parameter: weak mixing angle θ_W

	1.64.5
	1.47.4
1000	11.56
1100	1815.5
	1.6.5
6036	1.5
	1000
ge	100
50	1.1.1
-	1000
^	



Charged Current

70

Neutral Current

Neutral Weak Interaction Theory Glashow, Weinberg and Salam: SU(2)_LX U(1)_Y

The Z boson incorporated

Left-	Right-	
$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	γ Charge
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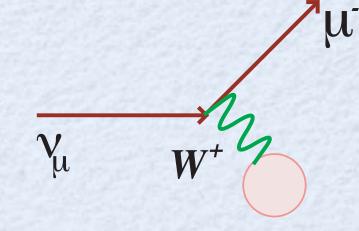
First table-top atomic parity violation searches: negative!

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One free parameter: weak mixing angle θ_W

,	
ge	
ge	12
e	6.50



Charged Current

70

Neutral Current

The Z boson incorporated

Left-	Right-	
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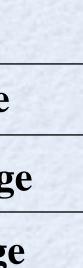
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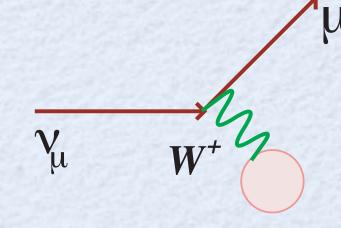
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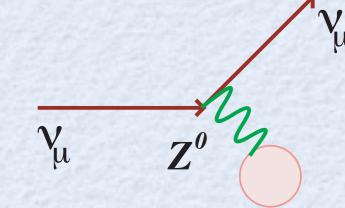
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One free parameter: weak mixing angle θ_W



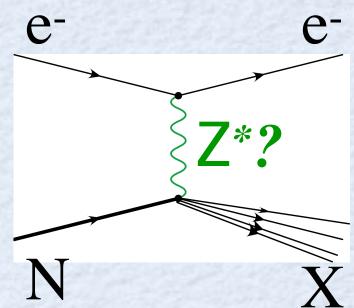


Charged Current



Neutral Current

electron-nucleon deep inelastic scattering



pressing problem in mid-70's

The Z boson incorporated

Left-	Right-	
$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	γ Charge
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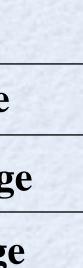
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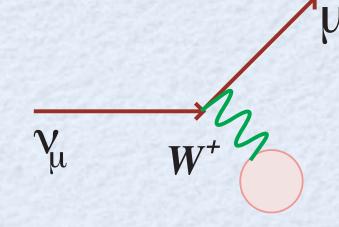
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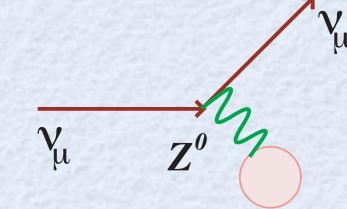
The Parity Violation Program At Jefferson Lab

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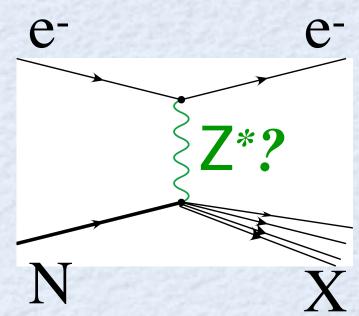


Charged Current



Neutral Current

electron-nucleon deep inelastic scattering



pressing problem in mid-70's

Seminal Experimental Measurement: E122 at the Stanford Linear Accelerator Center

The Z boson incorporated

Left-	Right-	
$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	γ Charge
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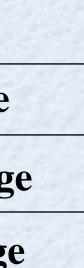
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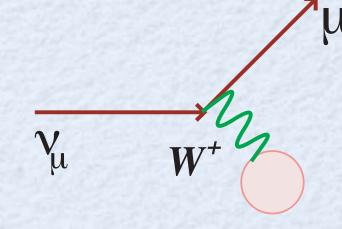
First table-top atomic parity violation searches: negative!

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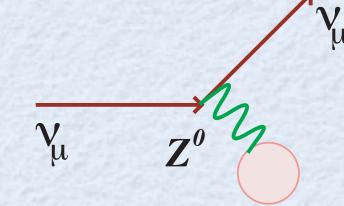
The Parity Violation Program At Jefferson Lab

One free parameter: weak mixing angle θ_W



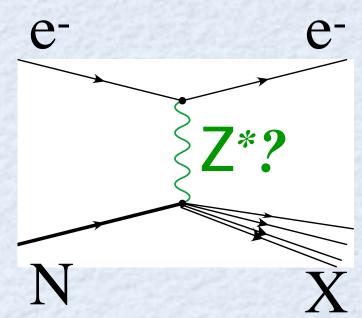


Charged Current



Neutral Current

electron-nucleon deep inelastic scattering



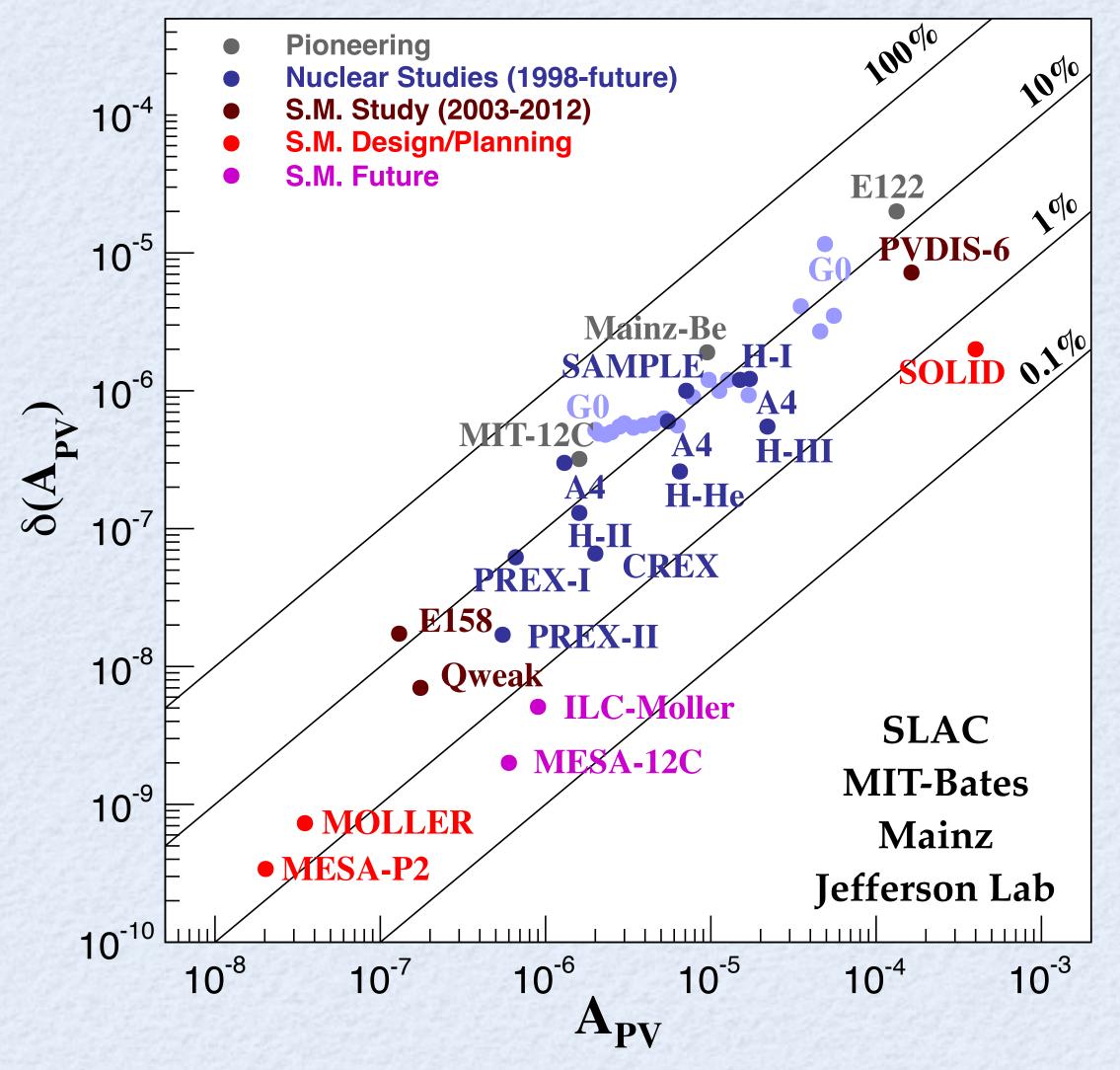
pressing problem in mid-70's

Seminal Experimental Measurement: E122 at the Stanford Linear Accelerator Center

•Parity Violation in Weak Neutral Current Interactions $\cdot \sin^2\theta_w = 0.224 \pm 0.020$: same as in neutrino scattering



4 Decades of Technical Progress



photocathodes, polarimetry, high power cryotargets, nanometer beam stability, precision beam diagnostics, low noise electronics, radiation hard detectors

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The Parity Violation Program At Jefferson Lab

Continuous interplay between probing hadron structure and electroweak physics

Parity-violating electron scattering has become a precision tool

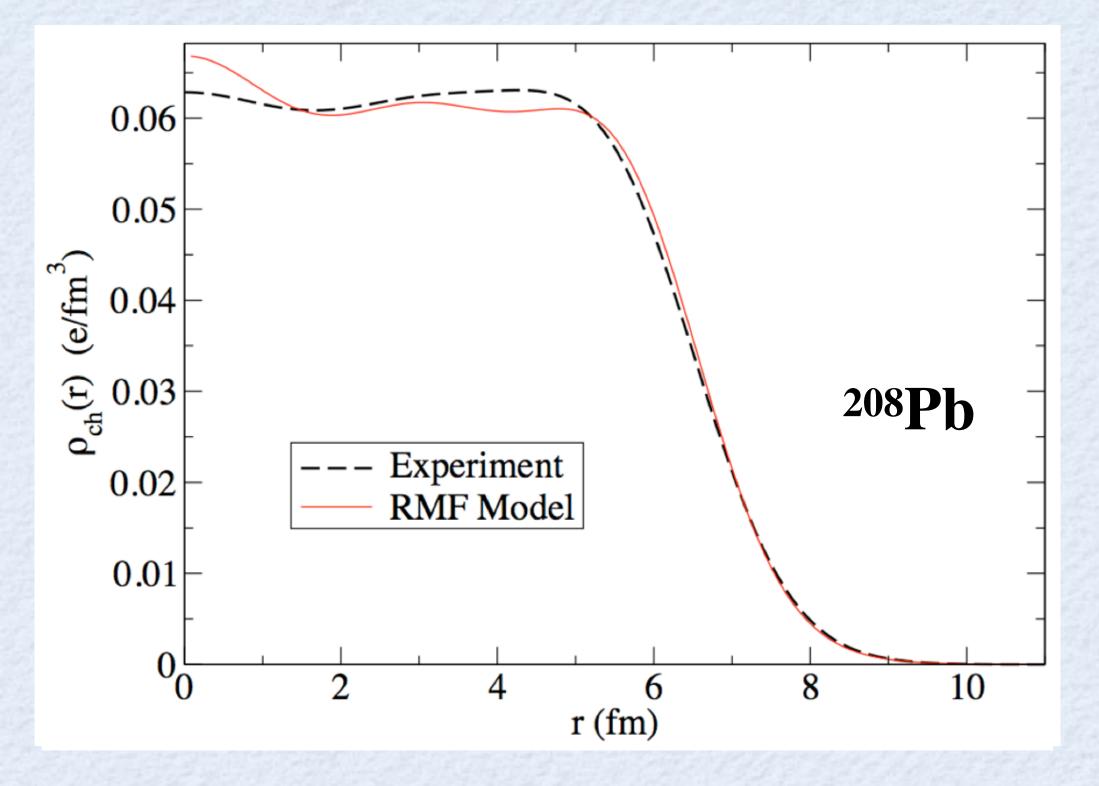
- Beyond Standard Model Searches
- Strange quark form factors
- Neutron skin of a heavy nucleus
- QCD structure of the nucleon

Mainz & MIT-Bates in the mid-80s JLab program launched in the mid-90s E158 at SLAC measured PV Møller scattering

State-of-the-art:

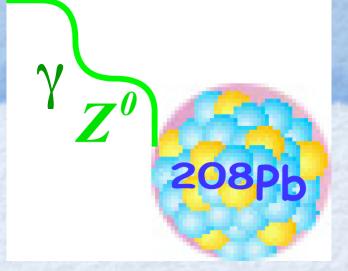
 sub-part per billion statistical reach and systematic control

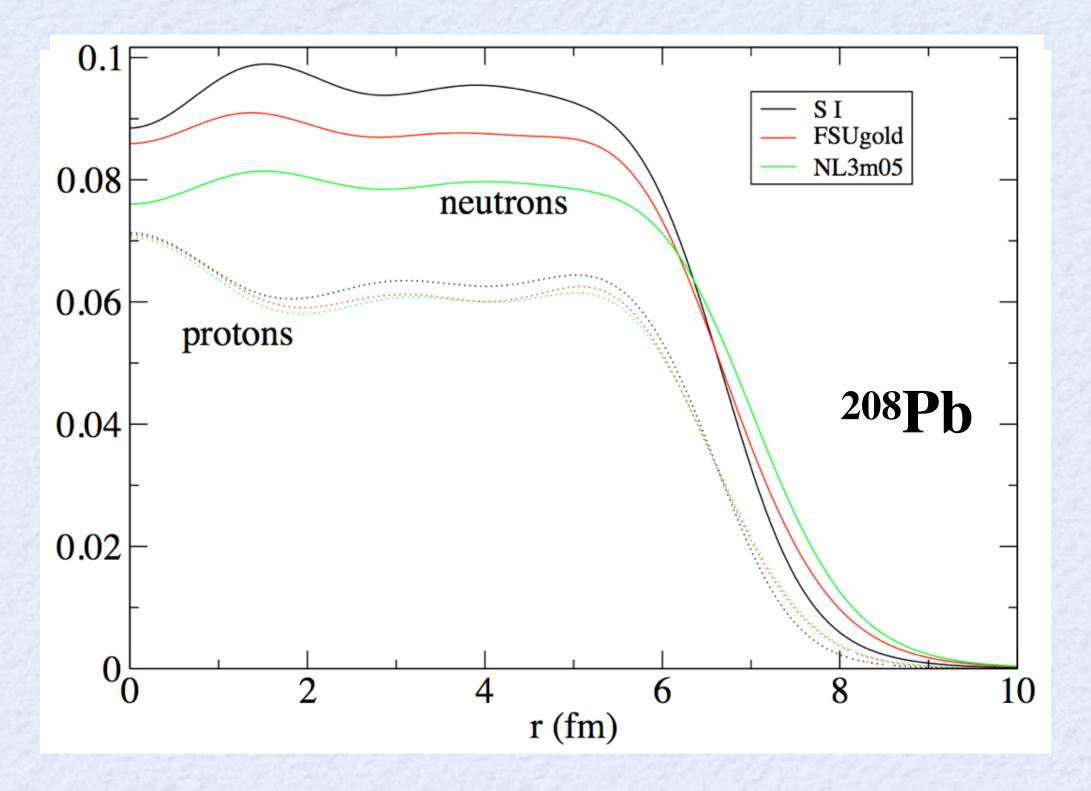
sub-1% normalization control



The Parity Violation Program At Jefferson Lab

PREX Concept

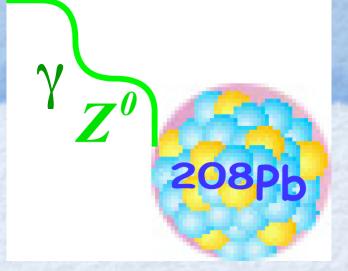




Krishna S. Kumar

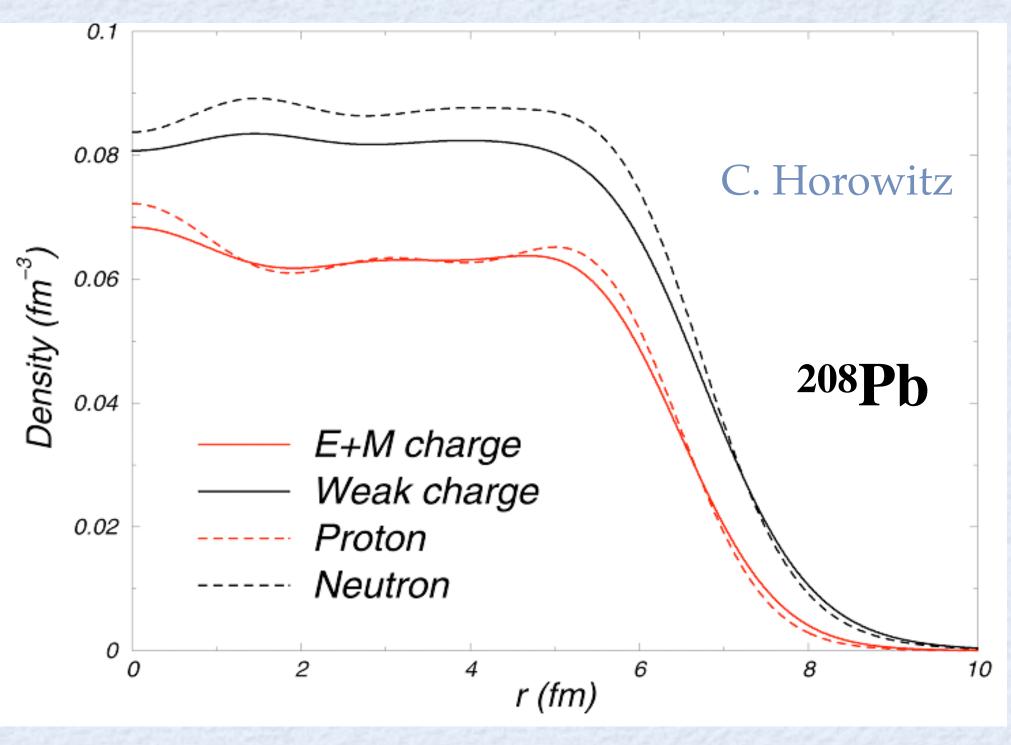
The Parity Violation Program At Jefferson Lab

PREX Concept



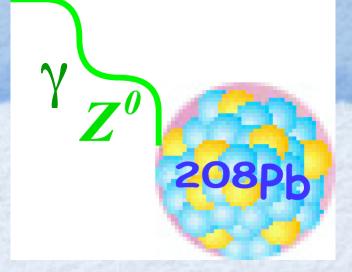
 $Q^n_W \sim -1$

Qⁿ_{EM} ~ 0 **Q**^{*p*}_{*EM*} ~ 1



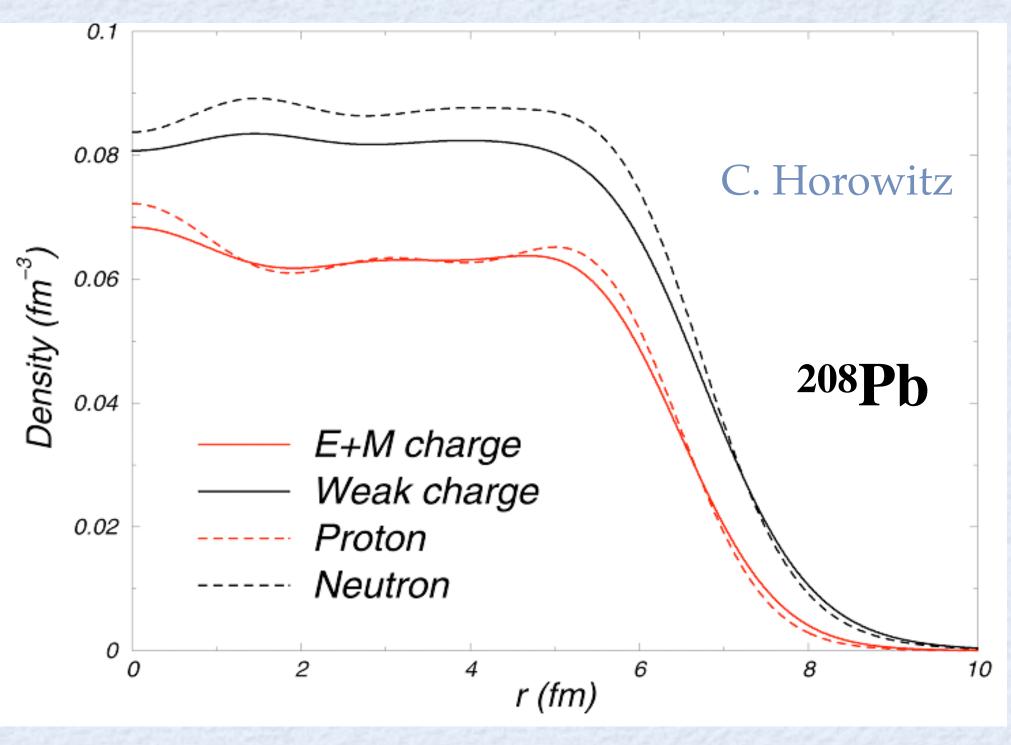
The Parity Violation Program At Jefferson Lab

PREX Concept $Q_W^P \sim 1 - 4 \sin^2 \theta_W$



 $Q^n_W \sim -1$

 $Q^n_{EM} \sim 0$ $Q^{p}_{EM} \sim 1$



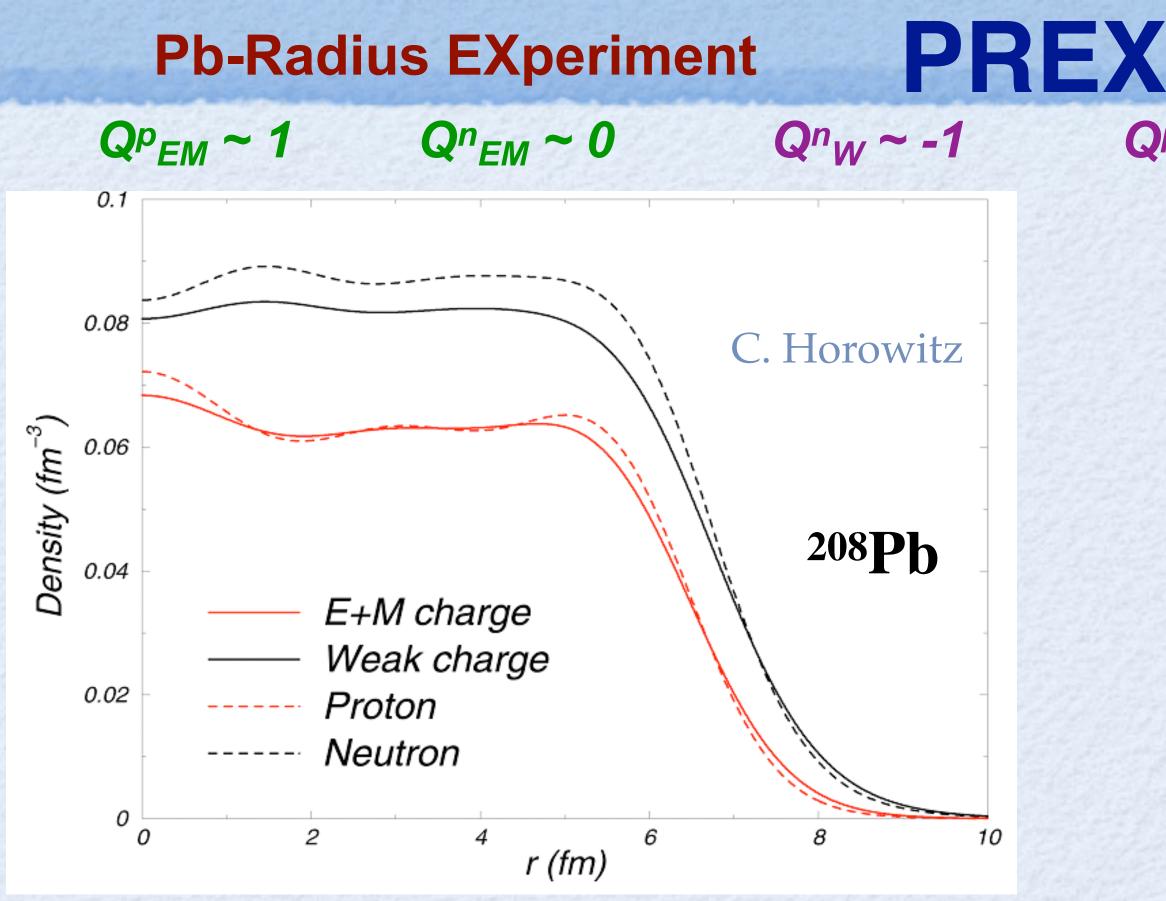
PREX Concept $Q^{p}_{W} \sim 1 - 4 \sin^2 \theta_{W}$



 $\mathcal{M}_{EM} \propto \frac{4\pi\alpha}{Q^2} F_{ch}(Q^2)$

 $\mathcal{M}_{Weak} \propto rac{G_F}{\sqrt{2}} \gamma_5 F_W(Q^2)$





F_{ch} and F_W: Functions of single nucleon form factors F_p and F_n

$$A_{PV} \approx \frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \frac{F_n(Q^2)}{F_p(Q^2)} + \dots$$

Small corrections involving electric form factors G_E (p,n,s)

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The Parity Violation Program At Jefferson Lab

PREX Concept $Q^{p}_{W} \sim 1 - 4 \sin^2 \theta_{W}$

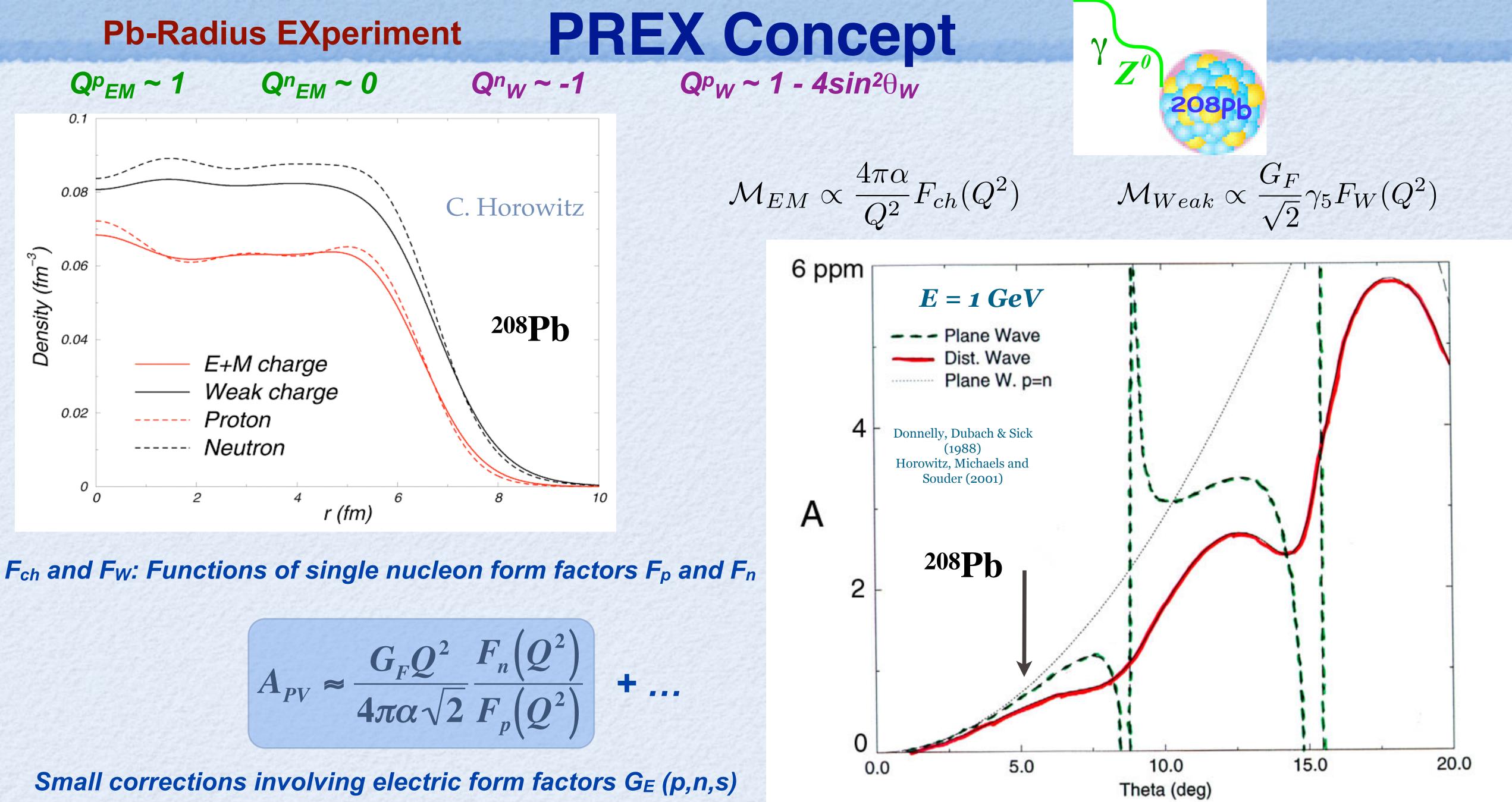
 \mathbf{N} 208D

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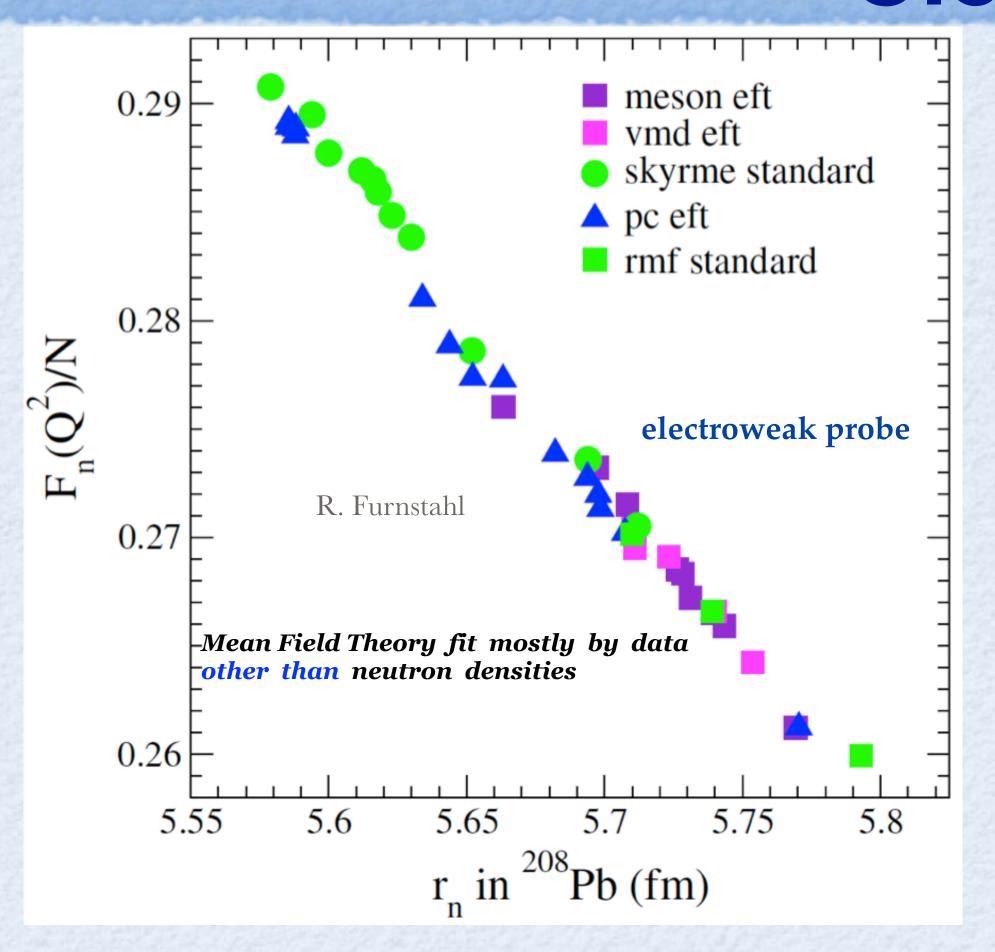




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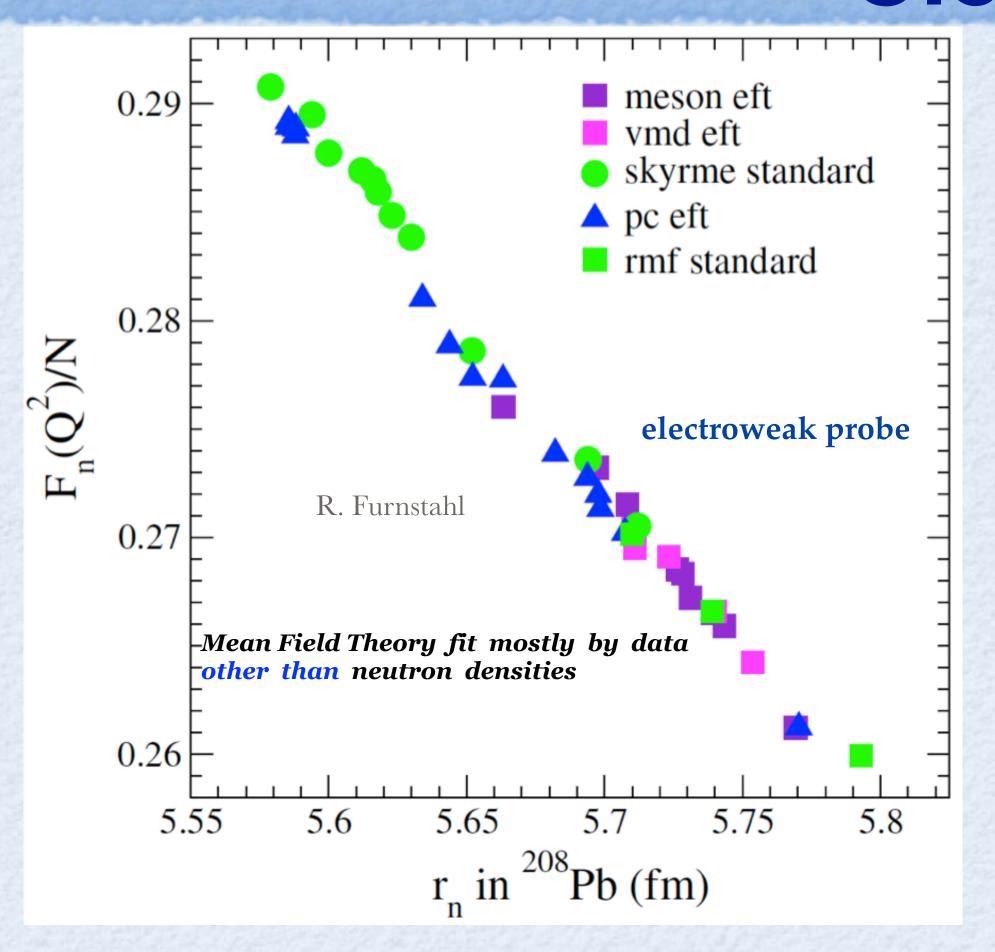
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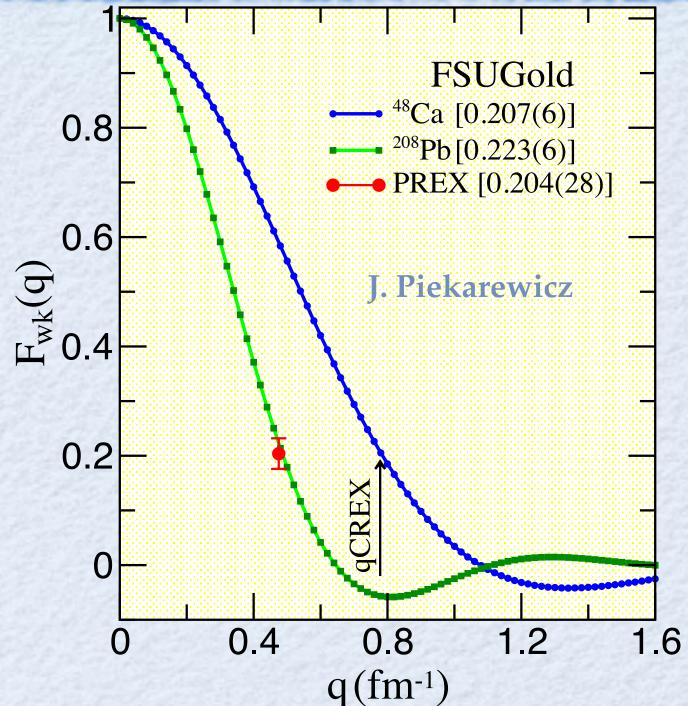
The Parity Violation Program At Jefferson Lab

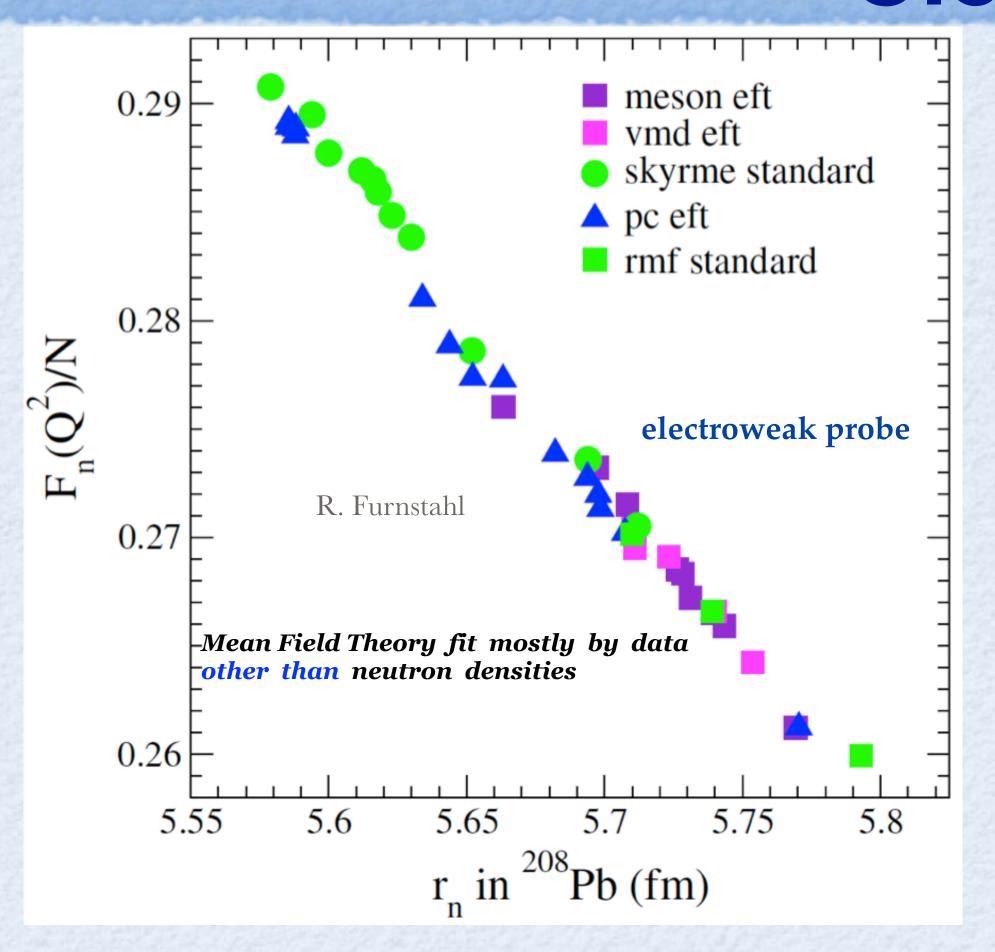
Cleanly Interpretable?



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Cleanly Interpretable?



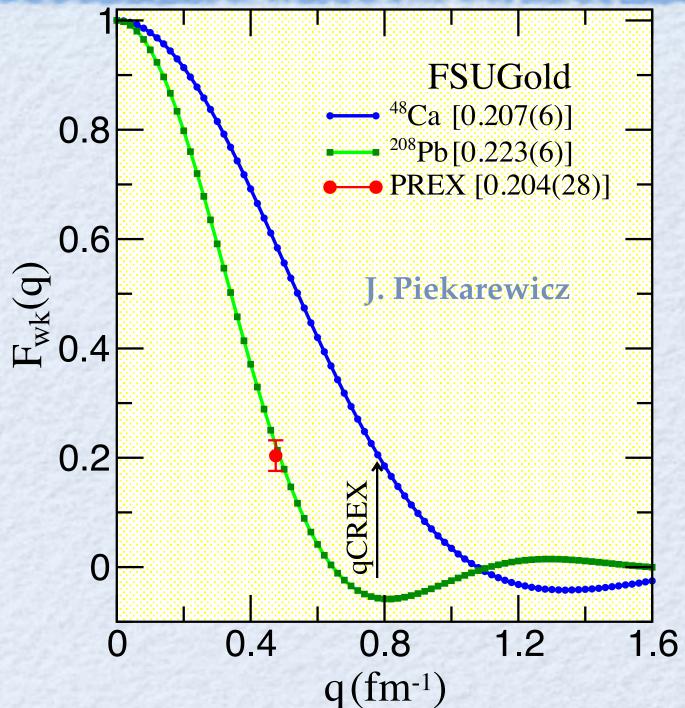


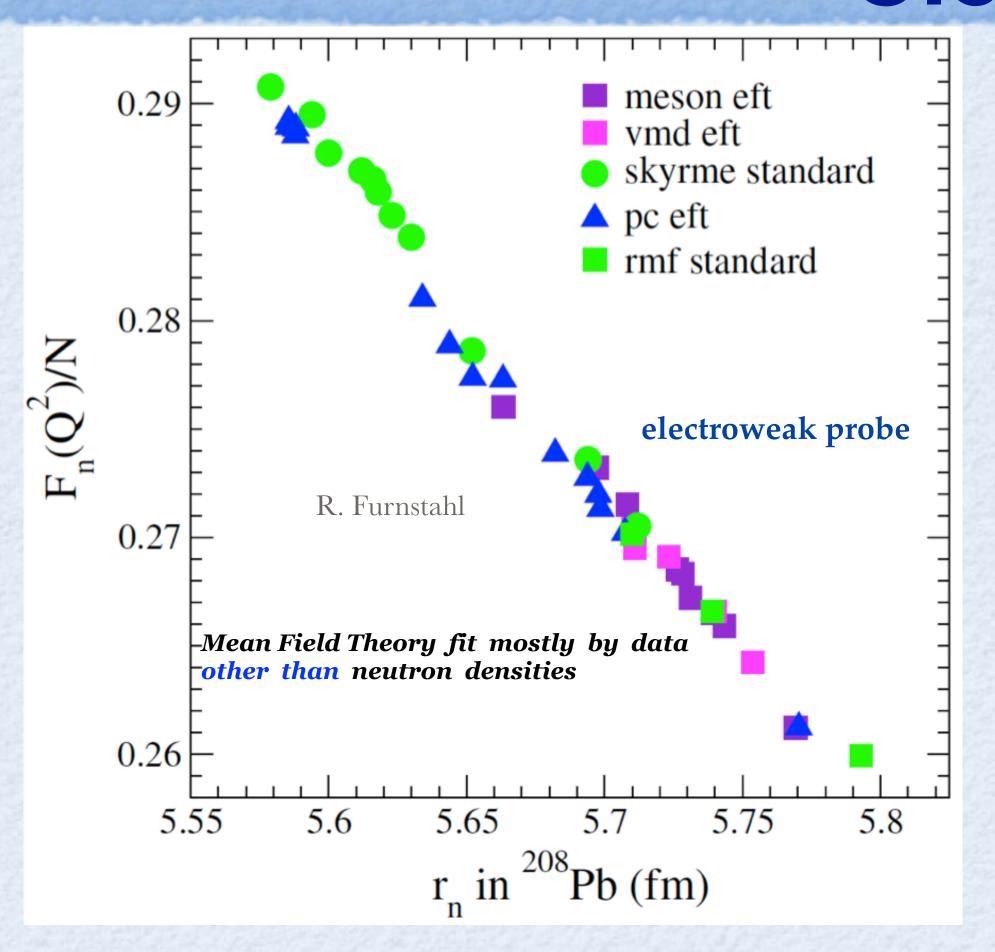
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Cleanly Interpretable?

Q² ~ 0.007 GeV² 5° scattering angle

A_{PV} ~ 0.5 ppm Rate ~ 4 GHz $\delta(A_{PV}) \sim 15 \text{ ppb!}$





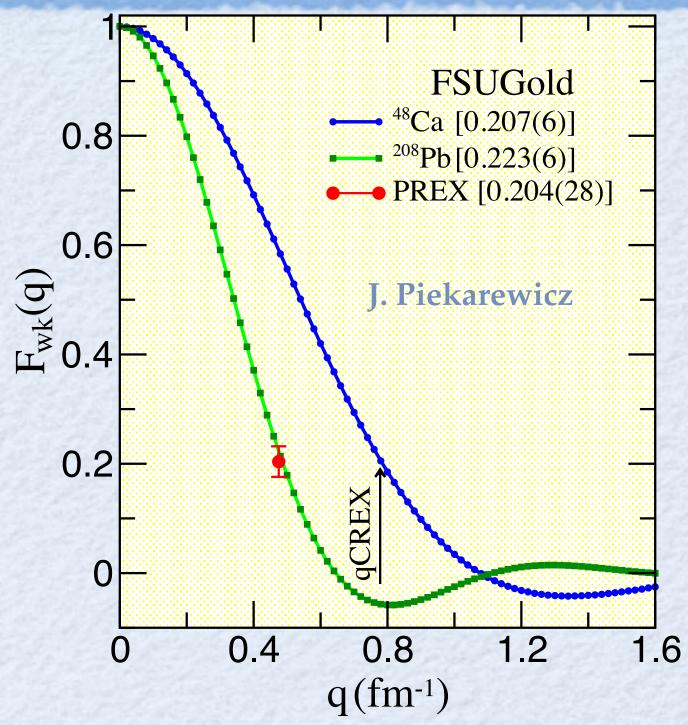
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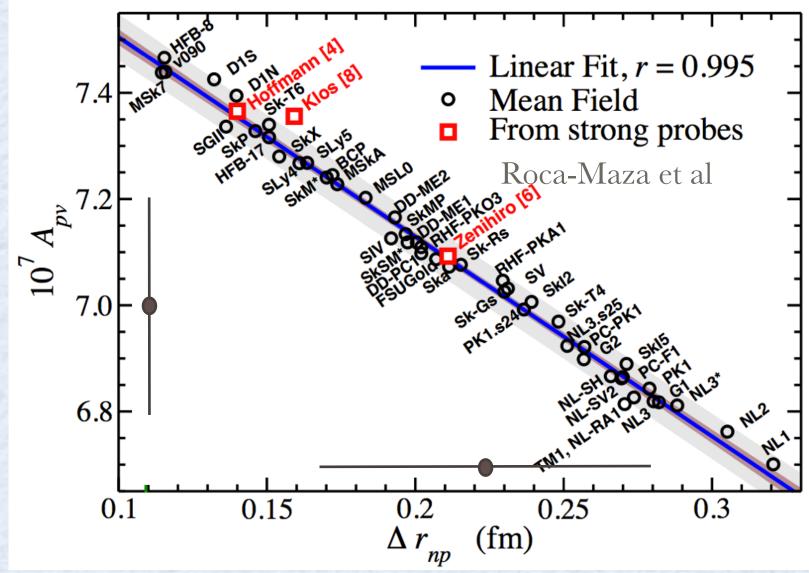
The Parity Violation Program At Jefferson Lab

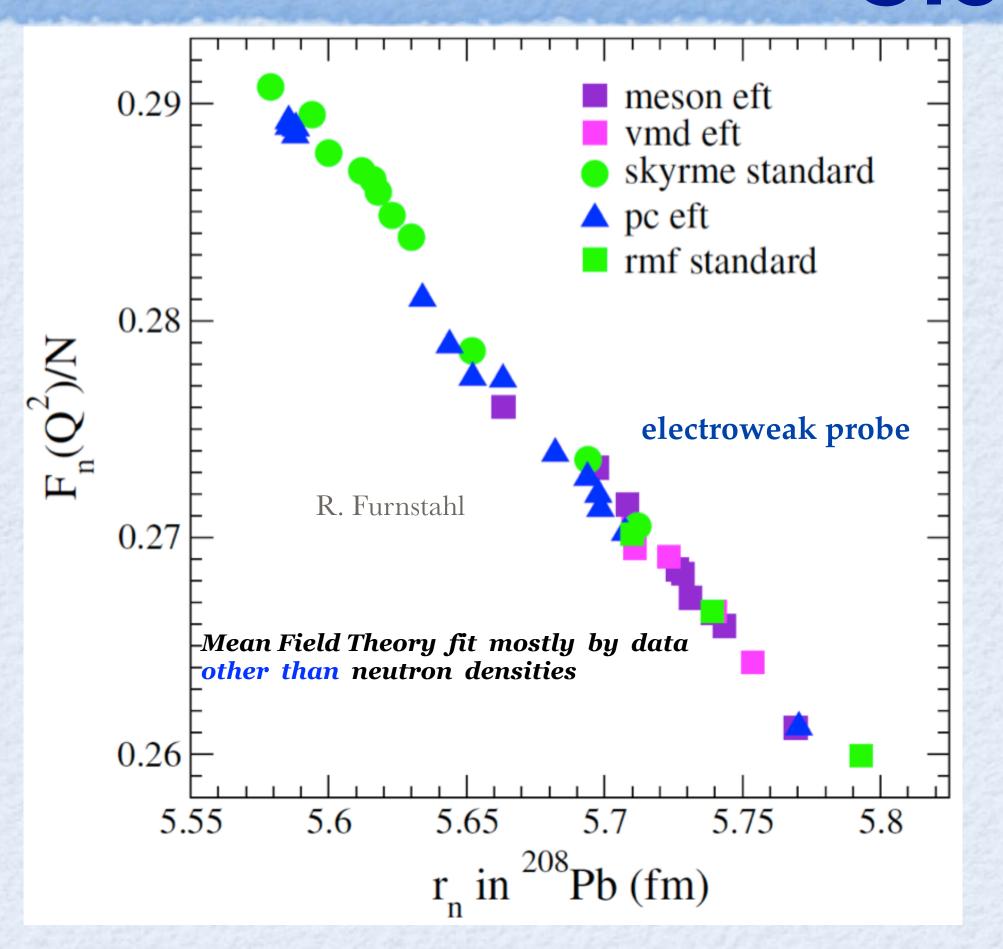
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 Clean "translation" from A_{PV} uncertainty to neutron RMS radius uncertainty

• Apv uncertainty dominated by statistics!

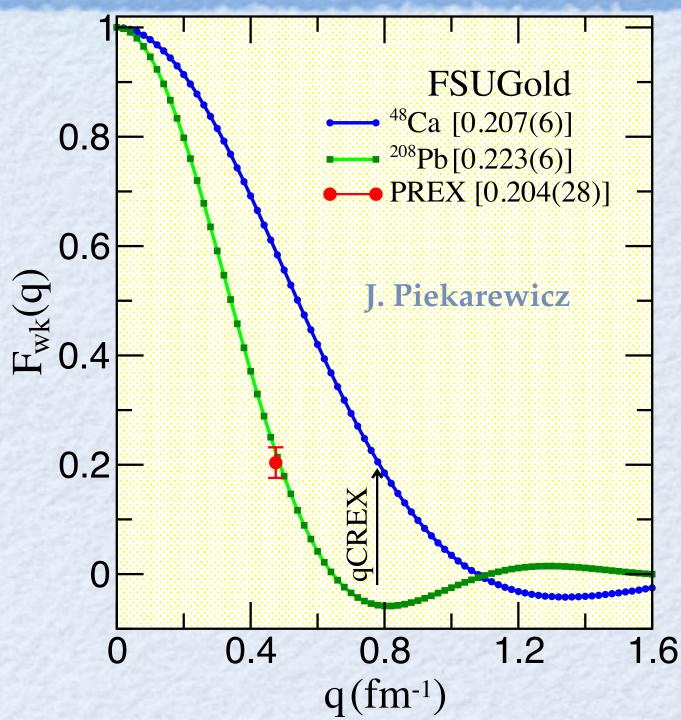
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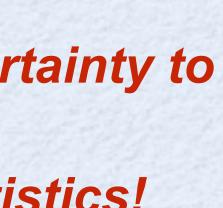
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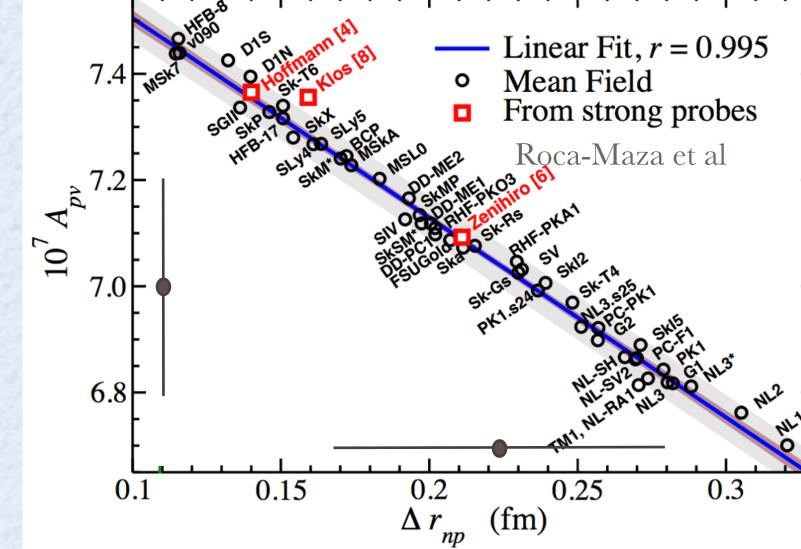
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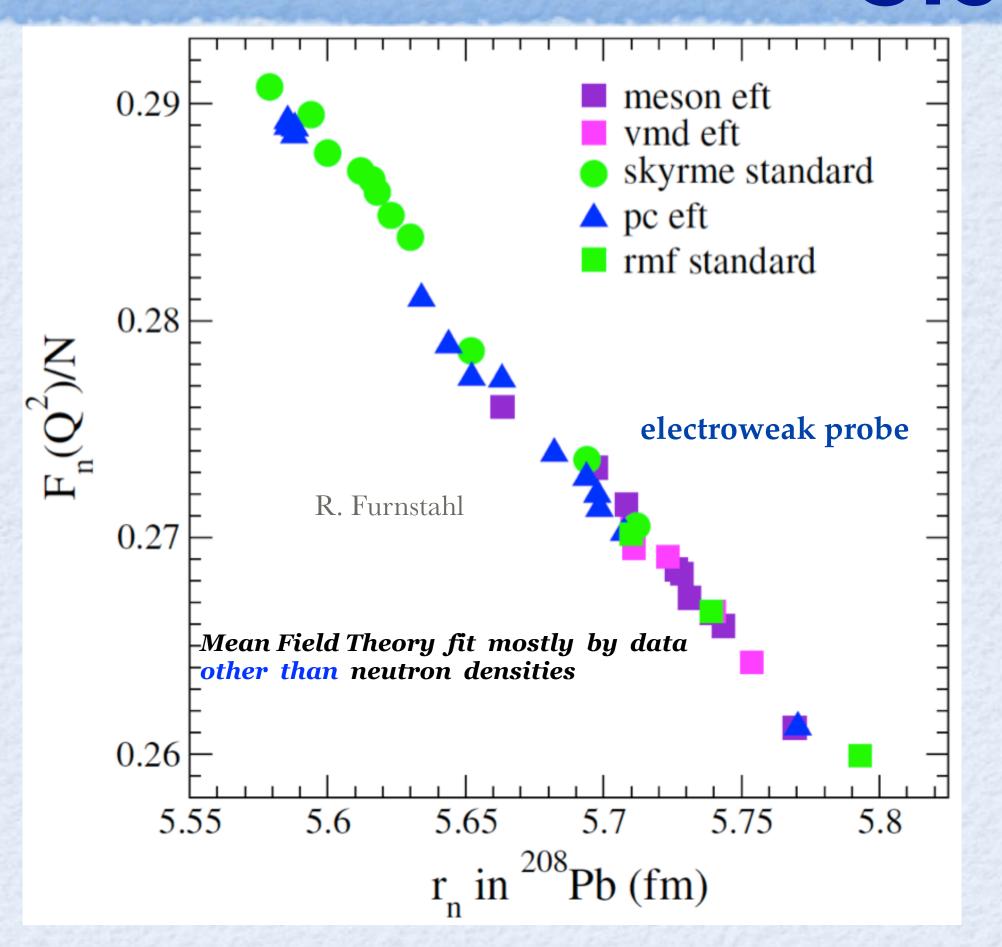


A_{PV} ~ 0.5 ppm Rate ~ 4 GHz $\delta(A_{PV}) \sim 15 \text{ ppb!}$









 Clean "translation" from A_{PV} uncertainty to neutron RMS radius uncertainty

• A_{PV} uncertainty dominated by statistics!

At this level of precision, one must account carefully for radiative corrections

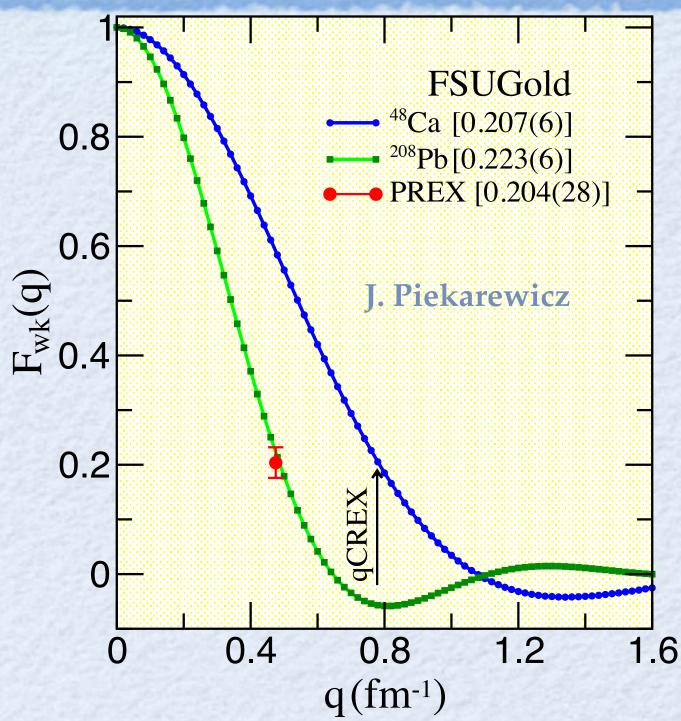
Krishna S. Kumar

The Parity Violation Program At Jefferson Lab

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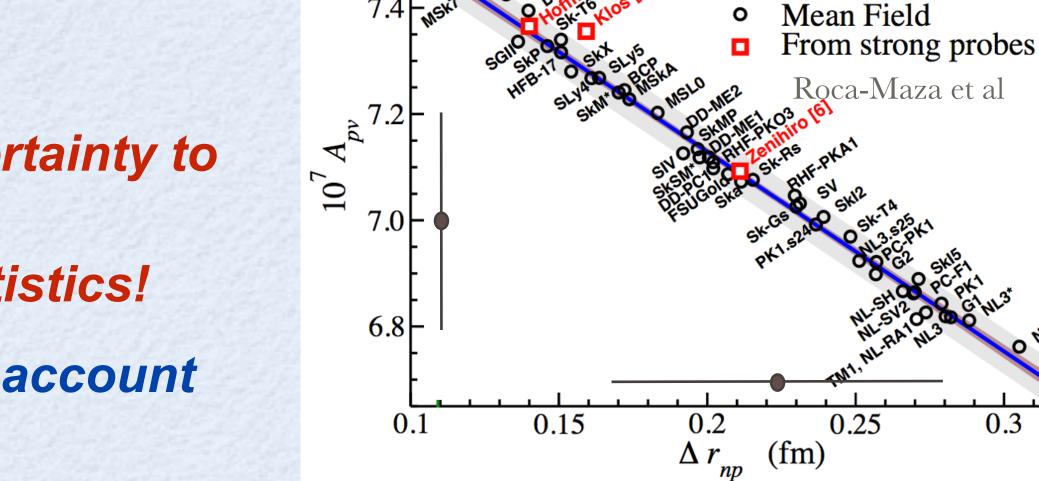


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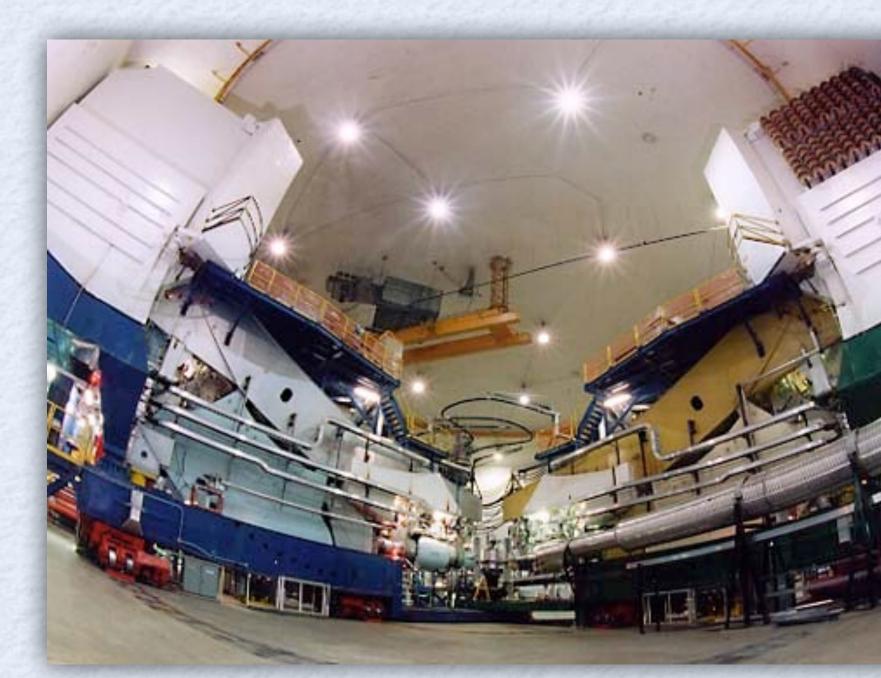
Linear Fit, r = 0.995

0.3





1 GeV electron beam, 50-70 µA high polarization, ~89% helicity reversal at 120 Hz



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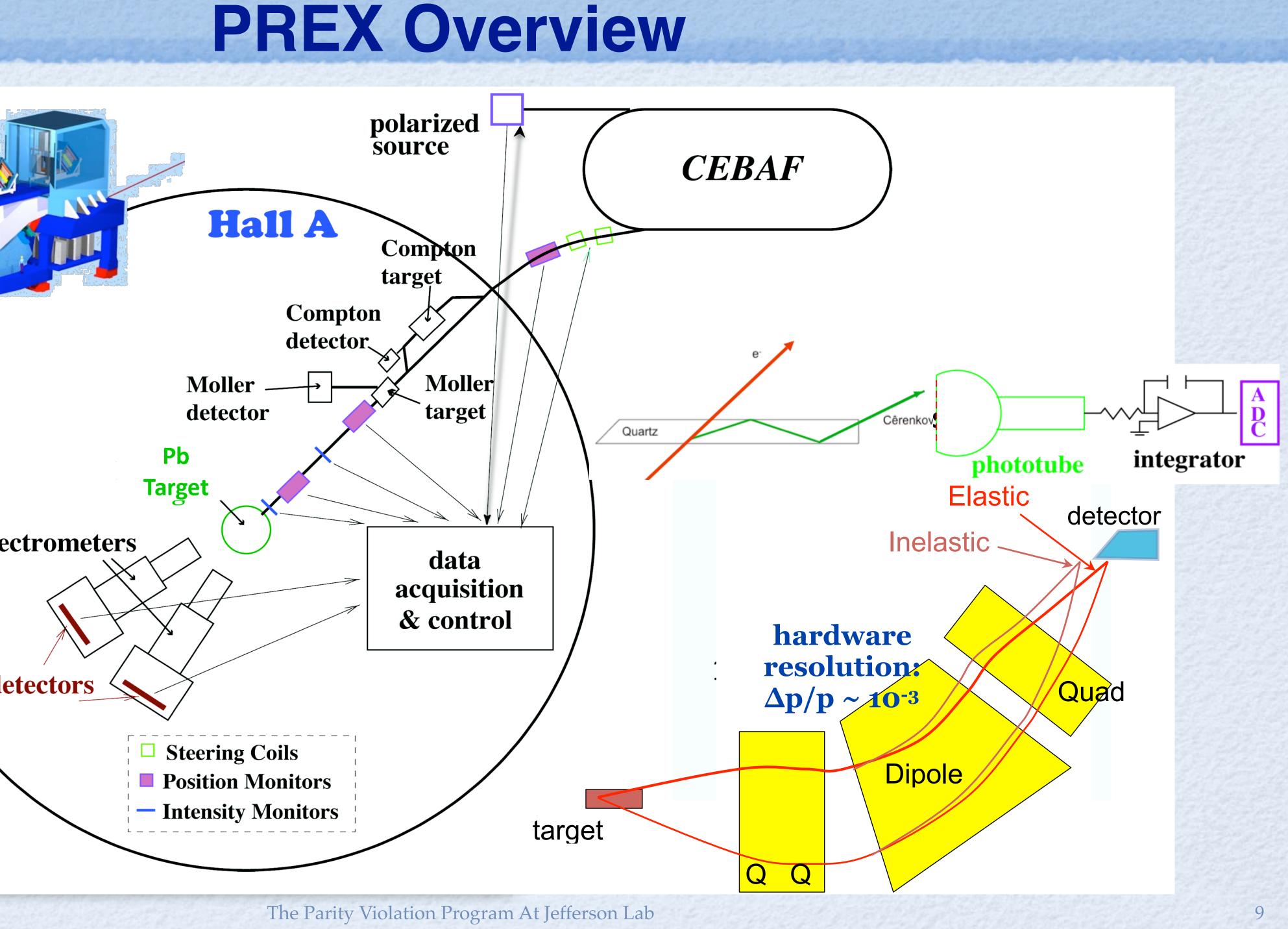
The Parity Violation Program At Jefferson Lab

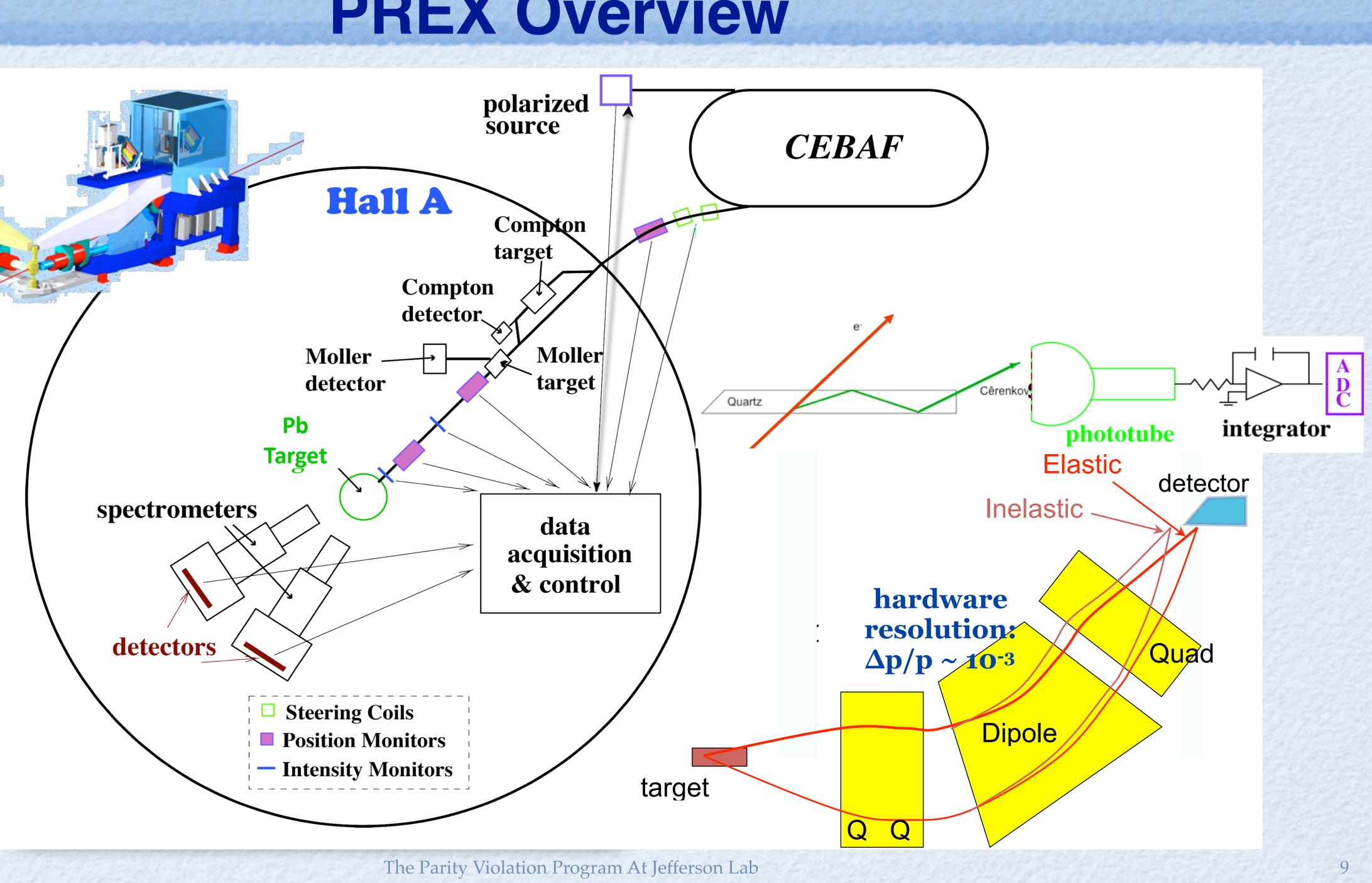
PREX Overview



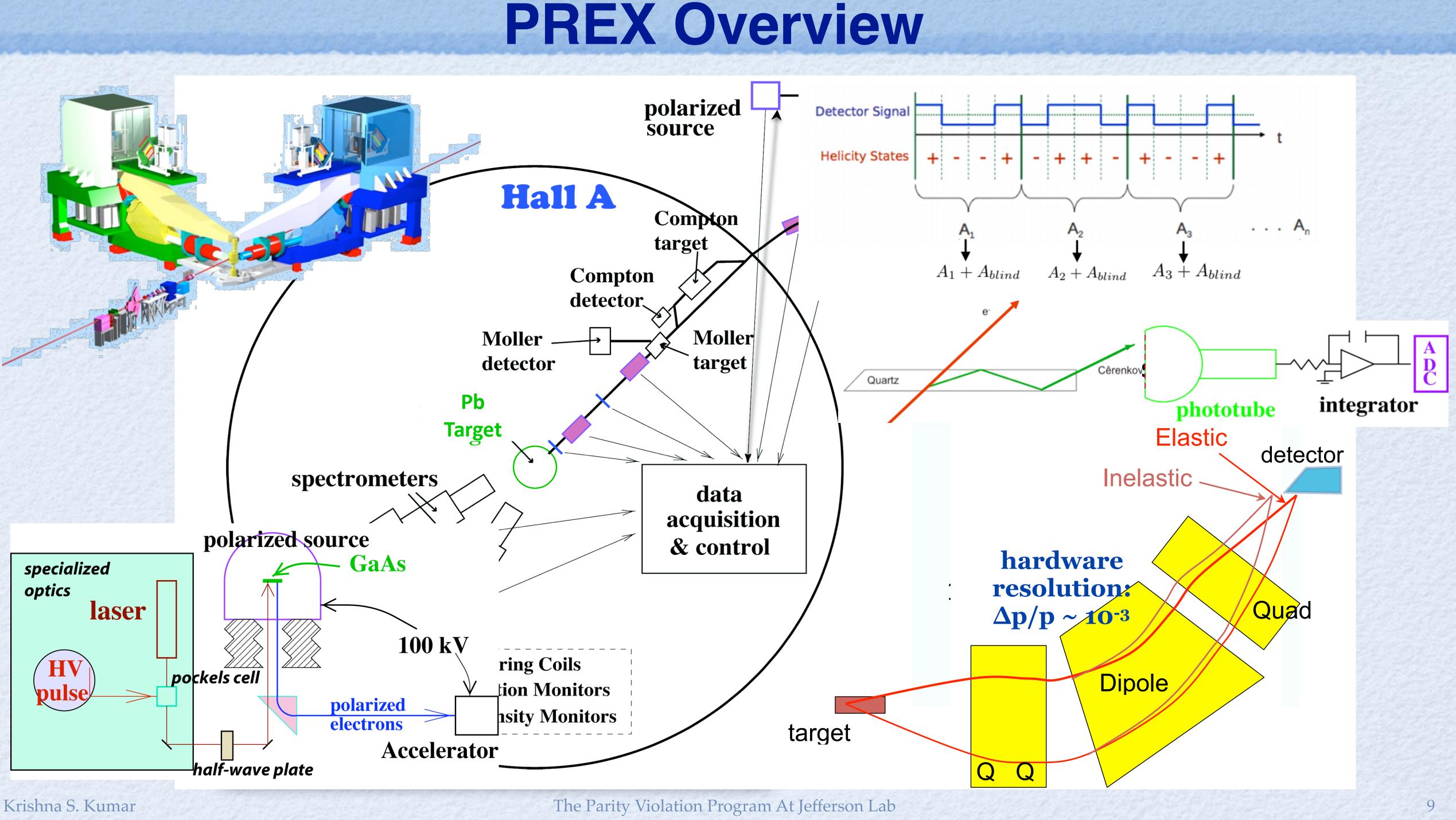
0.5 mm isotopically pure ²⁰⁸Pb target 5° scattered electrons $Q^2 = 0.007 \, GeV^2/c^2$ new thin quartz detectors



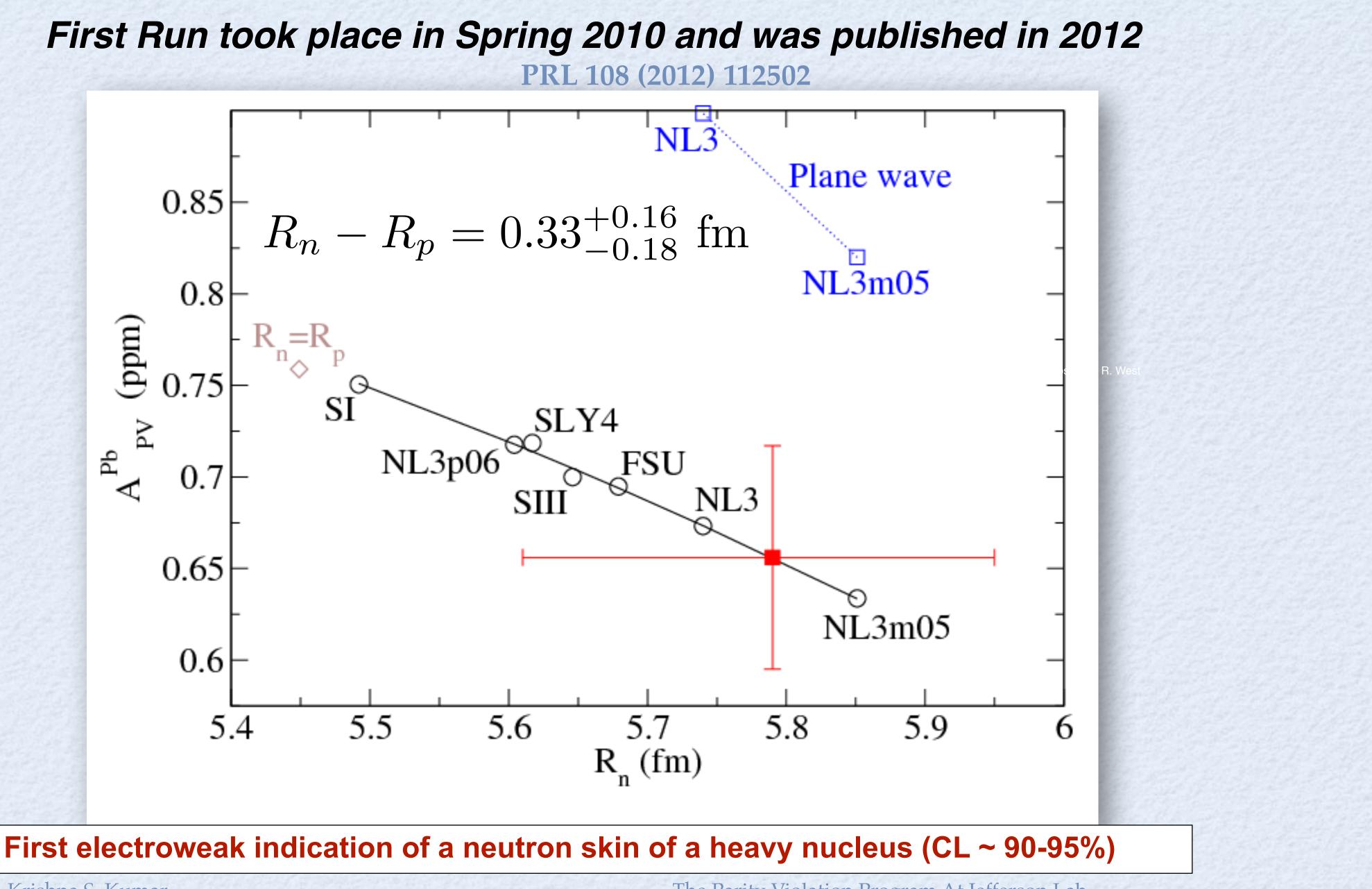








The Neutron Skin



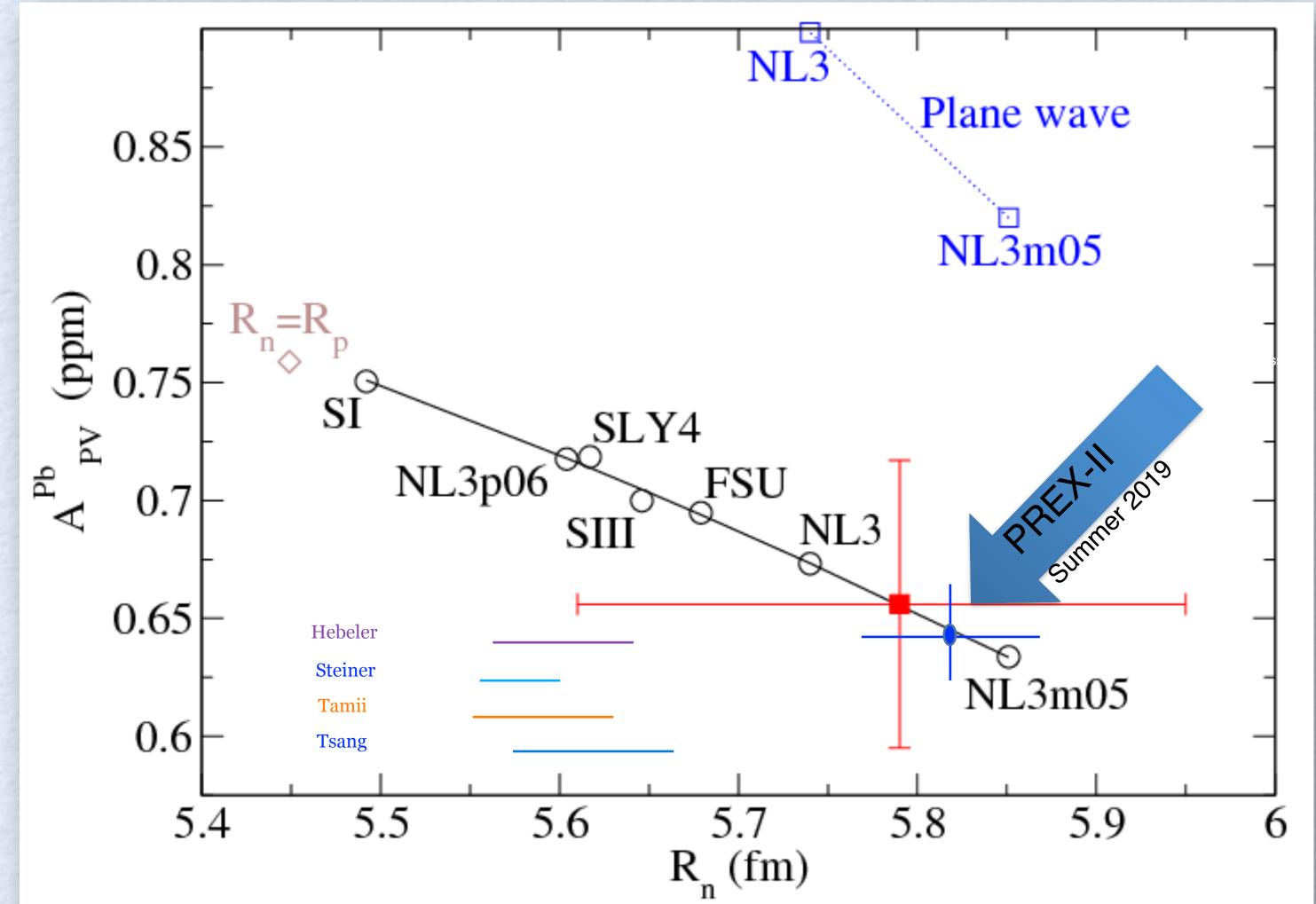
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The Parity Violation Program At Jefferson Lab

The Neutron Skin

First Run took place in Spring 2010 and was published in 2012

PRL 108 (2012) 112502



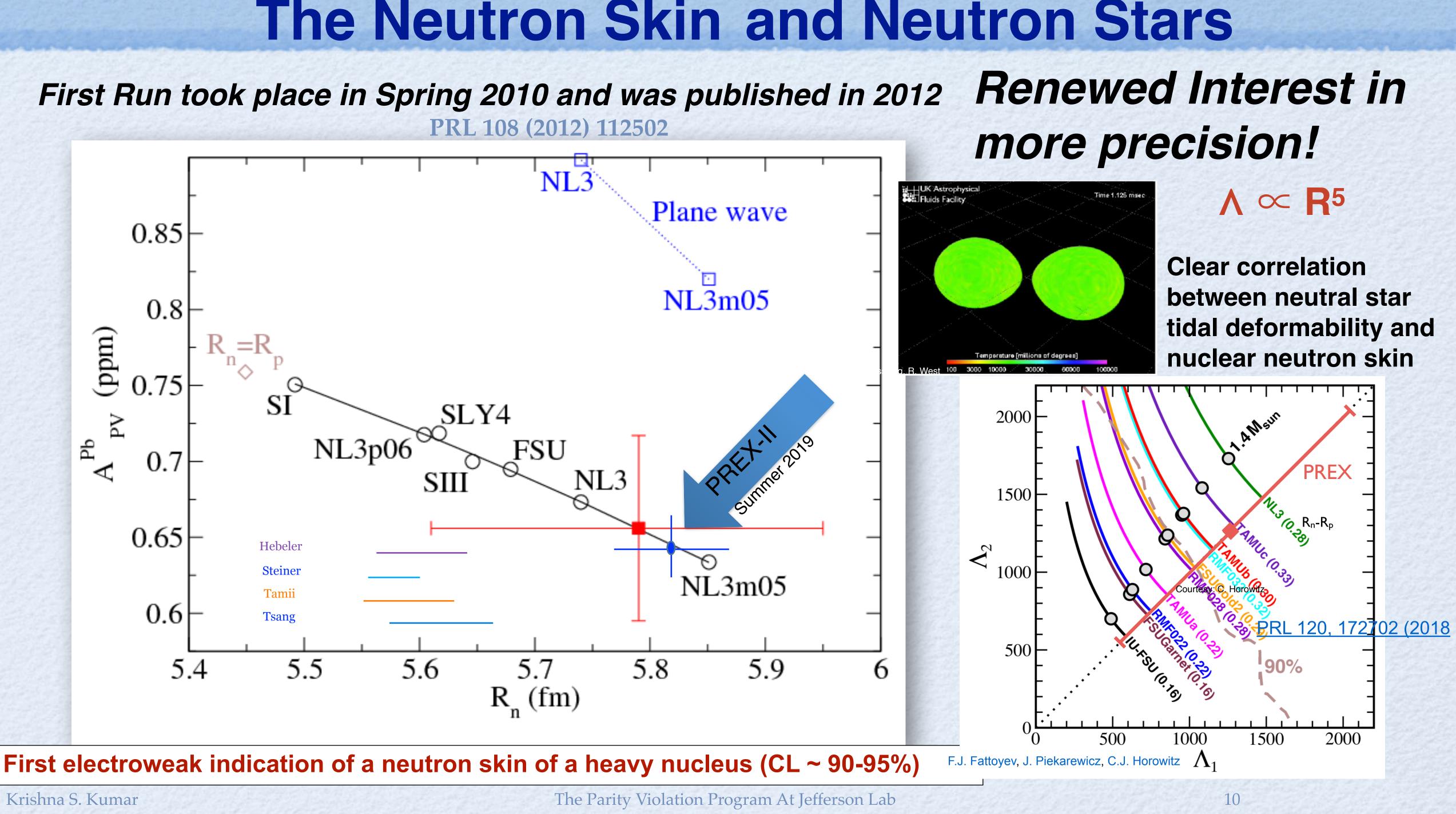
First electroweak indication of a neutron skin of a heavy nucleus (CL ~ 90-95%)

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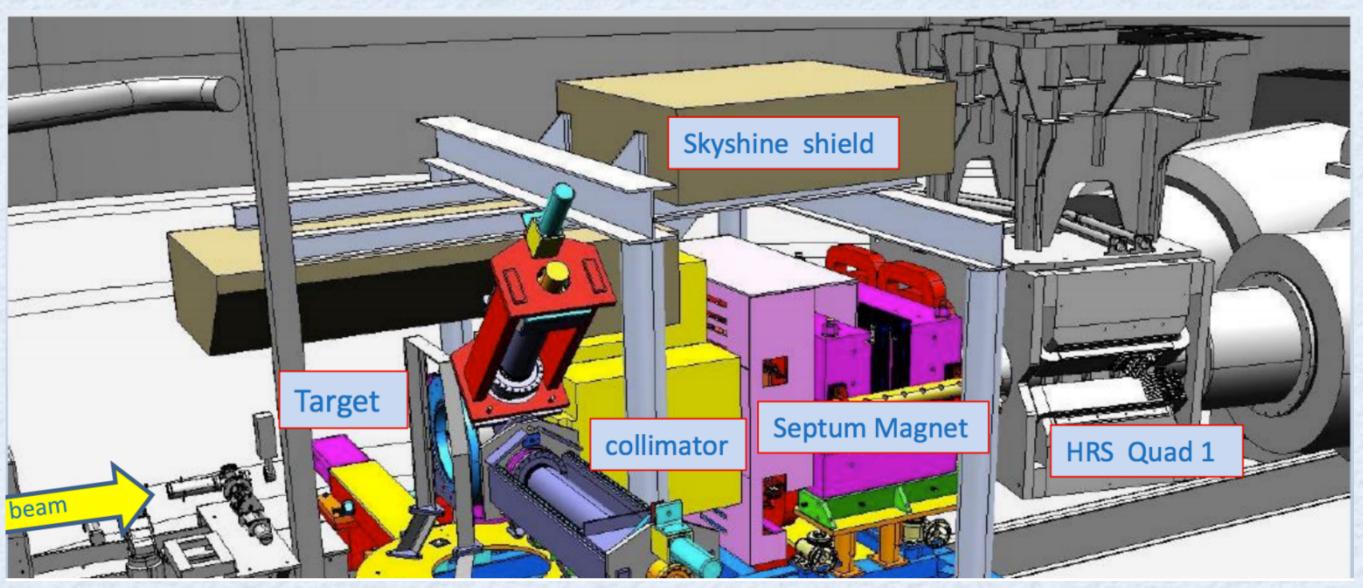
The Parity Violation Program At Jefferson Lab

The Neutron Skin and Neutron Stars First Run took place in Spring 2010 and was published in 2012

(2012) 112502



Krishna S. Kumar

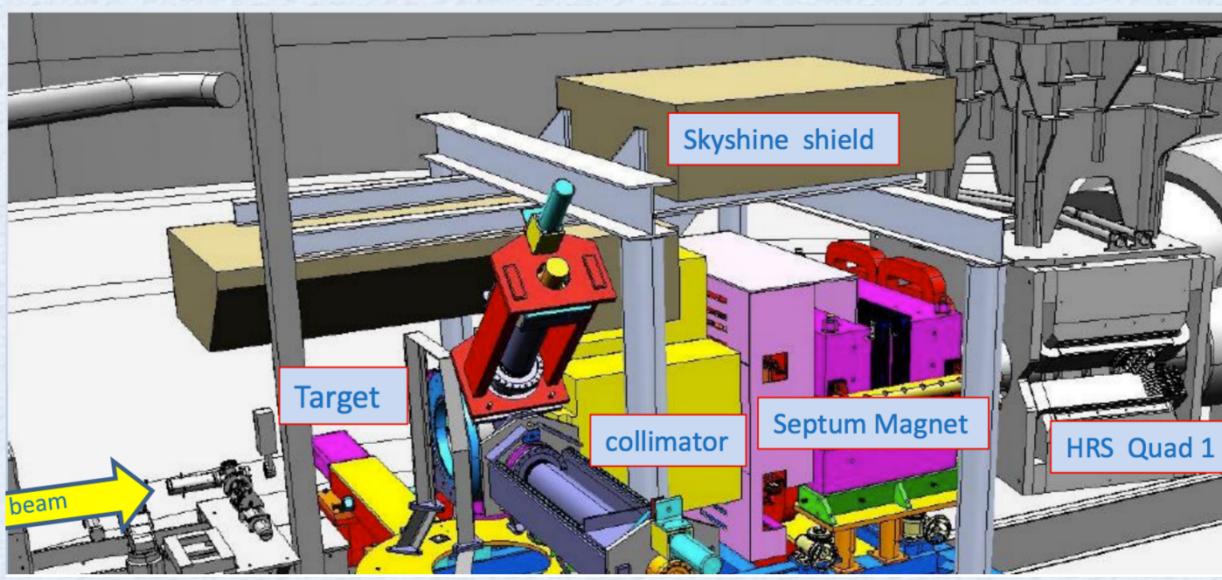


- Physics running took place in July/August 2019
- •126 Coulombs on target; 116 passed online cuts
- •On track to achieve proposal goal:
- •about 3% stat and 1.5% syst. uncertainties

My view! Gordon Conference rules!

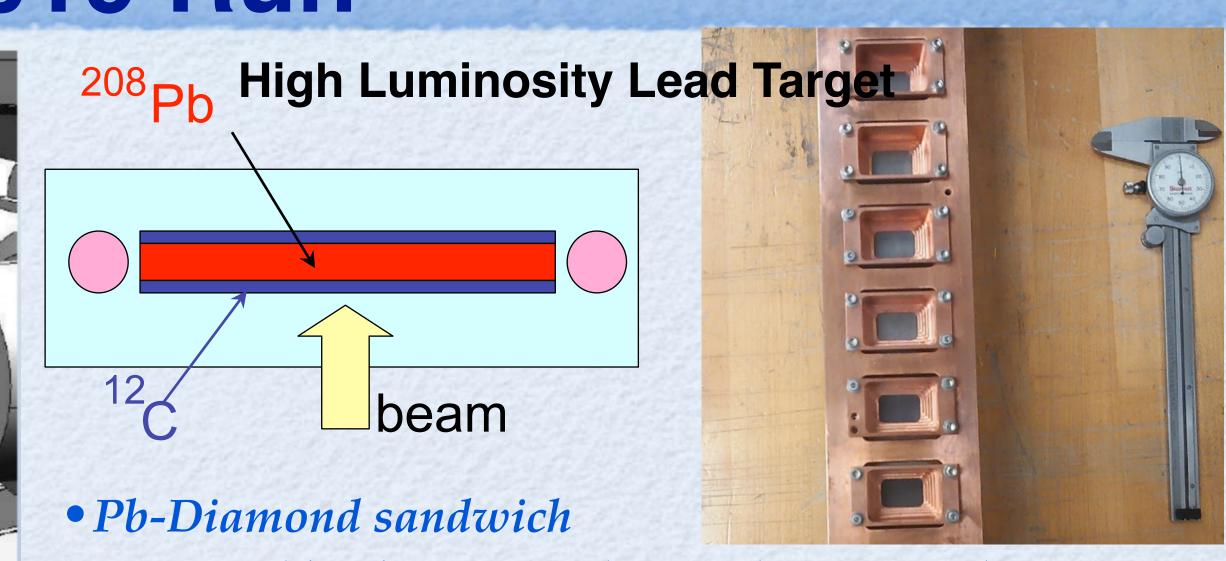
The Parity Violation Program At Jefferson Lab





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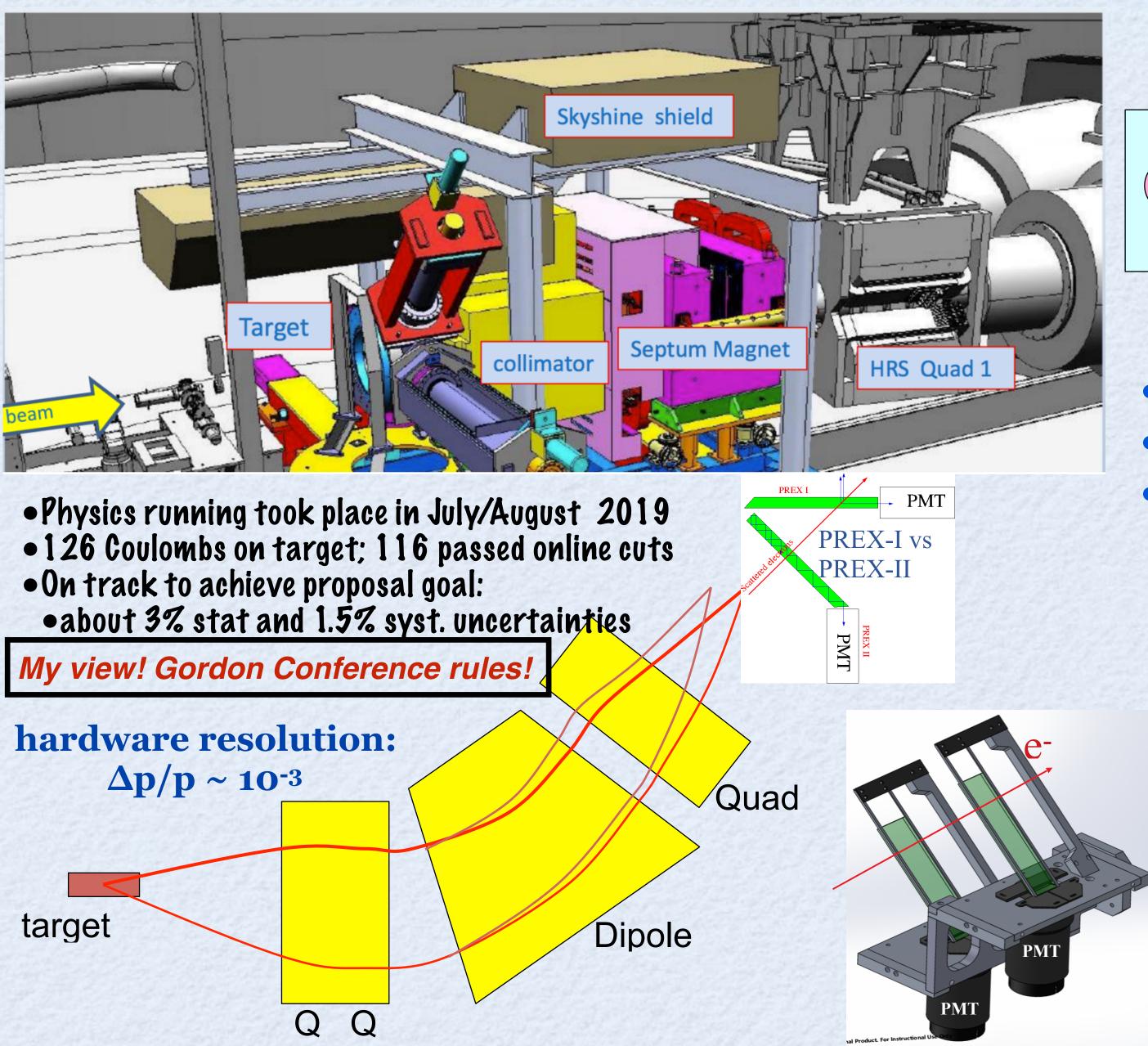
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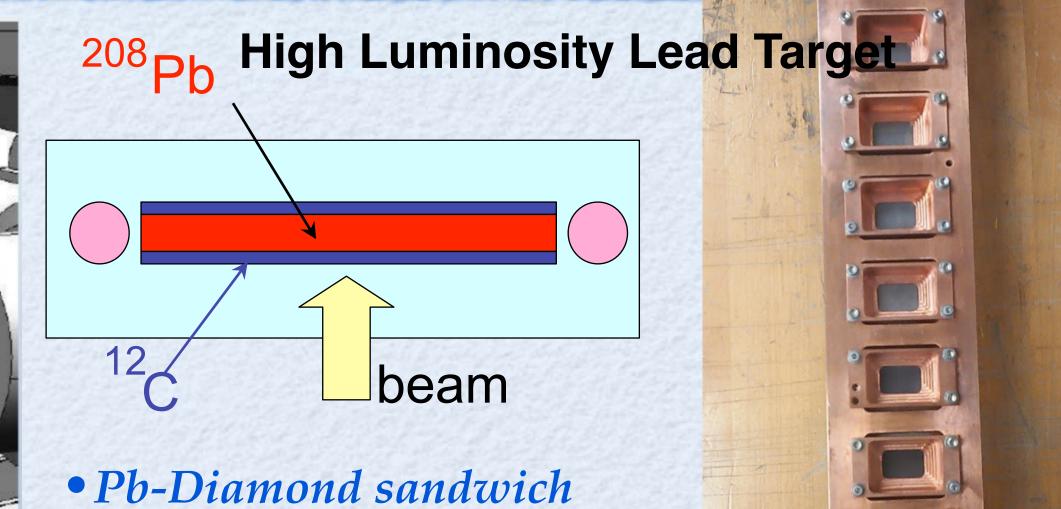
- Diamond backing provides conductive cooling
- Active cryo-cooling with available He lines

The Parity Violation Program At Jefferson Lab





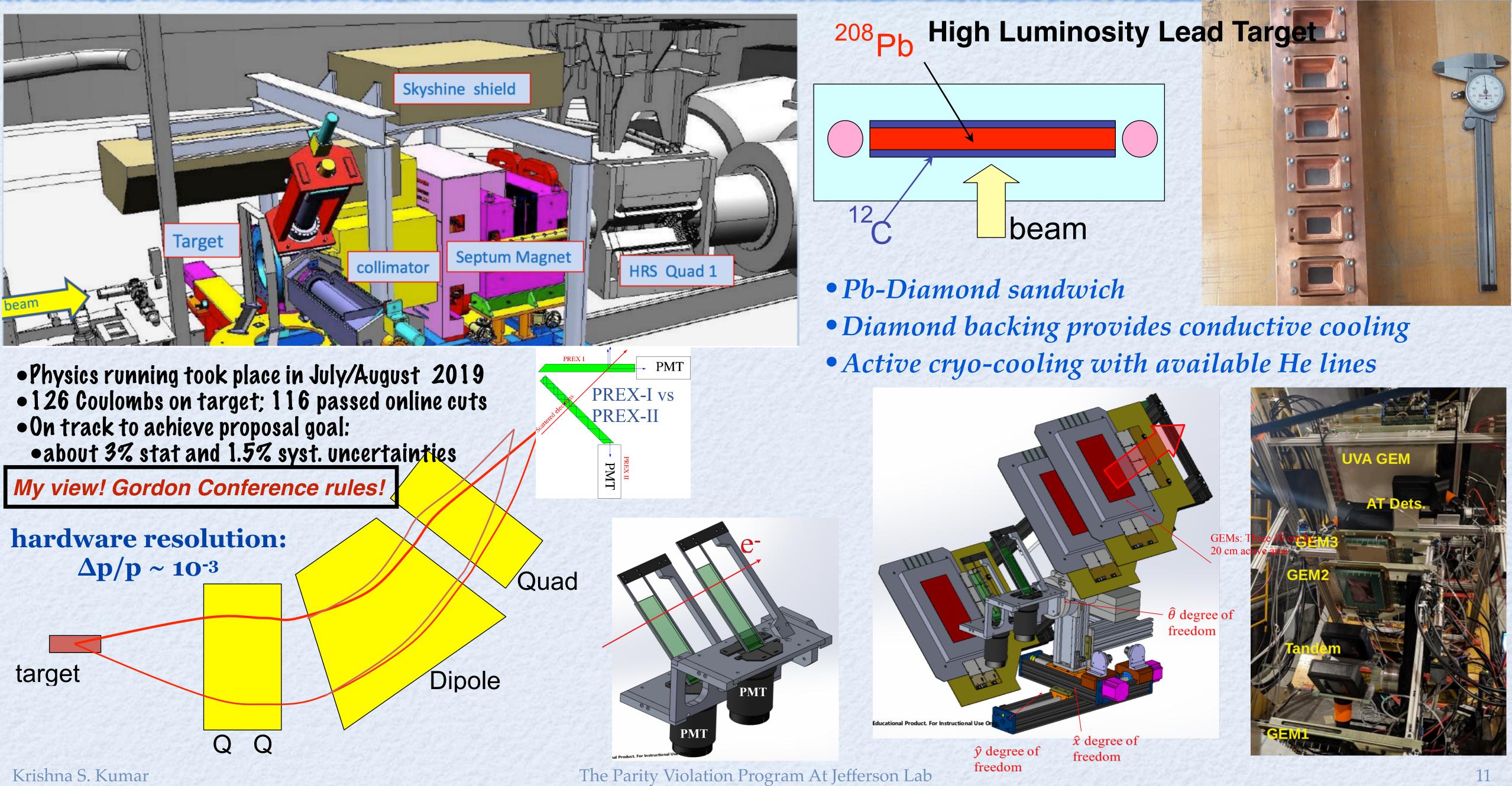
Krishna S. Kumar

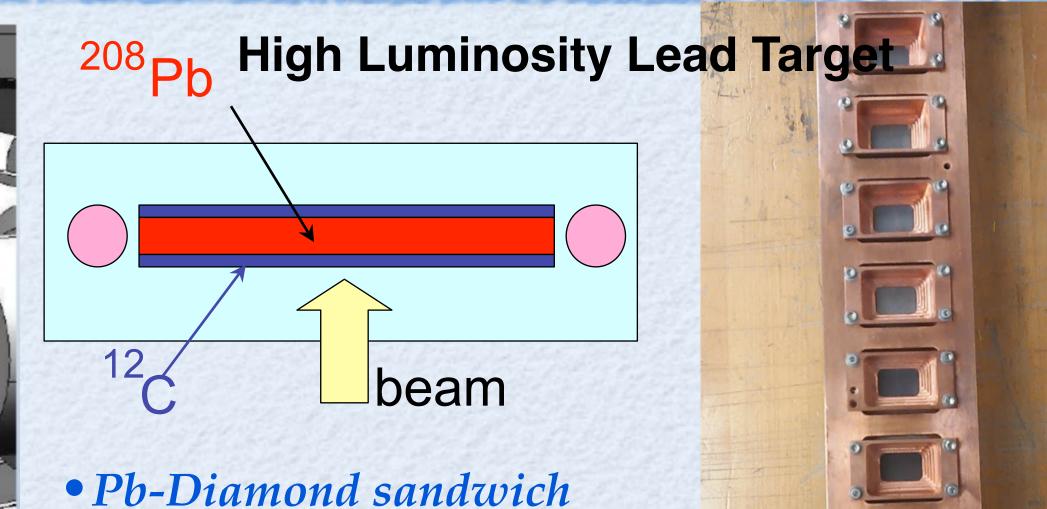


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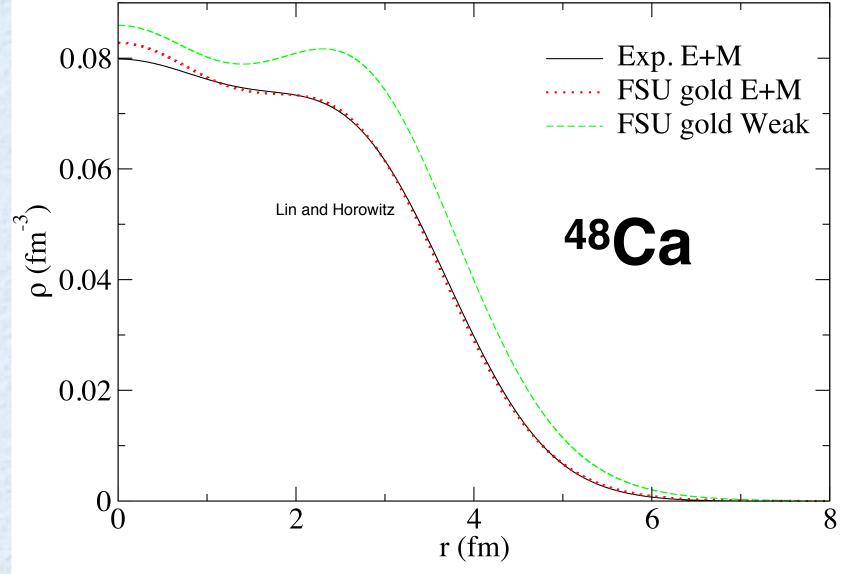








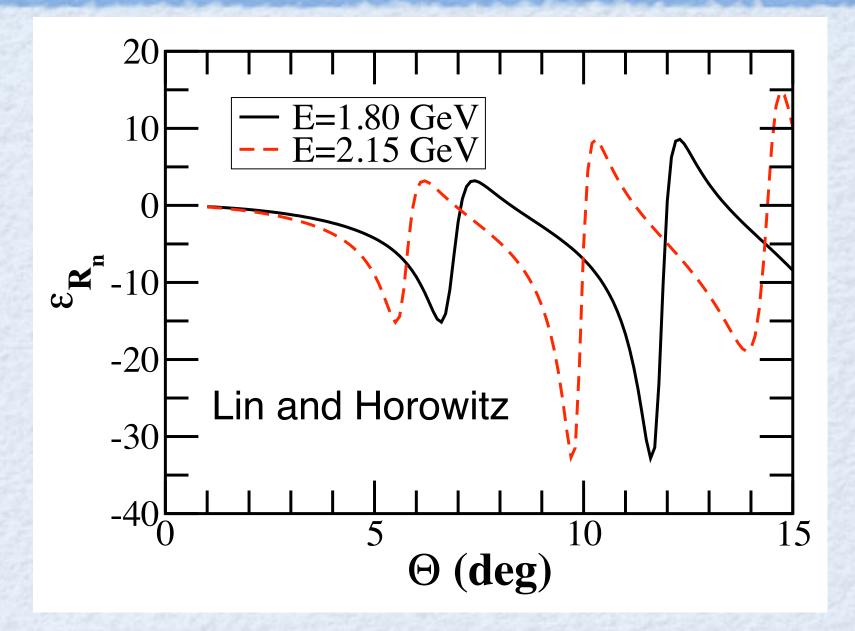
CREX Motivation



Optimum Q ~ 160 MeV $2 \text{ GeV} \rightarrow 5 \text{ degrees}$ A_{PV} ~ 3.8 ppm Rate ~ 40 MHz $\delta(A_{PV}) \sim 3\%$ $\delta(R_n) \sim \pm 0.02 \ fm$

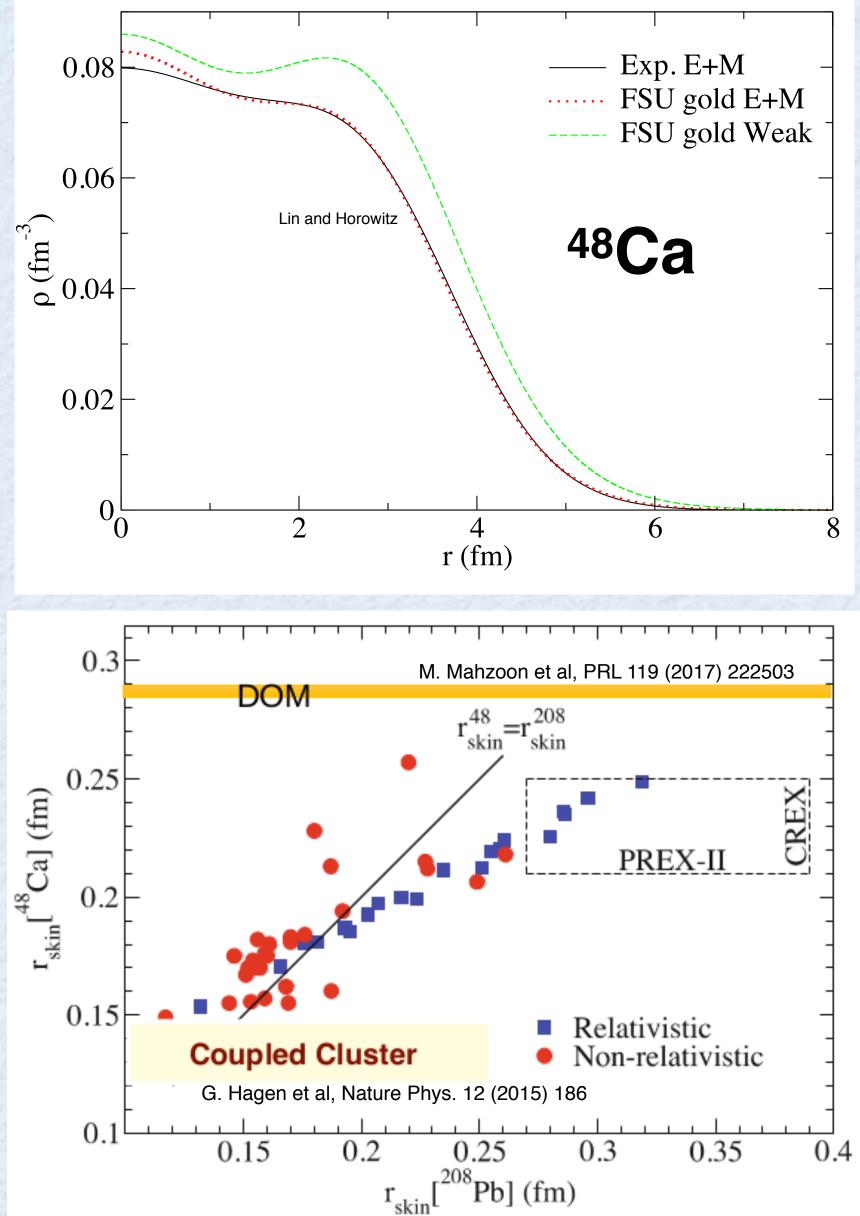
Krishna S. Kumar

Starts December 2!









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3 neutron forces

Diffractive Optical Model (DOM) and Coupled Cluster (CC) predictions very different!

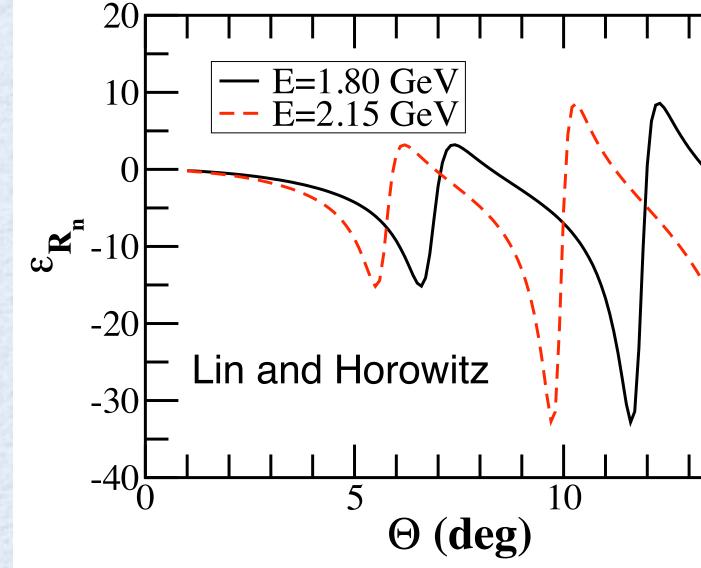
QMCpi-II prediction (Martinez et al, PRC 100 (2019) 024333) similar to CC

The Parity Violation Program At Jefferson Lab

Krishna S. Kumar

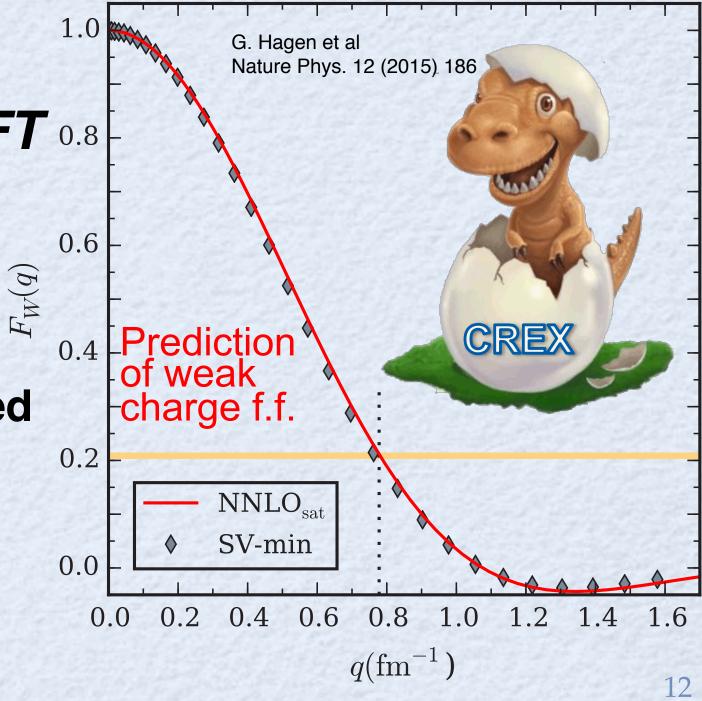
CREX Motivation

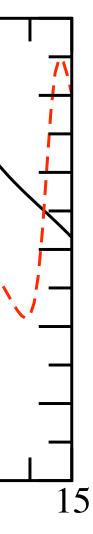
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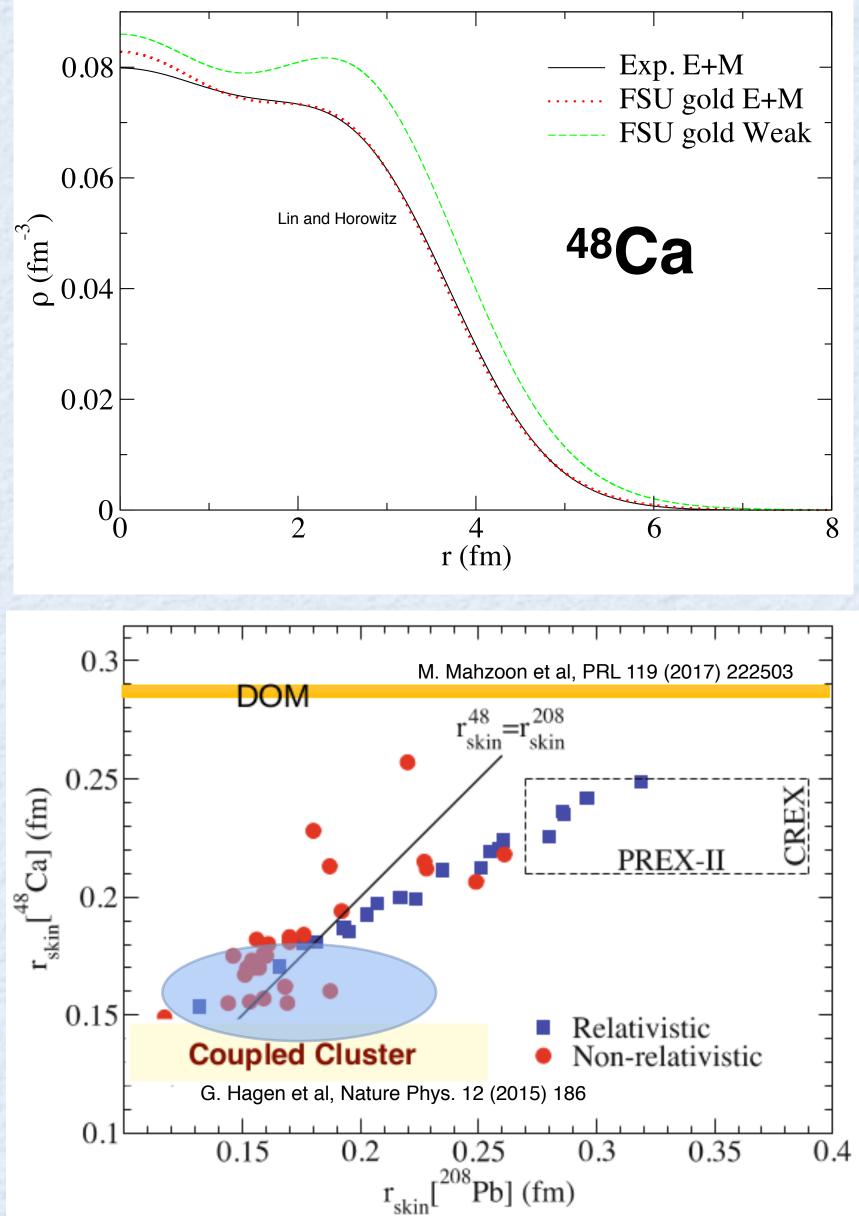
Bridge ab initio calculations and DFT 0.8

Three firm predictions!









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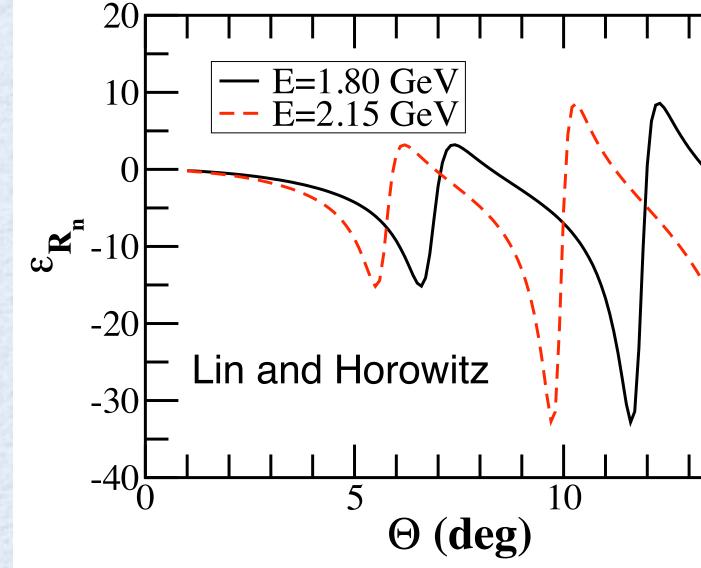
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The Parity Violation Program At Jefferson Lab

Krishna S. Kumar

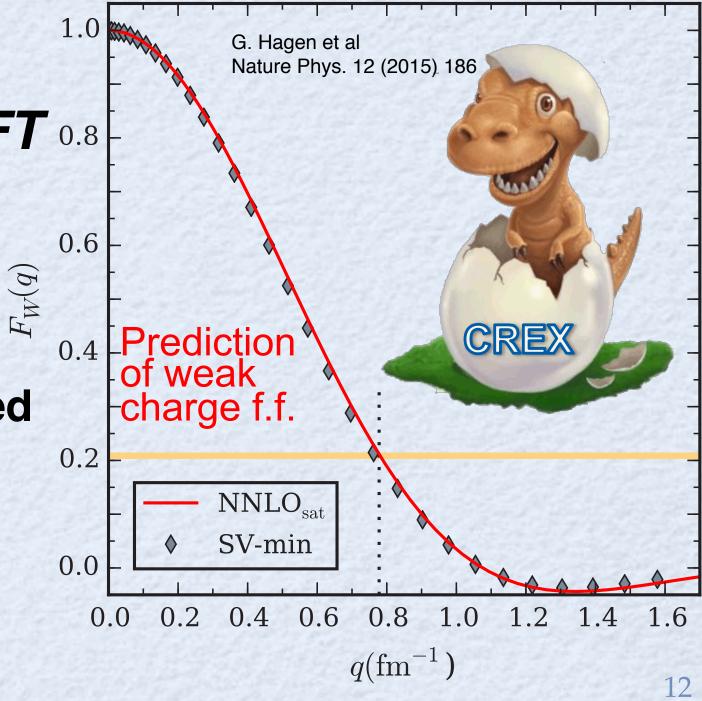
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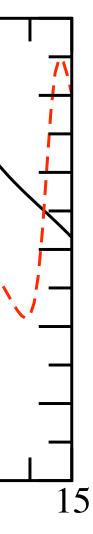
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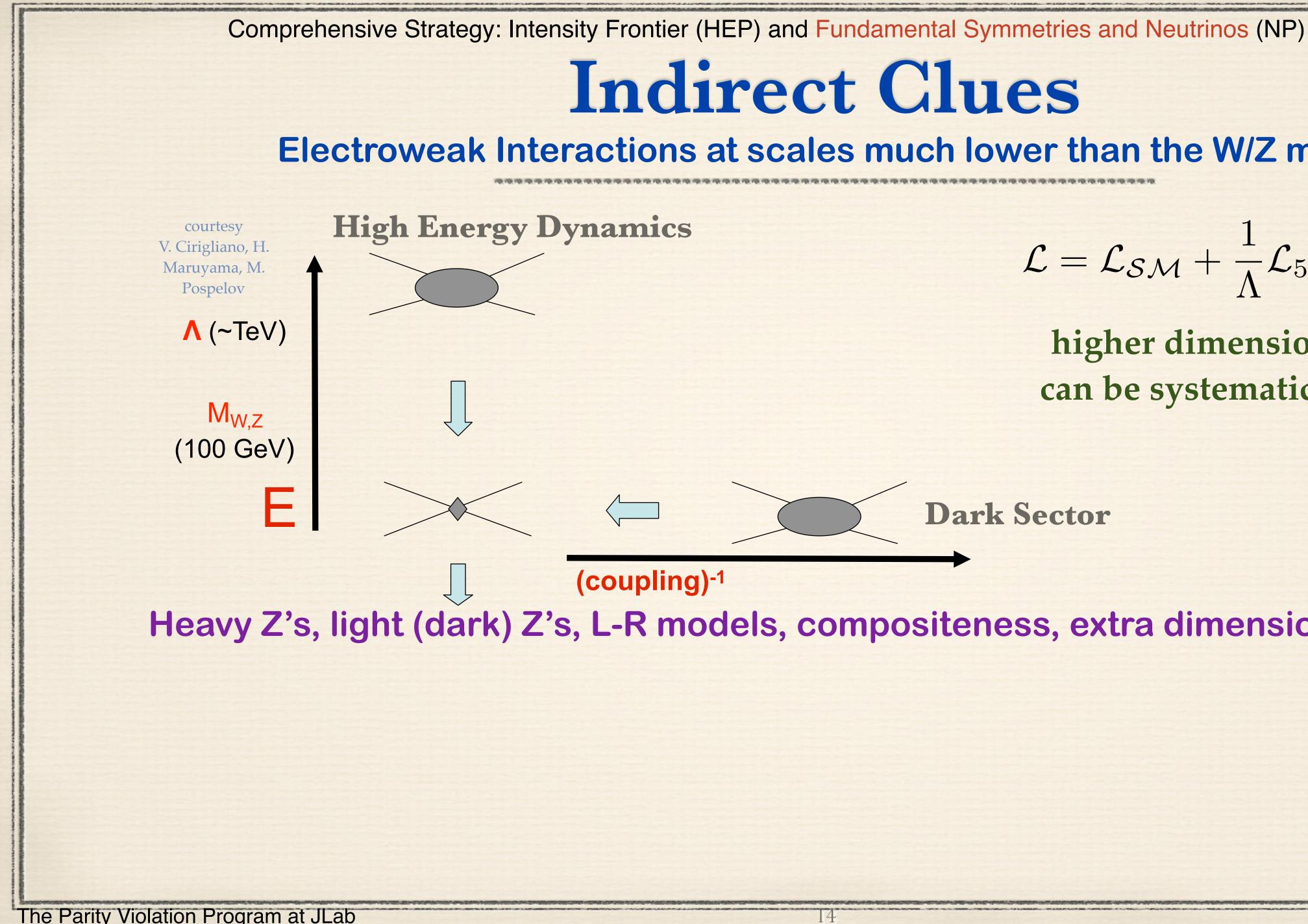
Physics down to a length scale of 10⁻¹⁹ m well understood but.... **Modern Electroweak Physics** Many questions still unanswered.... **The High Energy Frontier: Collider Physics** The Cosmic Frontier: Particle, Nuclear and Gravitational Astrophysics A comprehensive search for clues requires, in addition: **The Intensity/Precision Frontier**



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The Parity Violation Program at JLsb





Indirect Clues Electroweak Interactions at scales much lower than the W/Z mass

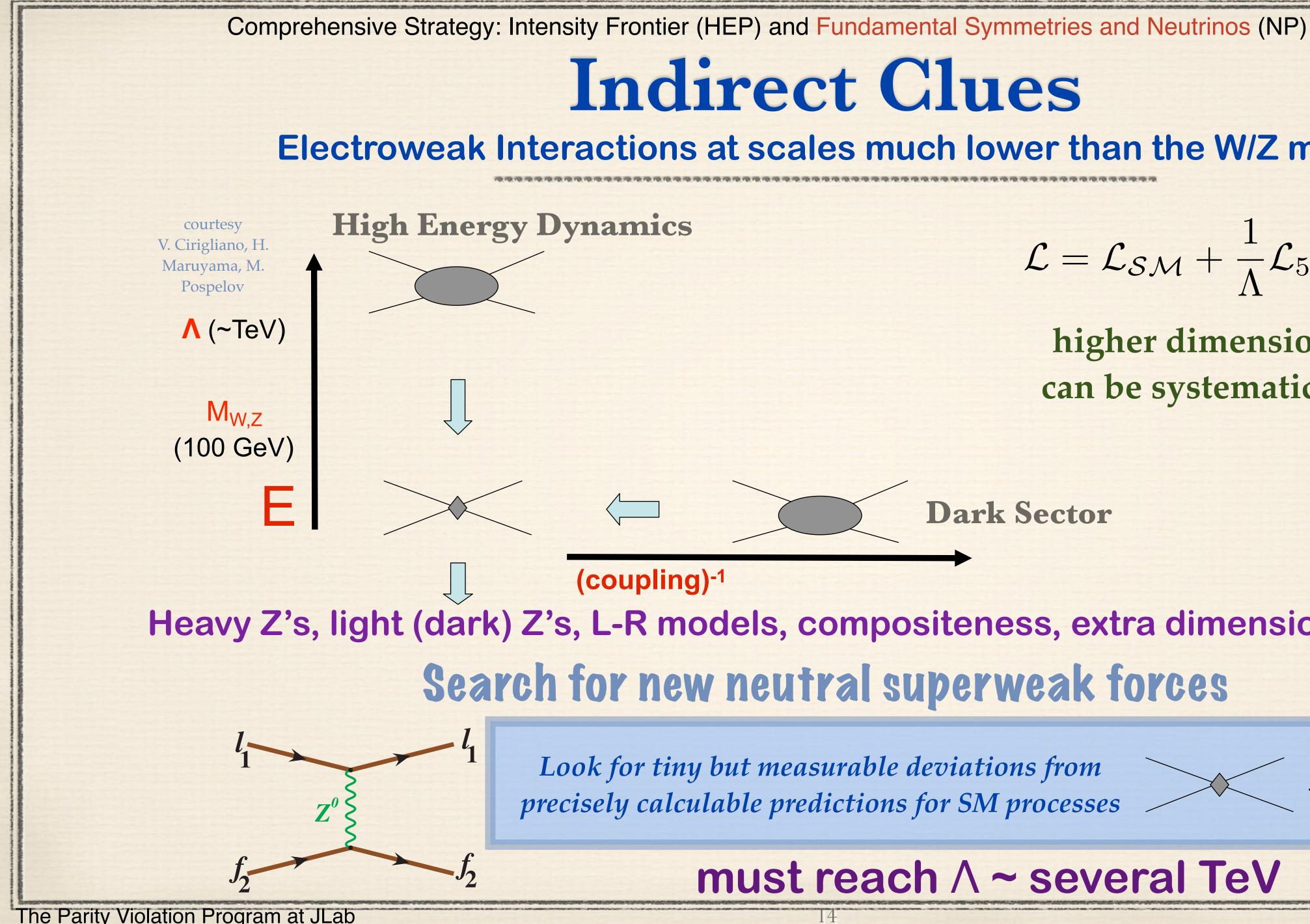
 $\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{\Lambda}\mathcal{L}_5 + \frac{1}{\Lambda^2}\mathcal{L}_6 + \cdots$

higher dimensional operators can be systematically classified



Heavy Z's, light (dark) Z's, L-R models, compositeness, extra dimensions, SUSY...





Indirect Clues Electroweak Interactions at scales much lower than the W/Z mass

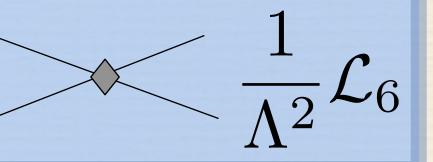
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Dark Sector

Heavy Z's, light (dark) Z's, L-R models, compositeness, extra dimensions, SUSY... Search for new neutral superweak forces

Look for tiny but measurable deviations from precisely calculable predictions for SM processes

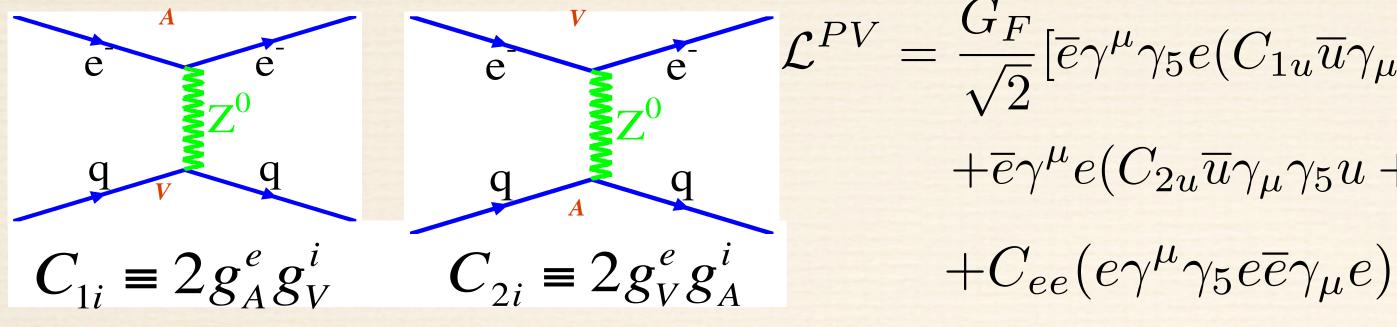


must reach $\Lambda \sim several TeV$

Krishna Kumar, October 30, 2019



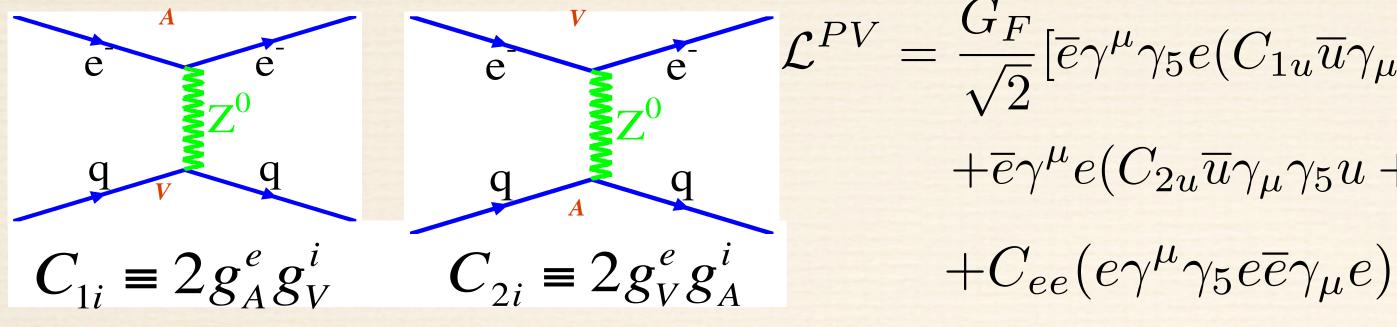




 $\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\overline{e}\gamma^{\mu}\gamma_5 e(C_{1u}\overline{u}\gamma_{\mu}u + C_{1d}\overline{d}\gamma_{\mu}d)]$ $+\overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_{5}u+C_{2d}\overline{d}\gamma_{\mu}\gamma_{5}d)]$

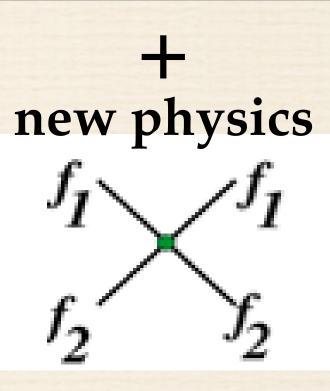






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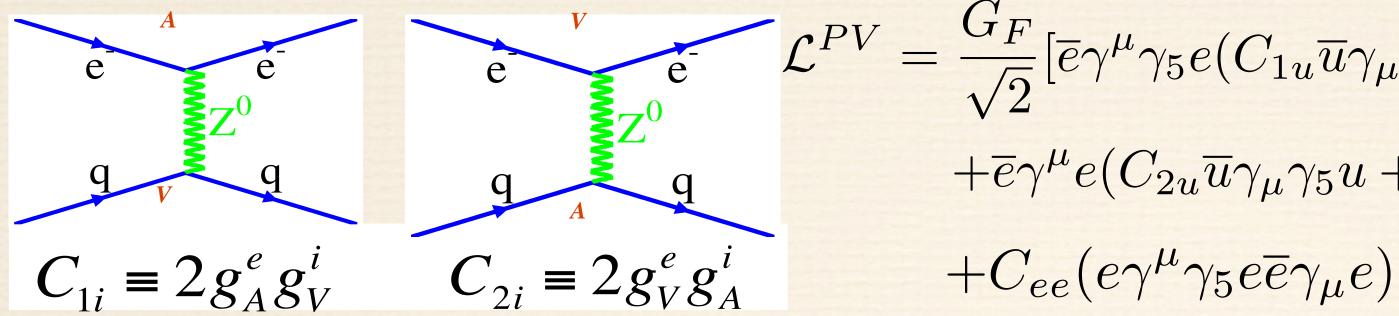
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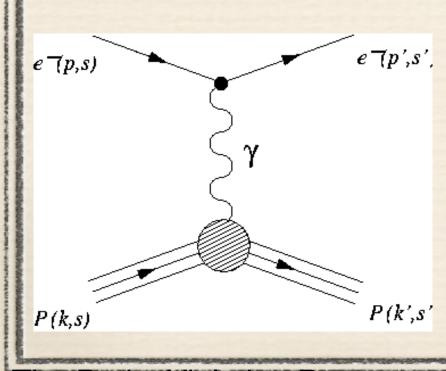
 $\mathcal{L}_{f_1 f_2}$

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma_\mu f_{2j}$





$C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$

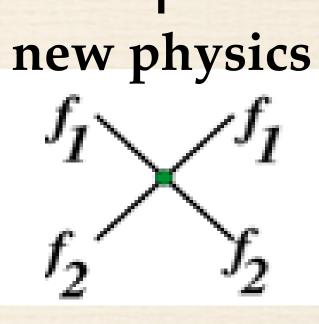


The Parity Violation Program at JLab

 $\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} \left[\overline{e} \gamma^{\mu} \gamma_5 e(C_{1u} \overline{u} \gamma_{\mu} u + C_{1d} \overline{d} \gamma_{\mu} d) \right] \begin{pmatrix} C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1d} = -\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W \approx -0.35 \end{pmatrix}$ $+\overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_{5}u+C_{2d}\overline{d}\gamma_{\mu}\gamma_{5}d)] C_{2u} = -\frac{1}{2}+2\sin^{2}\theta_{W} \approx -0.04$

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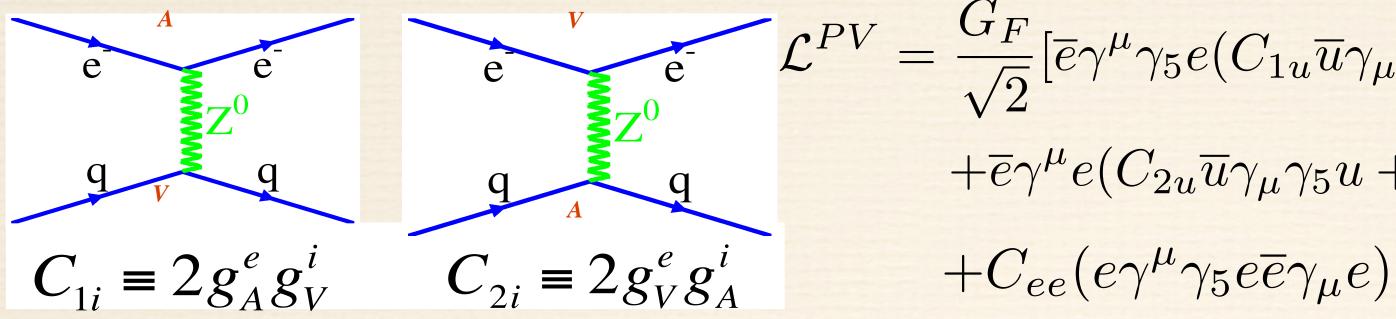
PV elastic e-N scattering, **Atomic parity violation**



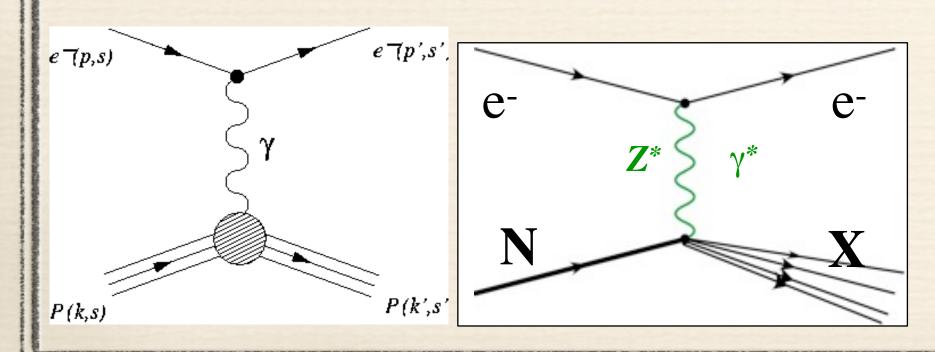
 \mathcal{L}_{f_1,f_2}

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma_\mu f_{2j}$





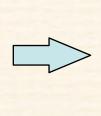
 $C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$ $C_{2q} \propto (g_{RR}^{eq})^2 - (g_{RL}^{eq})^2 + (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$



The Parity Violation Program at JLab

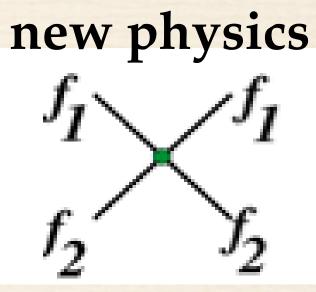
 $\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} \left[\overline{e} \gamma^{\mu} \gamma_5 e(C_{1u} \overline{u} \gamma_{\mu} u + C_{1d} \overline{d} \gamma_{\mu} d) \right] \begin{pmatrix} C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \end{pmatrix}$ $+\overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_{5}u+C_{2d}\overline{d}\gamma_{\mu}\gamma_{5}d)] C_{2u} = -\frac{1}{2}+2\sin^{2}\theta_{W} \approx -0.04$

 $C_{1d} = \frac{1}{2} - \frac{2}{3} \sin^2 \theta_W \approx 0.35$ $C_{2d} = \frac{1}{2} - 2 \sin^2 \theta_W \approx$



PV elastic e-N scattering, **Atomic parity violation**

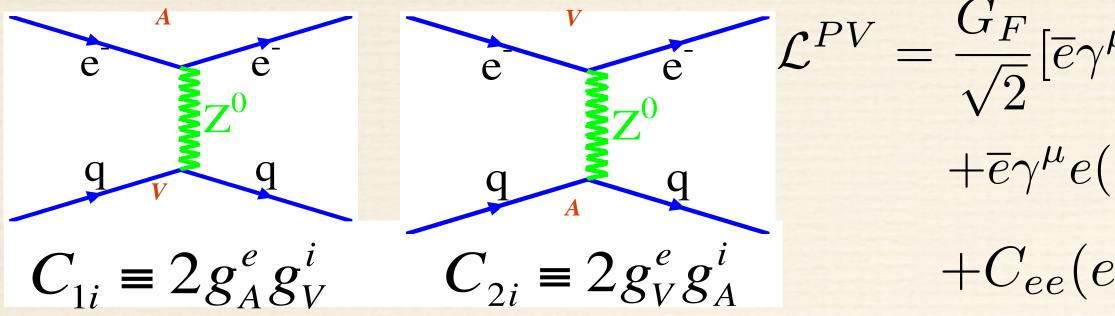
PV deep inelastic scattering



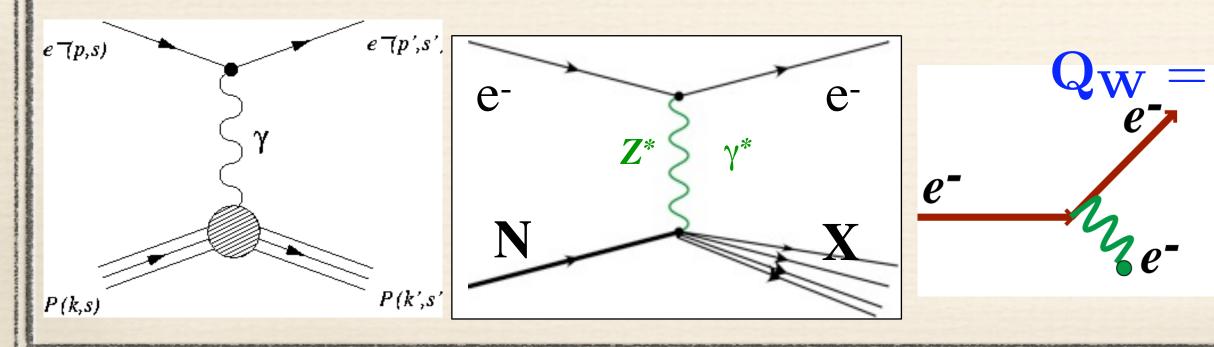
 $\mathcal{L}_{f_1f_2}$

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_{\mu} f_{1i} \bar{f}_{2j} \gamma_{\mu} f_{2j}$





 $C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$ $C_{2q} \propto (g_{RR}^{eq})^2 - (g_{RL}^{eq})^2 + (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$ $C_{ee} \propto (g_{RR}^{ee})^2 - (g_{LL}^{ee})^2 \implies \text{PV Møller scattering}$



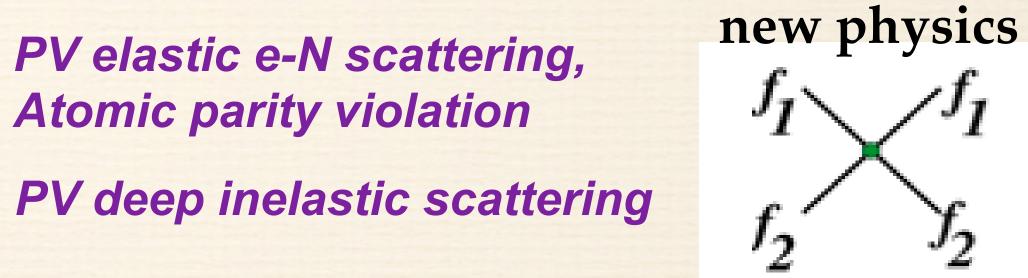
The Parity Violation Program at JLab

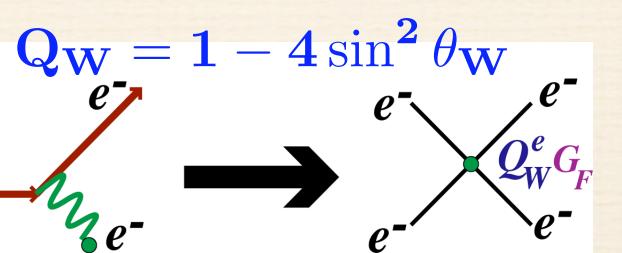
WNC Couplings

 $\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} \left[\overline{e} \gamma^{\mu} \gamma_5 e(C_{1u} \overline{u} \gamma_{\mu} u + C_{1d} \overline{d} \gamma_{\mu} d) \right] \begin{pmatrix} C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.25 \\ C_{1u} = -\frac{1}{2} + \frac{1}{3} + \frac$ $+\overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_{5}u+C_{2d}\overline{d}\gamma_{\mu}\gamma_{5}d)] C_{2u} = -\frac{1}{2}+2\sin^{2}\theta_{W} \approx -0.04$ $+C_{ee}(e\gamma^{\mu}\gamma_{5}e\overline{e}\gamma_{\mu}e)$

 $C_{1d} = \frac{1}{2} - \frac{2}{3} \sin^2 \theta_W \approx 0.35$ $C_{2d} = \frac{1}{2} - 2 \sin^2 \theta_W \approx$

PV elastic e-N scattering, **Atomic parity violation**

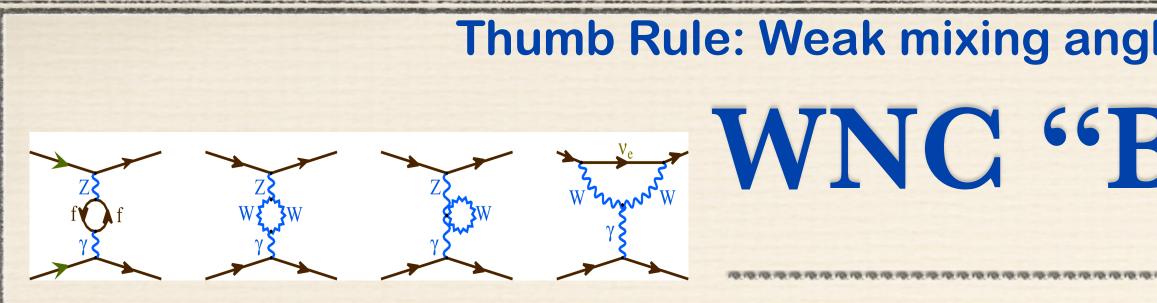




 $(g_{i\,j}^{12})^2$ $f_{1i}\gamma_{\mu}f_{1i}f_{2j}\gamma_{\mu}f_{2j}$ i, j = L, R

 $\mathcal{L}_{f_1f_2}$





Electroweak Radiative Corrections causes weak mixing angle to "run"

Atomic Parity Violation: Cs-133
future measurements and theory challenging
Neutrino Deep Inelastic Scattering:
future measurements and theory challenging
PV Møller Scattering: E158 at SLAC
statistics limited, theory robust
next generation: MOLLER (factor of 5 better)
PV elastic e-p scattering: Qweak
theory robust at low beam energy
next generation: P2 (factor of 3 better)
PV Deep Inelastic Scattering: PVDIS
theory robust for ² H in valence quark region
factor of 5 to 8 improvement possible: SOLID

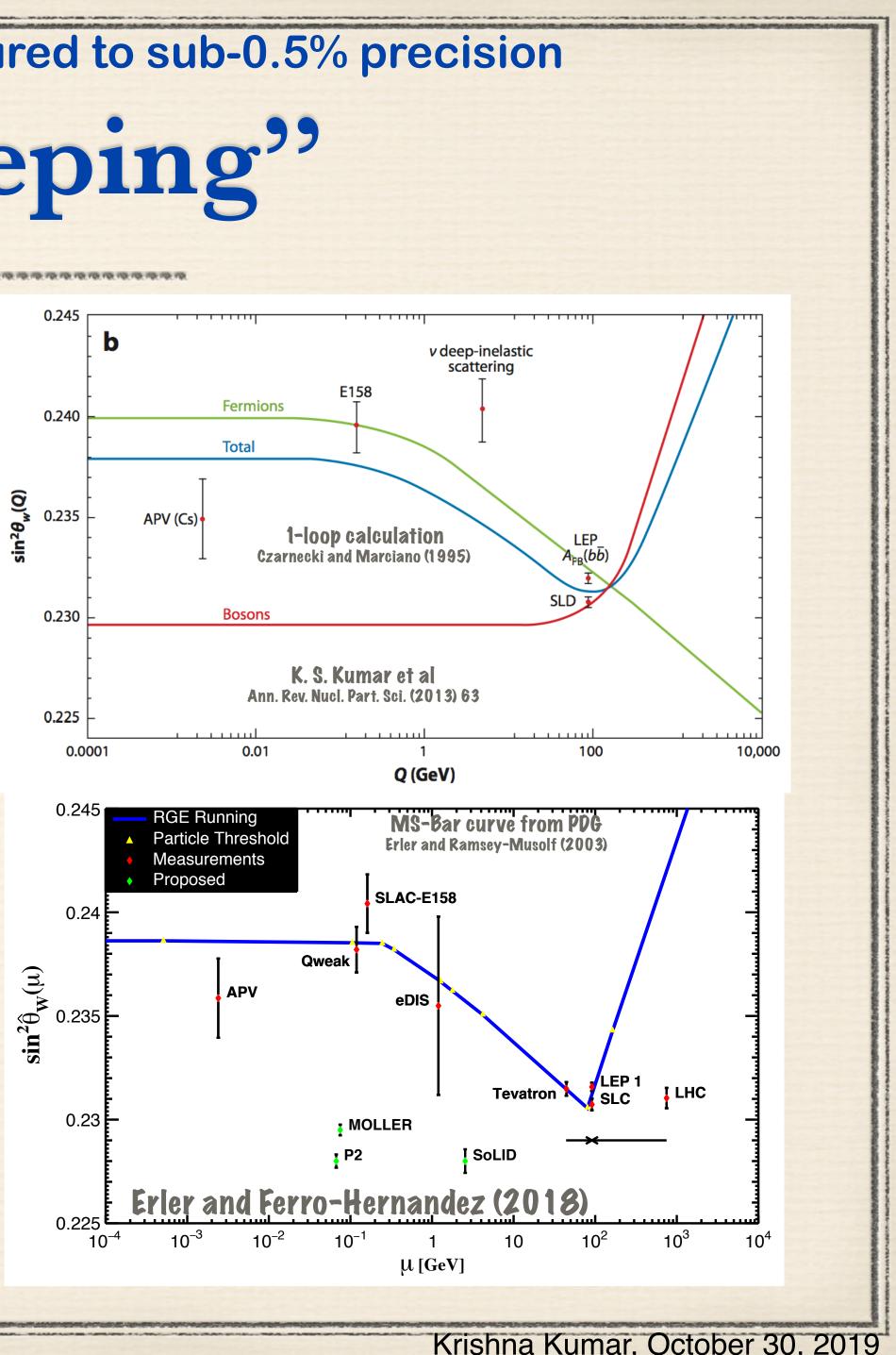
The Parity Violation Program at JLab

Thumb Rule: Weak mixing angle must be measured to sub-0.5% precision

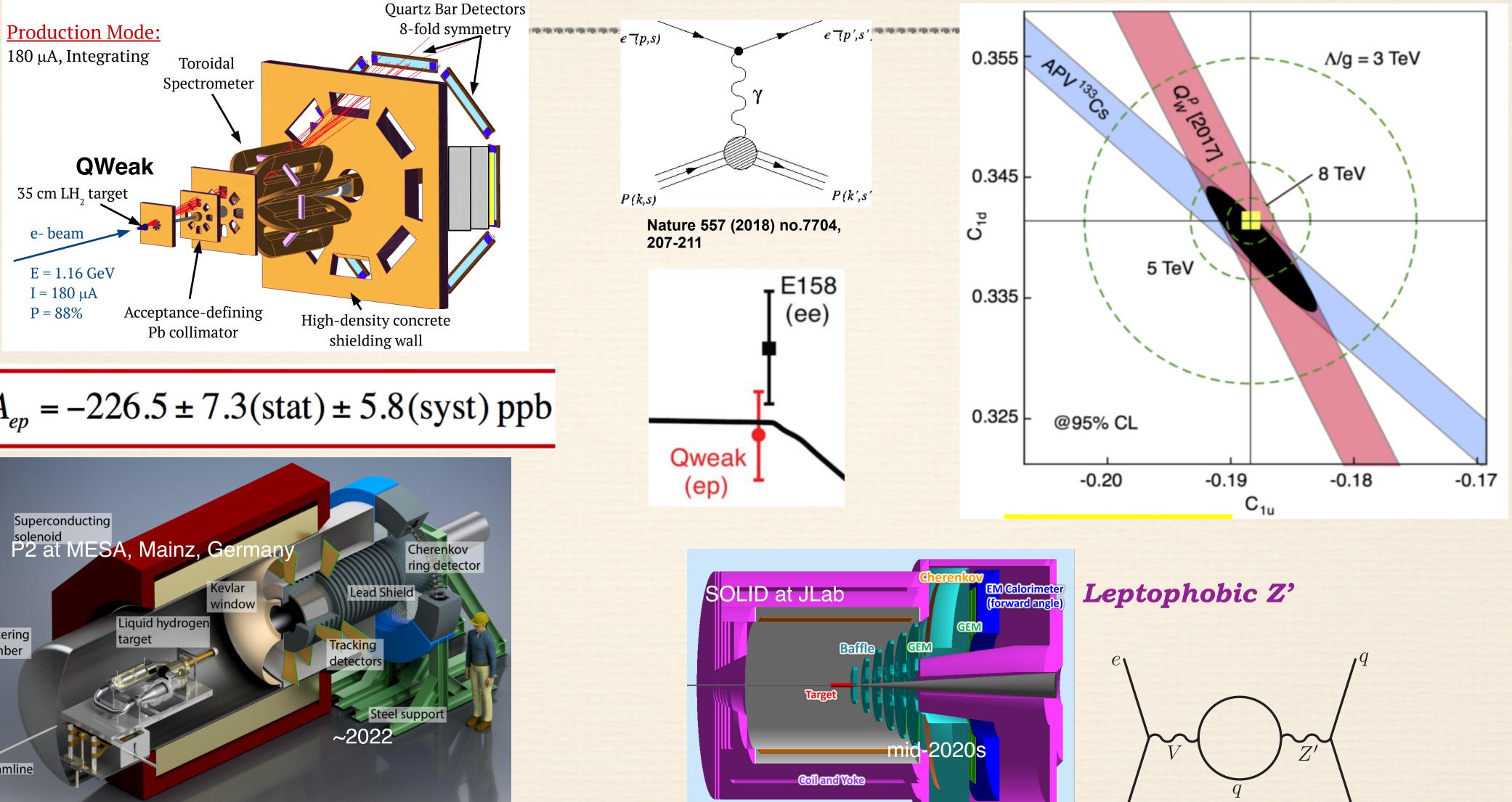
WNC "Bookkeeping"

NuTeV

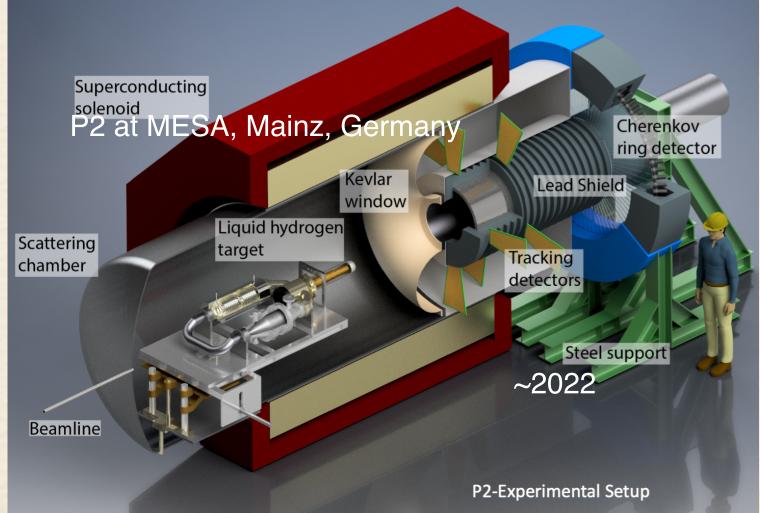




Semi-Leptonic: Recent Past and Future

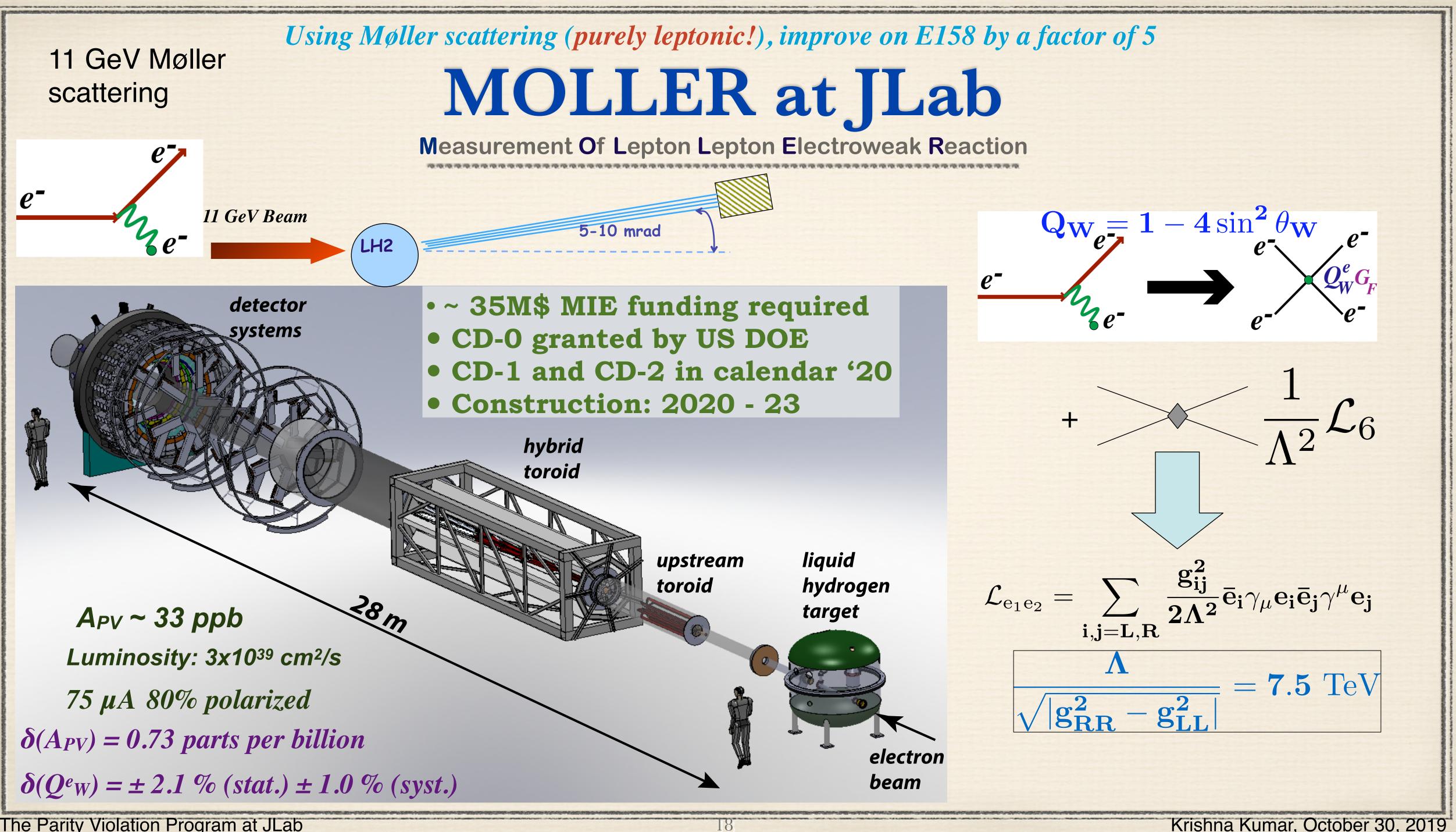


$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$



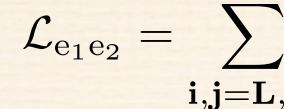
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95% C. L. Reach **Comparison with e⁺e⁻ Collisions**

Best reach on purely leptonic contact interaction amplitudes: LEP200



$g_{ij} = 4\pi\eta_{ij}$						
Model	η^f_{LL}	η_{RR}^{f}	η_{LR}^f	η^f_{RL}		
LL^{\pm}	± 1	0	0	0		
RR^{\pm}	0	± 1	0	0		
VV^{\pm}	± 1	± 1	± 1	± 1		

MOLLER is accessing discovery space that cannot be reached until the advent of a new lepton collider or neutrino factory

The Parity Violation Program at JLab

$$\sum_{\mathbf{q},\mathbf{R}} rac{\mathbf{g}_{\mathbf{ij}}^2}{\mathbf{2}\Lambda^2} \mathbf{\bar{e}}_{\mathbf{i}} \gamma_{\mu} \mathbf{e}_{\mathbf{i}} \mathbf{\bar{e}}_{\mathbf{j}} \gamma^{\mu} \mathbf{e}_{\mathbf{j}}$$

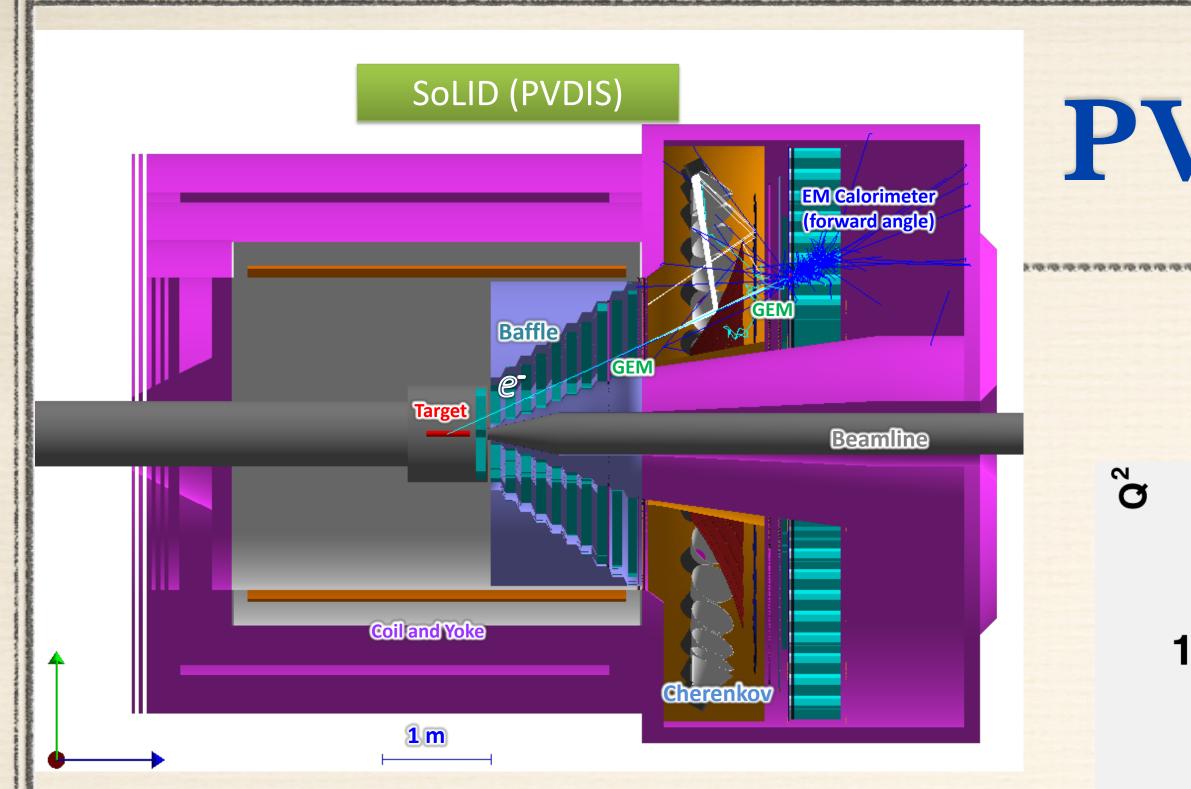
 $\Lambda^{
m ee}_{
m LL}\sim 8.3~{
m TeV}$ **LEP200 Reach**

E158 Reach

 $\Lambda^{
m ee}_{
m LL} \sim 12 \; {
m TeV}$

 $\Lambda^{
m ee}_{
m LL}\sim 27~{
m TeV}$ **MOLLER Reach**



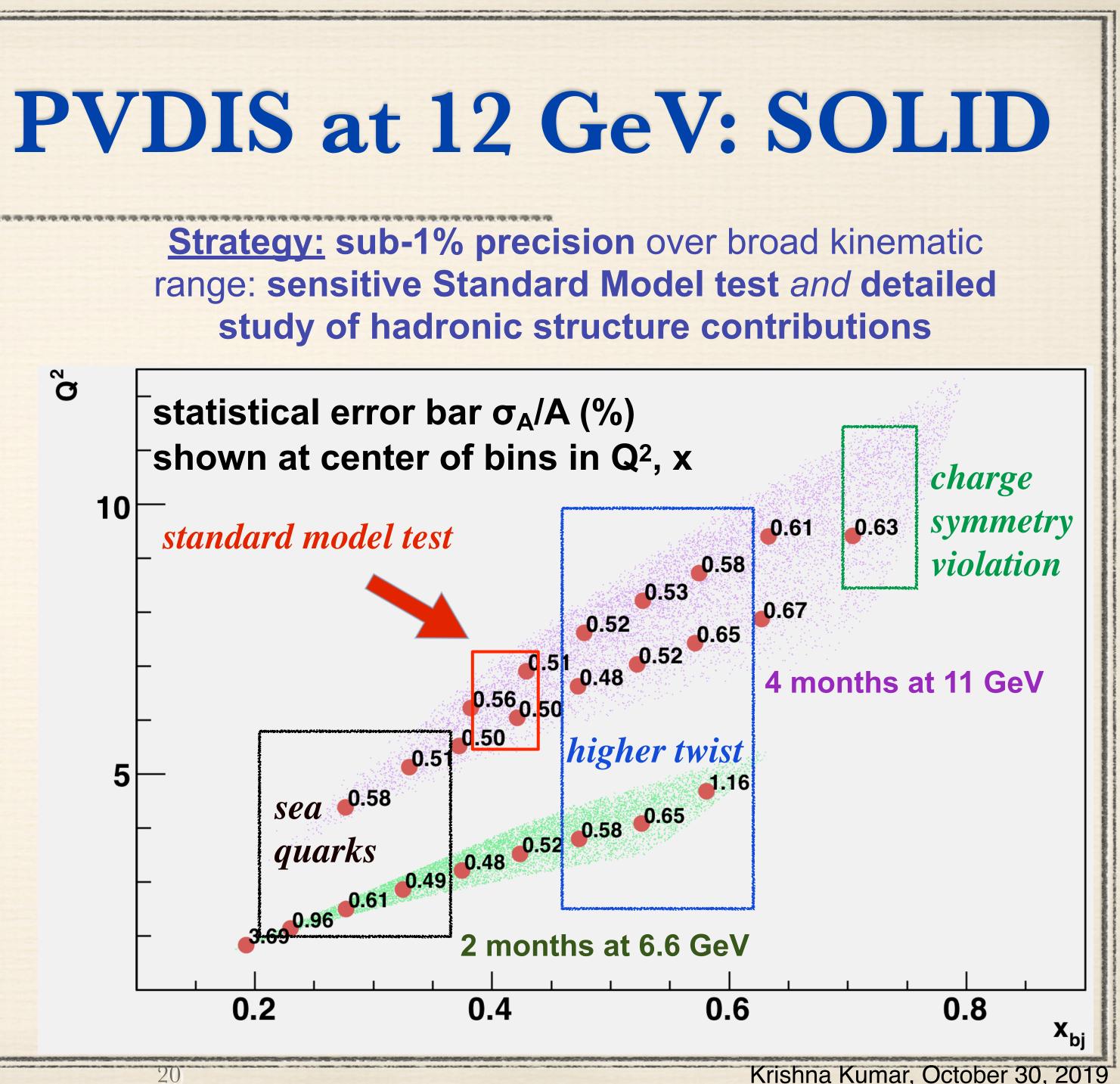


Requirements

- High Luminosity with E > 10 GeV
- Large scattering angles (for high x & y)
- Better than 1% errors for small bins
- x-range 0.25-0.75
- $W^2 > 4 \, GeV^2$
- Q² range a factor of 2 for each x -(Except at very high x)
- Moderate running times

The Parity Violation Program at JLab

Strategy: sub-1% precision over broad kinematic range: sensitive Standard Model test and detailed study of hadronic structure contributions



 $A_{PV} = \frac{G_F Q^2}{\sqrt{2}\pi\alpha} \left[a(x) + f(y)b(x) \right] \text{ first principles: using electroweak} \text{ } A_{PV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_3^{\gamma Z}}{F_1^{\gamma}} \right]$

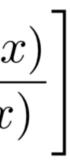


 $A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \left[a(x) + f(y)b(x) \right] \text{ first principles: using electroweak} \quad A_{PV} = \frac{G_F Q^2}{2\sqrt{2\pi\alpha}} \left[g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_3^{\gamma Z}}{F_1^{\gamma}} \right]$

quark-parton model and pdfs

$$a(x) \approx \frac{3}{4} \left[\frac{6C_{1u}u(x) - 3C_{1d}d(x)}{u(x) + \frac{1}{4}d(x)} \right] \sim \left[\frac{u(x) + 0.912d(x)}{u(x) + 0.25d(x)} \right]$$

The Parity Violation Program at JLab





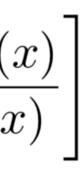
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 Nucleon scattering Slightly lower-x

 Statistics-dominated! • "For free": just replace LD2 with LH2



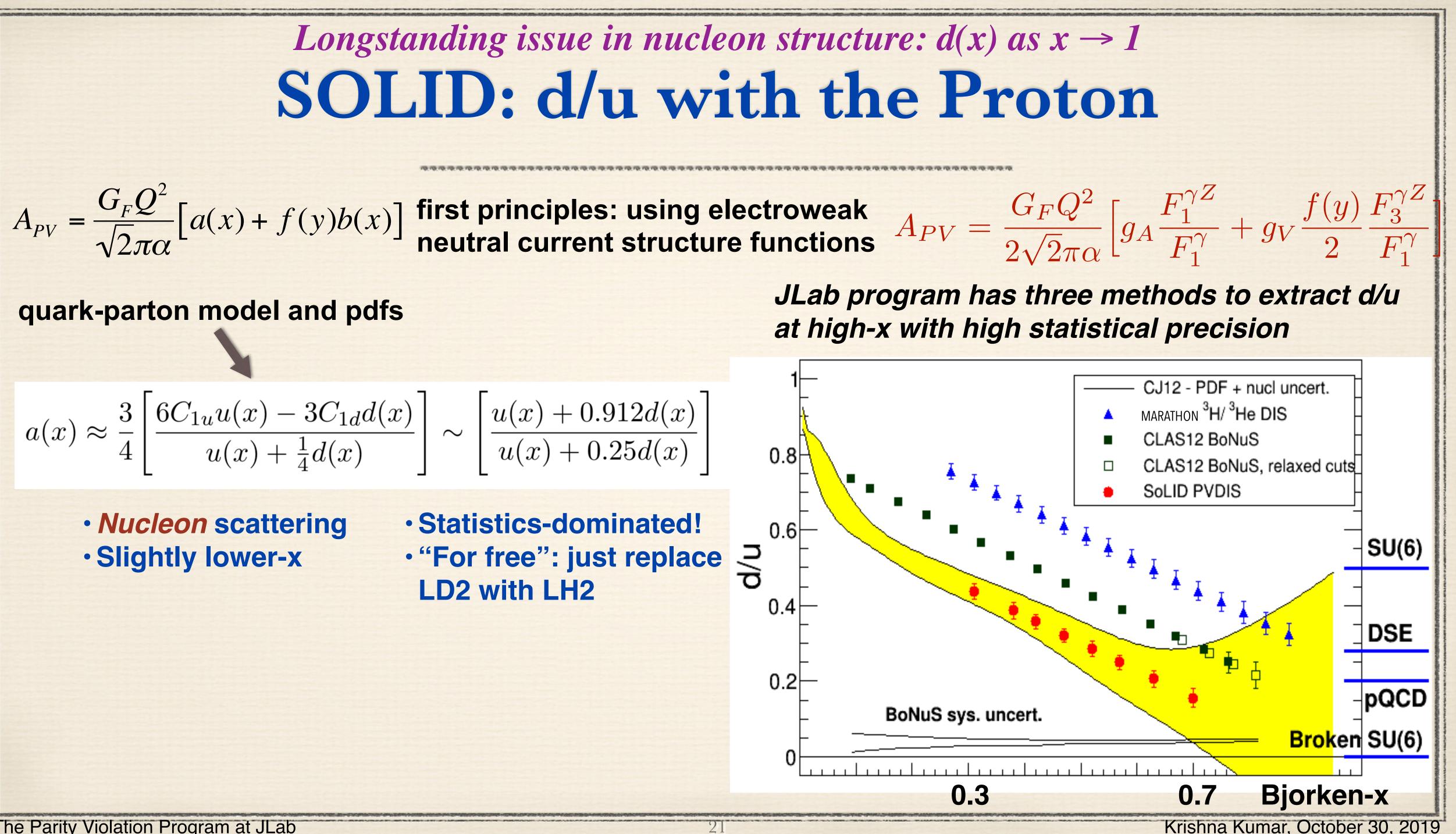


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$$A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \left[a(x) + f(y)b(x) \right]$$
 first print neutral of the formula of the second strength ot

quark-parton model and pdfs

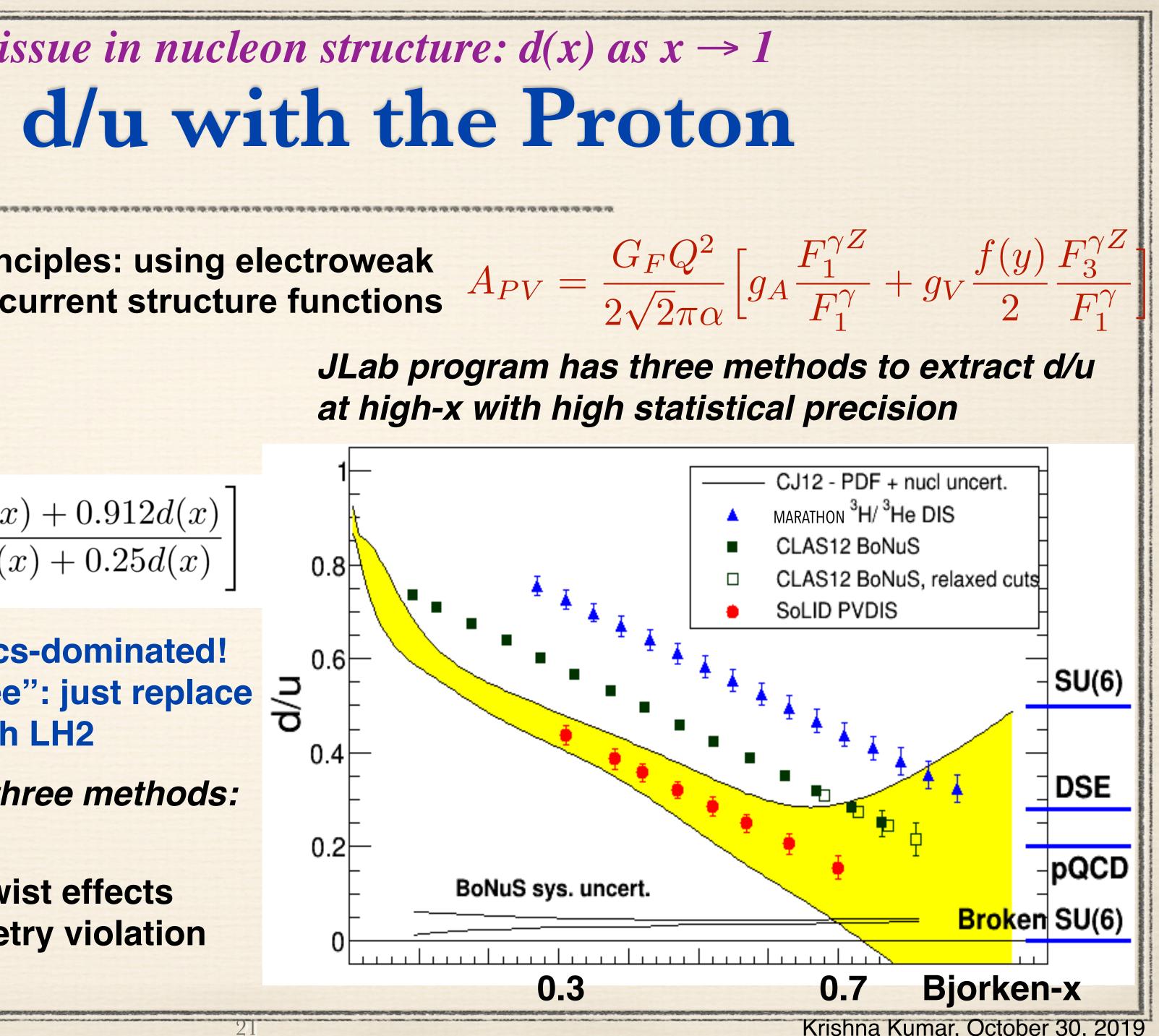
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 Statistics-dominated! Nucleon scattering • "For free": just replace Slightly lower-x LD2 with LH2

By comparing and contrasting the three methods:

- Robust extraction of d/u from JLab
- Possible new Insights into higher twist effects
- Possibly disentangle charge symmetry violation
- Nuclear dynamics at high-x

The Parity Violation Program at JLab



Cloet, Bentz, Thomas, PRL 109 (2012) 182301 48Ca PVDIS **Consider PVDIS on a heavy nucleus**

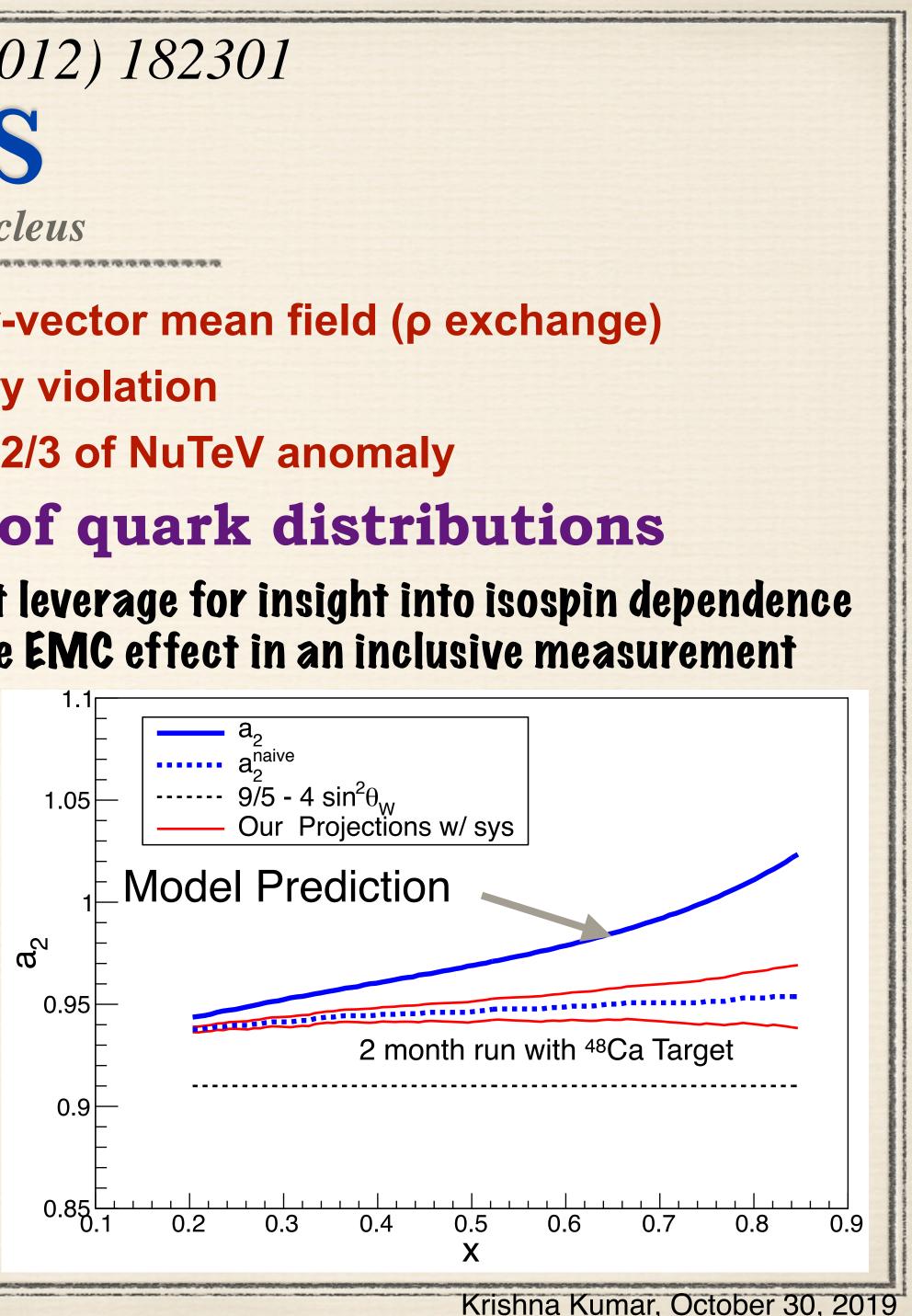
- shifts quark distributions: "apparent" charge symmetry violation
- Isovector EMC effect: could be responsible for at least 2/3 of NuTeV anomaly

$$a_2 \simeq \frac{9}{5} - 4\sin^2\theta_W - \frac{12}{25}\frac{u_A^+ - d_A^+}{u_A^+ + d_A^+} + \dots$$

Neutron or proton excess in nuclei leads to a isovector-vector mean field (p exchange)

new insight into medium modification of quark distributions

Great leverage for insight into isospin dependence of the EMC effect in an inclusive measurement



Cloet, Bentz, Thomas, PRL 109 (2012) 182301 48Ca PVDIS **Consider PVDIS on a heavy nucleus**

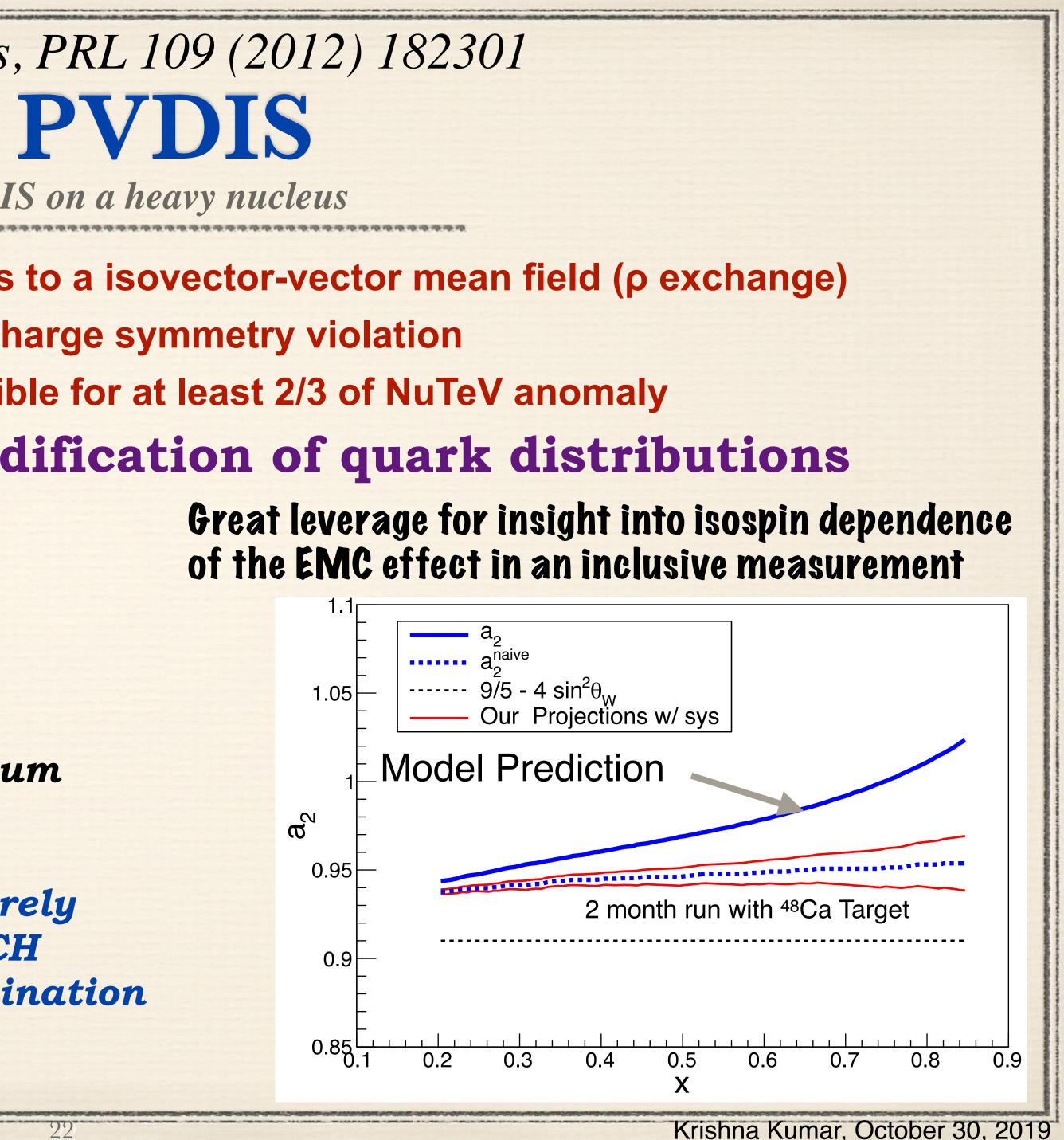
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$$a_2 \simeq \frac{9}{5} - 4\sin^2\theta_W - \frac{12}{25}\frac{u_A^+ - d_A^+}{u_A^+ + d_A^+} + .$$

- •methods of flavor decomposition of medium modifications challenging • must disentangle small effects
- Precise isotope cross-section ratios in purely electromagnetic electron scattering: MUCH reduced sensitivity to the isovector combination

Neutron or proton excess in nuclei leads to a isovector-vector mean field (p exchange)

new insight into medium modification of quark distributions



High luminosity: precision measurements of PV observables EIC Structure Functions $e^{-} \rightarrow e^{-} H, ^{2}H, ^{3}He$

The core physics topics of the EIC have driven designs that reach a new regime of extraordinarily high polarized luminosity, state-of-the-art collider detector technology and precision polarimetry

$$\begin{split} \frac{1}{2m_N} W^i_{\mu\nu} &= -\frac{g_{\mu\nu}}{m_N} F^i_1 + \frac{p_\mu p_\nu}{m_N (p \cdot q)} F^i_2 \stackrel{\text{Ji, Vogelsang, Blümlein, ...}}{\text{Anselmino, Efremov \& Leader, Phys. Rep. 261 (1995)} \\ &+ i \frac{\epsilon_{\mu\nu\alpha\beta}}{2(p \cdot q)} \left[\frac{p^\alpha q^\beta}{m_N} F^i_3 + 2q^\alpha S^\beta g^i_1 - 4xp^\alpha S^\beta g^i_2 \right] \\ &- \frac{p_\mu S_\nu + S_\mu p_\nu}{2(p \cdot q)} g^i_3 + \frac{S \cdot q}{(p \cdot q)^2} p_\mu p_\nu g^i_4 + \frac{S \cdot q}{p \cdot q} g_{\mu\nu} g^i_5 \end{split}$$



High luminosity: precision measurements of PV observables **EIC Structure Functions** $e^{-} \longrightarrow \leftarrow ^{-1}H, ^{2}H, ^{3}He$ The core physics topics of the EIC have driven designs that reach a new regime of extraordinarily high polarized luminosity, state-of-the-art collider detector technology and precision polarimetry polarized electron, unpolarized hadron deuteron proton $A_{PV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_3^{\gamma Z}}{F_1^{\gamma}} \right]$ unpolarized electron, polarized hadron $A_{TPV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[g_V \frac{g_5^{\gamma Z}}{F_1^{\gamma}} + g_A f(y) \frac{g_1^{\gamma Z}}{F_1^{\gamma}} \right]$ The Parity Violation Program at JLab

$$\begin{split} \frac{1}{2m_N}W^i_{\mu\nu} &= -\frac{g_{\mu\nu}}{m_N}\,F^i_1 + \frac{p_\mu p_\nu}{m_N(p\cdot q)}\,F^i_2 \stackrel{\text{Ji, Vogelsang, Blümlein, ...}}{\text{Anselmino, Efremov \& Leader, Phys. Rep. 261 (1995)} \\ &+ i\frac{\epsilon_{\mu\nu\alpha\beta}}{2(p\cdot q)} \left[\frac{p^\alpha q^\beta}{m_N}\,F^i_3 + 2q^\alpha S^\beta\,g^i_1 - 4xp^\alpha S^\beta\,g^i_2\right] \\ &- \frac{p_\mu S_\nu + S_\mu p_\nu}{2(p\cdot q)}\,g^i_3 + \frac{S\cdot q}{(p\cdot q)^2}\,p_\mu p_\nu\,g^i_4 + \frac{S\cdot q}{p\cdot q}\,g_{\mu\nu}\,g^i_5 \end{split}$$

$$egin{aligned} F_1^{\gamma Z} &\propto u+d+s \ F_3^{\gamma Z} &\propto 2u_v+d_v \ g_1^{\gamma Z} &\propto \Delta u+\Delta d+\Delta s \ g_5^{\gamma Z} &\propto 2\Delta u_v+\Delta d_v \end{aligned}$$

 $F_1^{\gamma Z} \propto u + d + 2s$ $F_3^{\gamma Z} \propto u_v + d_v$ $g_1^{\gamma Z} \propto \Delta u + \Delta d + \Delta s$ $g_5^{\gamma Z} \propto \Delta u_v + \Delta d_v$



High luminosity: precision measurements of PV observables **EIC Structure Functions** $e^{-} \longrightarrow \leftarrow ^{-1}H, ^{2}H, ^{3}He$ The core physics topics of the EIC have driven designs that reach a new regime of extraordinarily high polarized luminosity, state-of-the-art collider detector technology and precision polarimetry polarized electron, unpolarized hadron proton $A_{PV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_3^{\gamma Z}}{F_1^{\gamma}} \right]$ unpolarized electron, polarized hadron $A_{TPV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[g_V \frac{g_5^{\gamma Z}}{F_1^{\gamma}} + g_A f(y) \frac{g_1^{\gamma Z}}{F_1^{\gamma}} \right]$ First measurements of novel electroweak spin structure functions 6-flavor separation of polarized quark pdfs using just inclusive measurements exploit both polarized proton and deuteron (or helium-3) weak mixing angle (averaging for deuteron polarization) at high Q²

The Parity Violation Program at JLab

$$\begin{split} \frac{1}{2m_N}W^i_{\mu\nu} &= -\frac{g_{\mu\nu}}{m_N}\,F^i_1 + \frac{p_\mu p_\nu}{m_N(p\cdot q)}\,F^i_2 \stackrel{\text{Ji, Vogelsang, Blümlein, ...}}{\text{Anselmino, Efremov \& Leader, Phys. Rep. 261 (1995)} \\ &+ i\frac{\epsilon_{\mu\nu\alpha\beta}}{2(p\cdot q)} \left[\frac{p^\alpha q^\beta}{m_N}\,F^i_3 + 2q^\alpha S^\beta\,g^i_1 - 4xp^\alpha S^\beta\,g^i_2\right] \\ &- \frac{p_\mu S_\nu + S_\mu p_\nu}{2(p\cdot q)}\,g^i_3 + \frac{S\cdot q}{(p\cdot q)^2}\,p_\mu p_\nu\,g^i_4 + \frac{S\cdot q}{p\cdot q}\,g_{\mu\nu}\,g^i_5 \end{split}$$

deuteron

$$F_1^{\gamma Z} \propto u + d + s$$

$$F_3^{\gamma Z} \propto 2u_v + d_v$$

$$g_1^{\gamma Z} \propto \Delta u + \Delta d + \Delta s$$

$$g_5^{\gamma Z} \propto 2\Delta u_v + \Delta d_v$$

 $F_1^{\gamma Z} \propto u + d + 2s$ $F_3^{\gamma Z} \propto u_v + d_v$ $g_1^{\gamma Z} \propto \Delta u + \Delta d + \Delta s$ $g_5^{\gamma Z} \propto \Delta u_v + \Delta d_v$



Summary and Outlook

Parity-Violating Electron Scattering * Enabled unique studies of the weak force * Technical progress has enabled unprecedented precision * flagship experiments at electron accelerators Fundamental Nuclear/Nucleon Physics * Neutron RMS radii of heavy nuclei (PREX, CREX, MREX...) * valence guark structure of protons and neutrons (SOLID) Fundamental Electroweak Physics * Search for new dynamics at the TeV scale (P2, MOLLER, SOLID) complementary to colliders; would help interpret potential anomalies precision measurement of the weak mixing angle

Logical progression of this physics into the EIC

Exciting times for the next decade and beyond!



MOLLER Context Summary

best contact interaction reach for leptons at low OR high energy: similar to LHC reach with semi-leptonic amplitudes To do better for a 4-lepton contact interaction would require: Giga-Z factory, linear collider, neutrino factory or muon collider

 $\delta(\sin^2 \theta_W) = \pm 0.00024 \ (stat.) \pm 0.00013 \ (syst.)$ $\square > \sim 0.1\%$

If LHC sees ANY anomaly in Runs 2 or 3 (~2022)

other sensitive probes (e.g. g-2 anomaly)

Discovery scenarios beyond LHC signatures

- ★ Purely Leptonic Contact Interactions
- ★ Lepton Number Violating Amplitudes
- ★ Light Dark Matter Mediators
- ★ Lorentz Violation

-

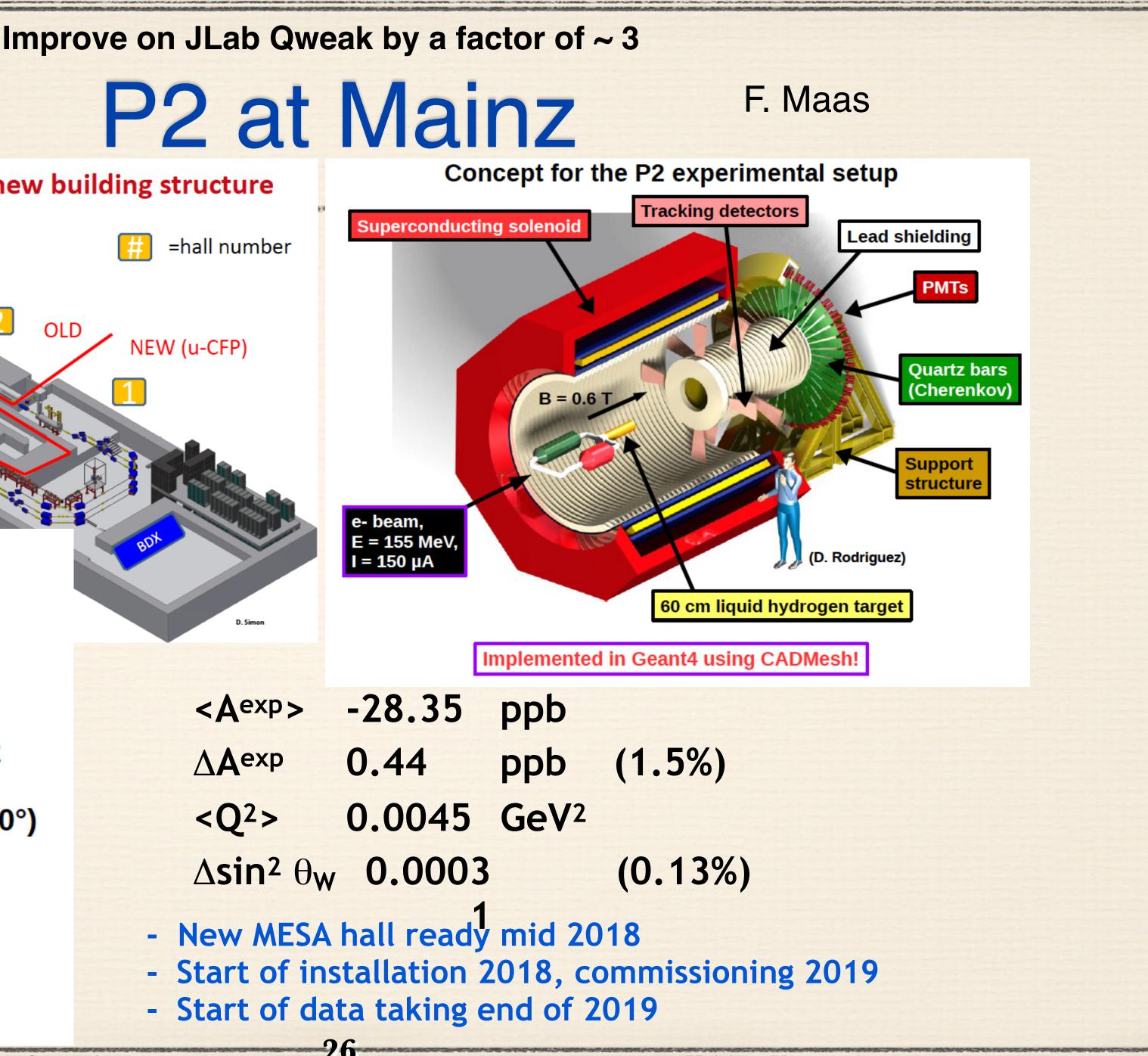


- Best projected uncertainty among projects being considered over next 10 years
 - * The unique discovery space probed by MOLLER will become a pressing need, like



MESA in old/new building structure 11m =hall number 2 OLD NEW (u-CFP) P2 Experiment: external beam, 155 MeV, Beam energy: 155 MeV Beam current: 150 µA Polarization: (85±0.5)% <Aexp> 60 cm IH2 Target: ∆Aexp <Q2> 2π·(35°±10°) Acceptance: 0.5 THz Rate: Runtime: 10000 h ΔA^{app}: 0.1 ppb

Physics with Parity-Violating Electron Scattering

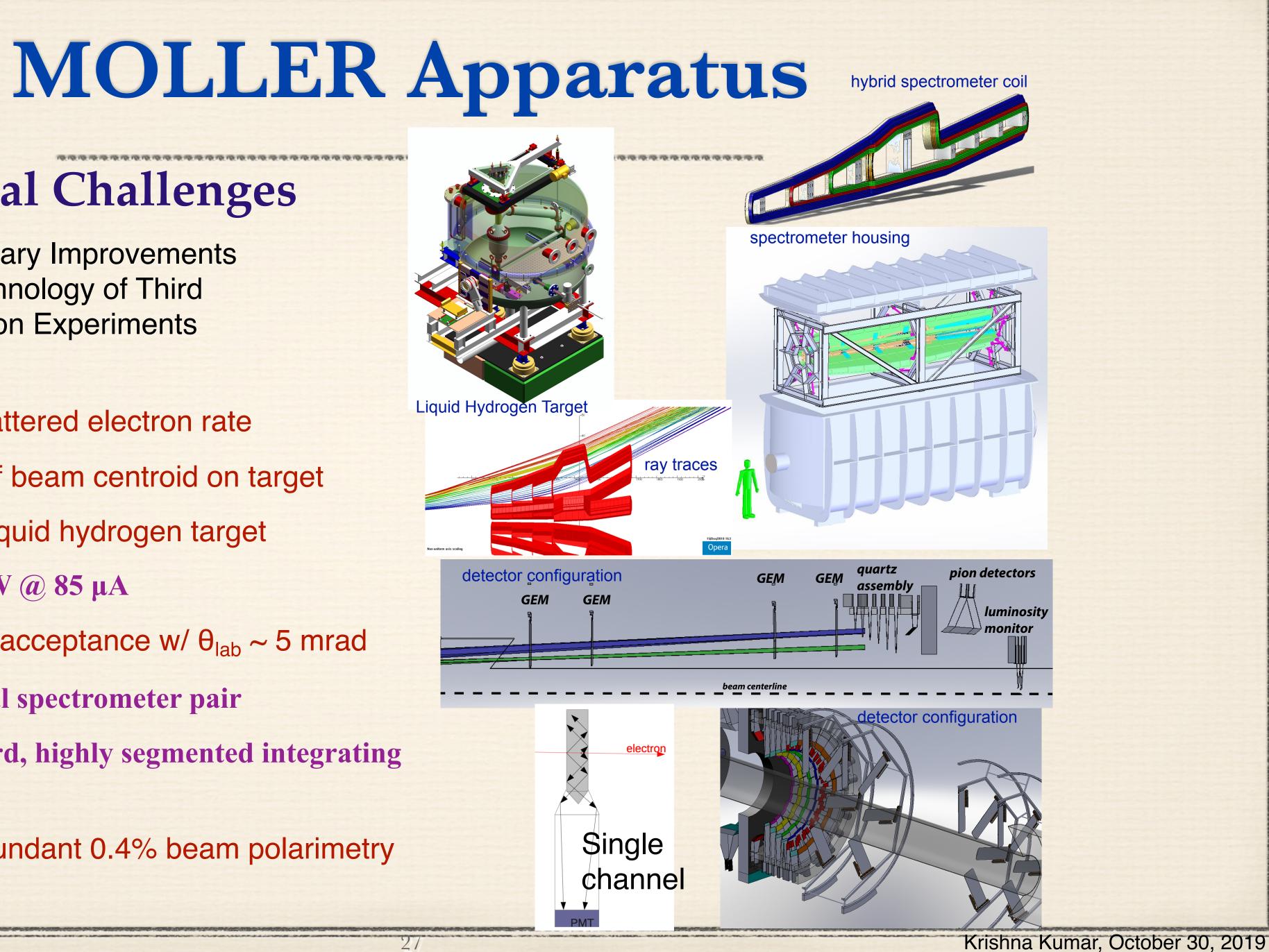




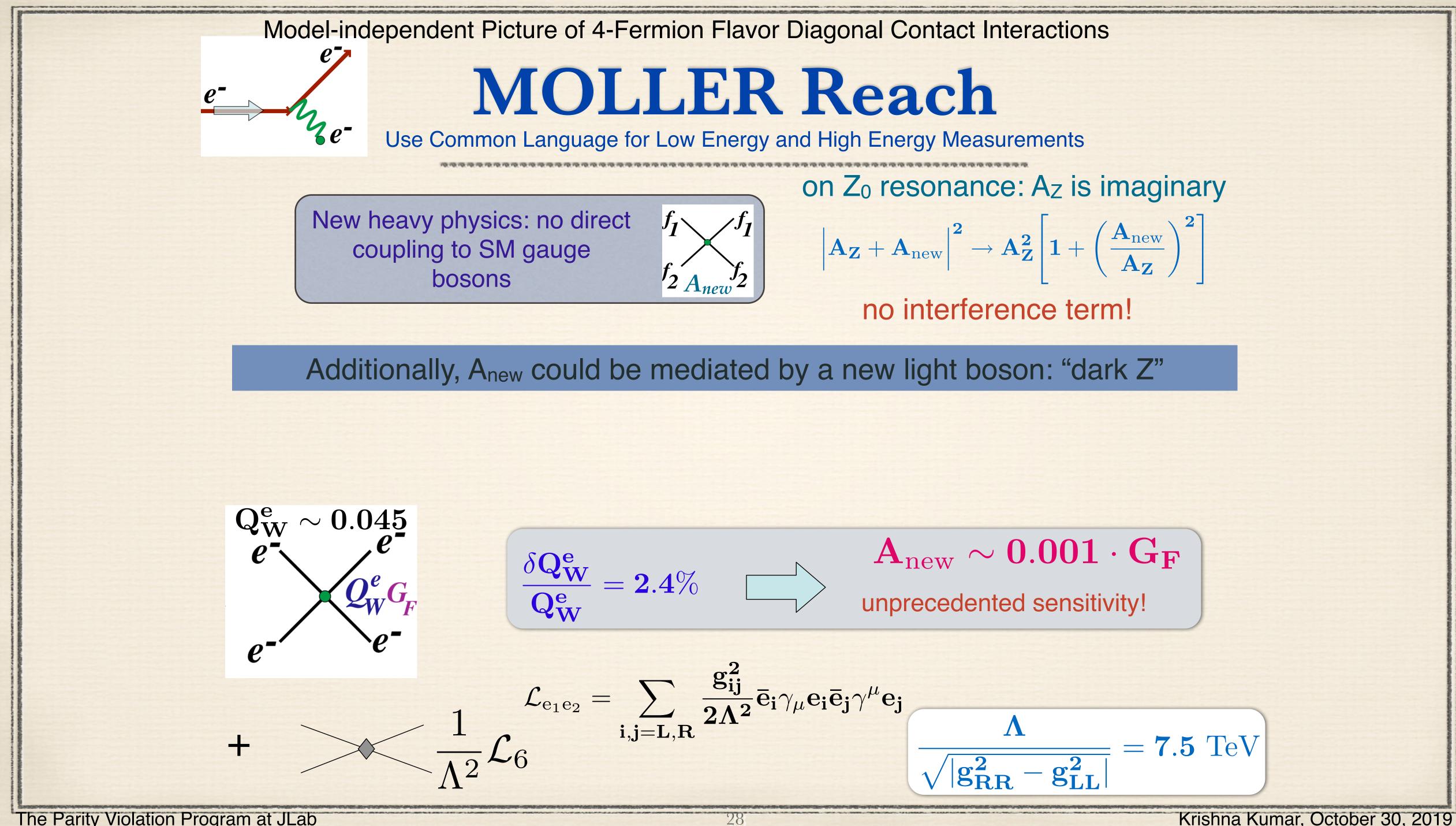
Technical Challenges

Evolutionary Improvements from Technology of Third **Generation Experiments**

- ~ 150 GHz scattered electron rate
- 1 nm control of beam centroid on target
- > 10 gm/cm² liquid hydrogen target
 - -1.5 m: ~5 kW @ 85 µA
- Full Azimuthal acceptance w/ $\theta_{lab} \sim 5$ mrad
 - novel toroidal spectrometer pair
 - radiation hard, highly segmented integrating detectors
- Robust & Redundant 0.4% beam polarimetry







$$\left|\mathbf{A_Z}+\mathbf{A_{new}}
ight|^2
ightarrow\mathbf{A_Z^2}\left[\mathbf{1}+\left(rac{\mathbf{A_{new}}}{\mathbf{A_Z}}
ight)^2
ight]$$



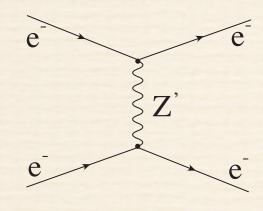
Unique Opportunity: Purely Leptonic Reaction at Q² << M_Z²

New Physics Examples

Deviations From Theory Prediction Interpretable as New Physics

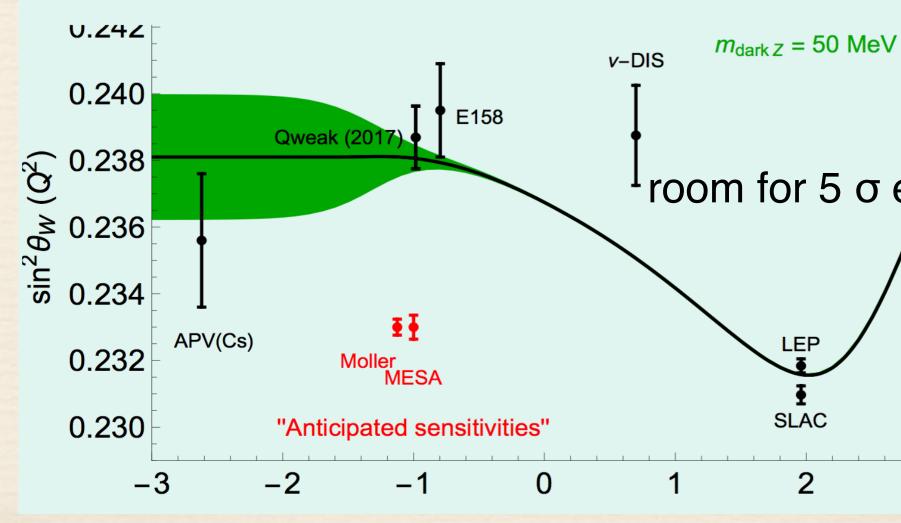
e

Many different scenarios give rise to effective 4-electron contact interaction amplitudes: significant discovery potential



Heavy Photons (A' mixed with Z₀): The Dark Z

H. Davoudiasl, H-S. Lee and W. Marciano



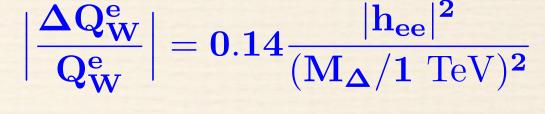
The Parity Violation Program at JLab

Lepton Number Violation

Doubly-Charged Scalar

H

 \sim



5 σ for h_{ee} ~ 1 and M_{Δ} ~ 1 TeV

See Michael's talk! (and send me the slides!)

Constraining Lorentz Invariance

Ralf Lehnert, J. Phys.: Conf. Ser. 952 (2018) 012008

$$\begin{aligned} \mathbf{\sigma} \, \text{effe} \, \delta A(t) &= \frac{G_F}{\sqrt{2}\pi\alpha} \, \frac{E_k \, y \, (1-y) \sin^2 \theta_W}{(y^2 - y + 1)^2} \, \vec{k}(t) \cdot \vec{\xi} \\ &= \frac{G_F}{\sqrt{2}\pi\alpha} \, \frac{E_k^2 \, y \, (1-y) \sin^2 \theta_W}{(y^2 - y + 1)^2} \times \\ & \left[\sqrt{\xi_X^2 + \xi_Y^2} \sqrt{1 - \cos^2 \alpha \sin^2 \chi} \, \cos \Omega_{\oplus} t + c_0 \right] \end{aligned}$$

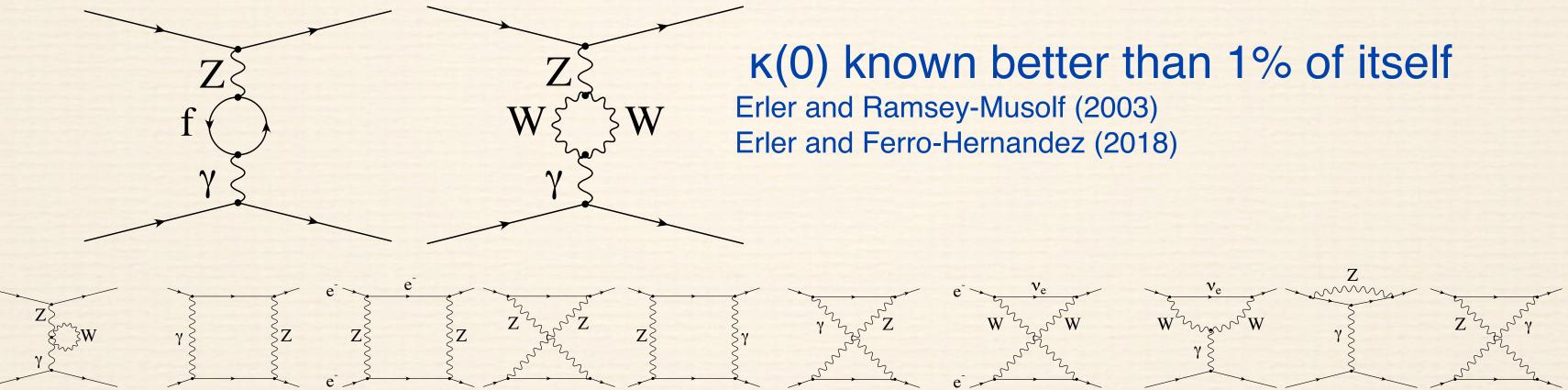
3



Electroweak Theory

EW Theory Prediction Uncertainty Well Below Projected Experimental Uncertainty

Dominant Contribution at 1-loop



 $\delta(Q^e_W)$ (theory) = 0.6%, another factor of 2 improvement with full two-loop calculation

MOLLER $\delta(Q^{e_{W}})$ goal = ± 2.1 % (stat.) ± 1.1 % (syst.)

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Unique Opportunity: Purely Leptonic Reaction at $Q^2 \ll M_Z^2$

Czarnecki and Marciano (1995)

 $A_{PV}(ee) \propto \rho G_F \left[1 - 4\kappa(0) \sin^2 \theta_W(m_Z)_{\overline{\text{MS}}} \right] + \cdots$

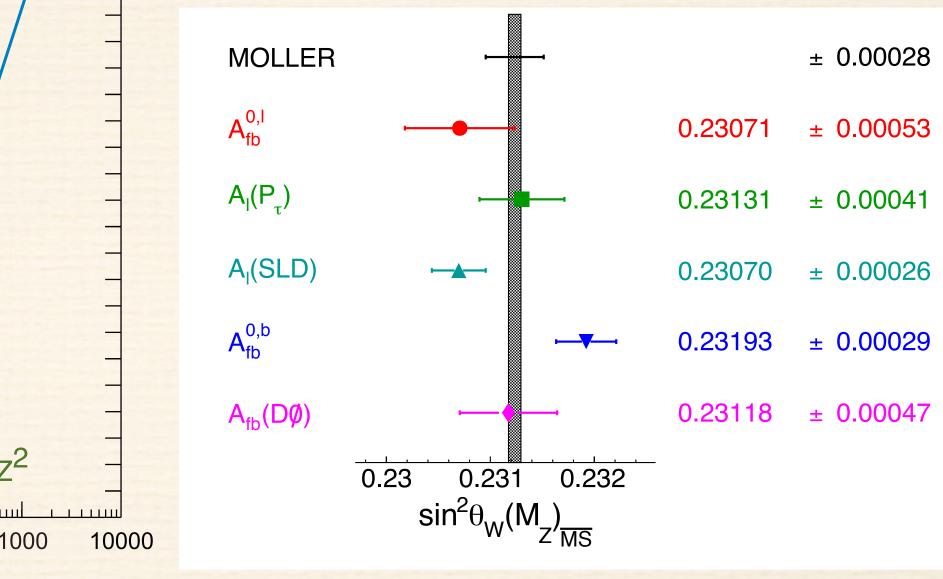
See talks by:

- A. Aleksejev
- A. Freitas
- R. Ferro Hernandez
- H. Patel
- M. Ramsey-Musolf



A Fundamental Parameter of the Electroweak Theory sin²0w MOLLER Projection: $\delta(sin^2\theta_W) = \pm 0.00024$ (stat.) ± 0.00013 (syst.) no interference term 0.245 MOLLER MOLLER NuTeV Q_w(e) goal 0.240 $A_{fb}^{0,l}$ 0.23071 $sin^2 \theta_W(\mu)$ $A_{I}(P_{\tau})$ 0.23131 Q_W(APV) 0.235 A_I(SLD) 0.23070 **------**EP $A_{fb}^{0,b}$ 0.230 SLD **MOLLER: 0.00028** $A_{fb}(D\emptyset)$ \pm 10 σ discovery potential at Q² << M_Z² 0.23 0.231 0.232 0.225 sin²θ_w(M_{z)} 1000 10000 0.0001 0.001 0.01 100 0.1 10 μ[GeV] Mainz P2: ~ 0.00032 **Future projections** Final Tevatron: ~ 0.00041 (similar time scale) LHC 14 TeV, 300 fb⁻¹ : ~ 0.00036 Note: systematics-dominated (pdf uncertainties) Old and New Physics with Electron-Electron Scattering

Z resonance measurements:





MOLLER Uncertainty Table

Beam	Assumed	Accuracy of	Required 2 kHz	Required cumulative	Systematic
Property	Sensitivity	Correction	random fluctuations	helicity-correlation	contribution
Intensity	1 ppb / ppb	~1%	< 1000 ppm	< 10 ppb	$\sim 0.1 \text{ ppb}$
Energy	-1.4 ppb / ppb	$\sim 10\%$	< 108 ppm	< 0.7 ppb	$\sim 0.05 \; \mathrm{ppb}$
Position	0.85 ppb / nm	$\sim 10\%$	$< 47~\mu{ m m}$	< 1.2 nm	$\sim 0.05 \text{ ppb}$
Angle	8.5 ppb / nrad	~10%	$< 4.7 \ \mu rad$	< 0.12 nrad	$\sim 0.05 \text{ ppb}$

Error Source	Fractional Error (%)
Statistical	2.1
Absolute Normalization of the Kinematic Factor	0.5
Beam (second order)	0.4
Beam polarization	0.4
$e + p(+\gamma) \rightarrow e + X(+\gamma)$ All systematics	0.4
Beam (position, angle, energy) required at	0.4
Beam (intensity) sub-1% level	0.3
$e + p(+\gamma) \rightarrow e + p(+\gamma)$	0.3
$\gamma^{(*)} + p \to (\pi, \mu, K) + X$	0.3
Transverse polarization	0.2
Neutral background (soft photons, neutrons)	0.1
Total systematic	1.1

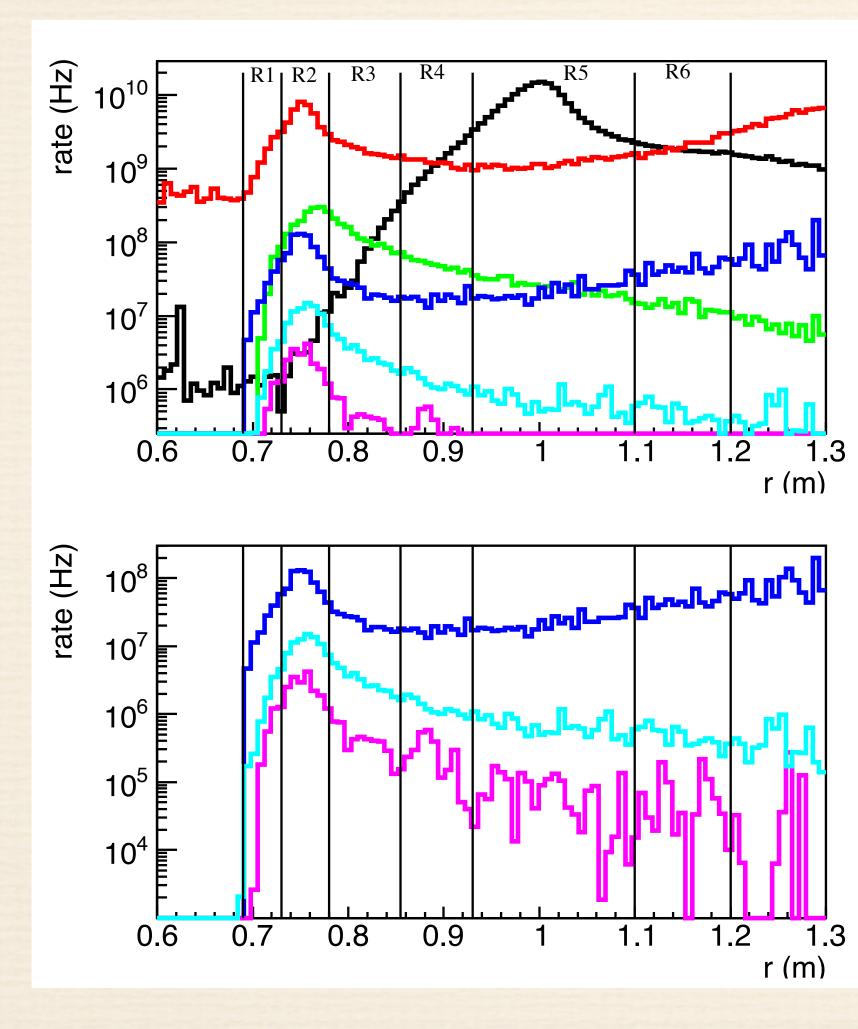
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Old and New Physics with Electron-Electron Scattering

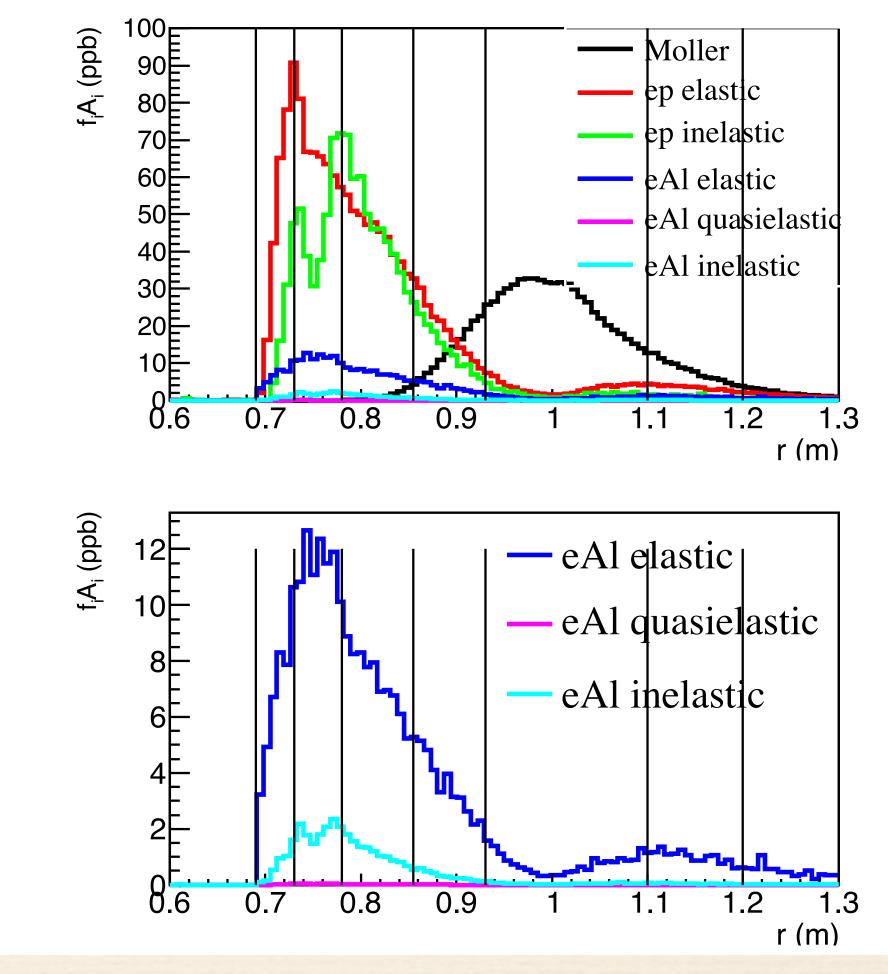


Rate and Asymmetry

33

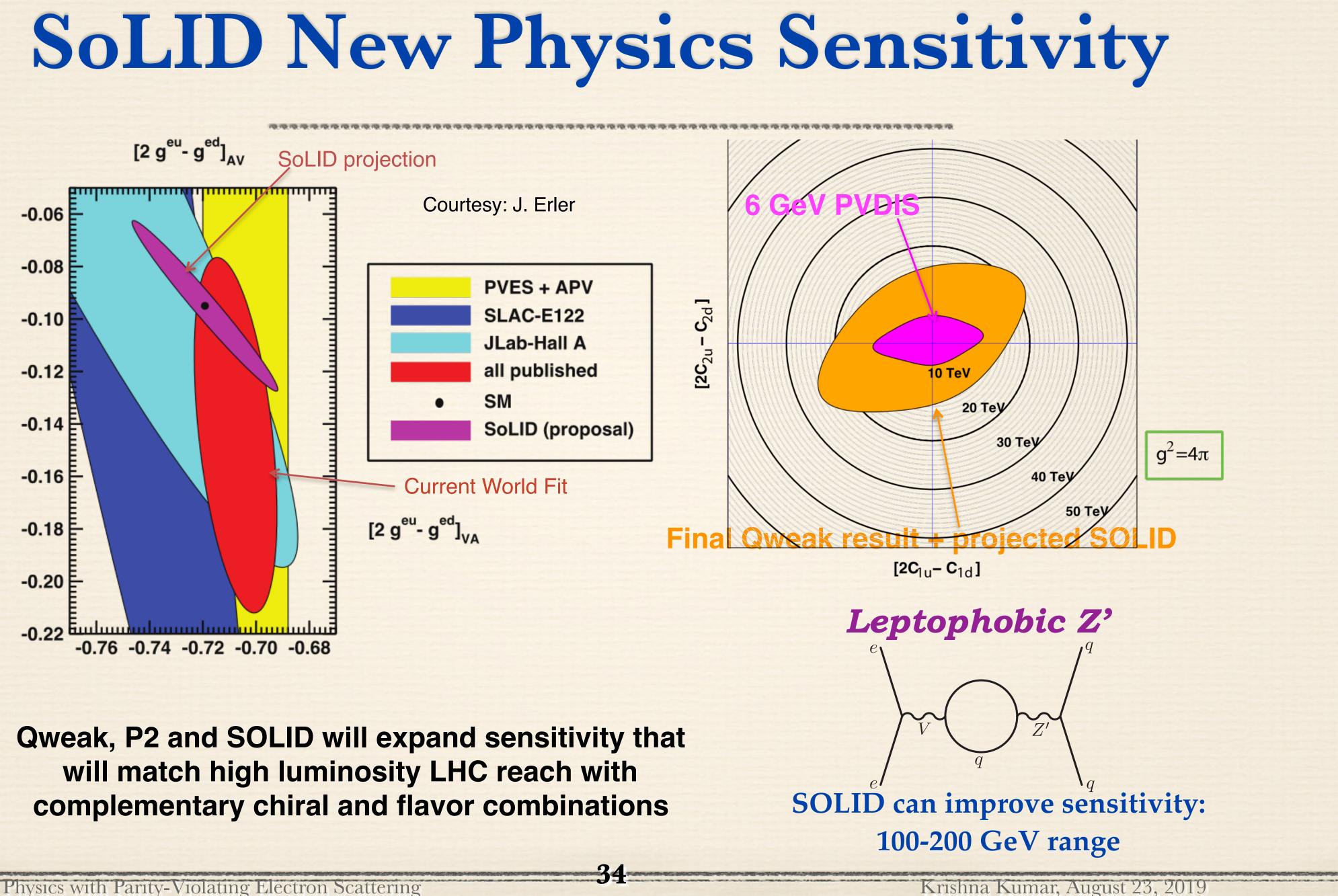


Backgrounds Teleconference



Krishna Kumar, November 24, 2015







Normalization Errors

Goal for total systematic error ~ 2% achieved!

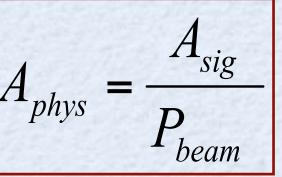
Systematic Error	Absolute (ppm)	Relative (%)	
Polarization	0.0083	1.3	1
Detector Linearity	0.0076	1.2	
Beam current normalization	0.0015	0.2	Two i
Rescattering	0.0001	0	polari
Transverse Polarization	0.0012	0.2	Scatte
Q ²	0.0028	0.4	Both 1.5 ^o
Target Backing	0.0026	0.4	sub-
Inelastic States	0	0	Sub-
TOTAL	0.0140	2.1	

4-momentum transfer $Q^2 = 4 E E' \sin^2 \frac{\theta}{2}$

calibration

E: spin precession in machine E': NMR in HRS B field scattering angle: survey ~ 1 mr

Q² distribution obtained by low rate runs; trigger on quartz pulse-height

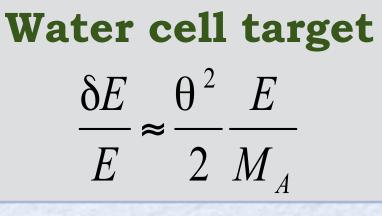


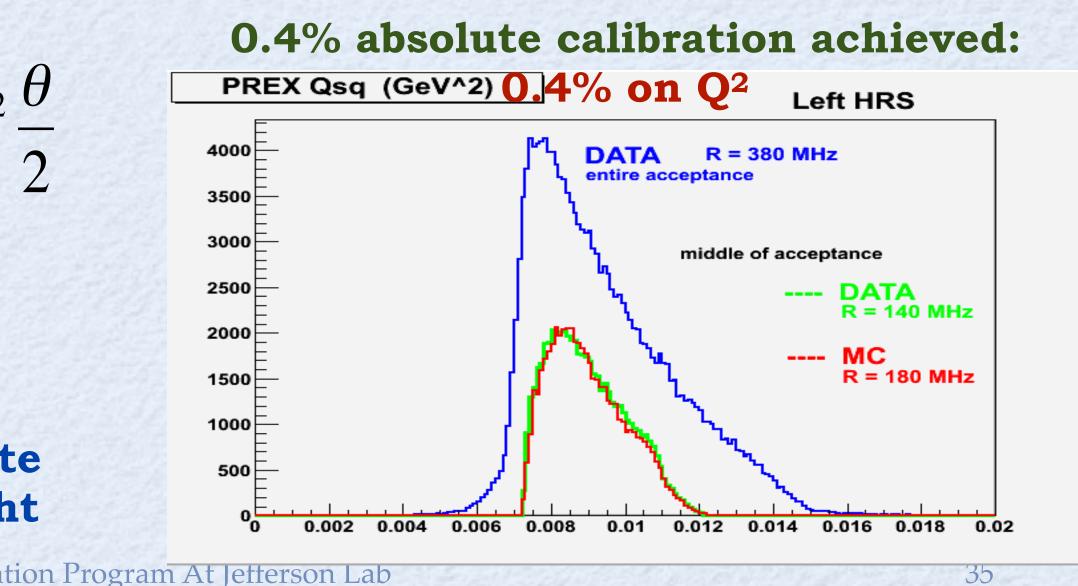
independent methods, rized Møller and Compton tering

h methods achieved ~ %: expected to reach **1% for PREX-II/CREX**

Absolute angle calibration via nuclear recoil variation

Recoil is large for H, small for nuclei





Final Result

$A_{PV} = 0.656 \ ppm \pm 0.060(stat) \pm 0.014(syst)$

Measured A_{PV}

PRL 108 (2012) 112502

 $\bar{q} = 0.475 \text{ fm}^{-1}$

Weak density at one Q²

Small corrections for $G_E^n \ G_E^s \ MEC$

Atomic Parity Violation

Neutron density at one Q²

Assume surface thickness good to 25% (MFT)

 $R_W - R_{ch} =$ 0.32 ± 0.12 (expt) ±0.03 (model) fm

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R_n

Correct for Coulomb Distortions

$F_W(\bar{q}) = 0.204 \pm 0.028(\exp) \pm 0.001(\text{model}) \text{ fm}$

PRC 85 (2012) 032501

Mean Field and Other Models

> Neutron Stars

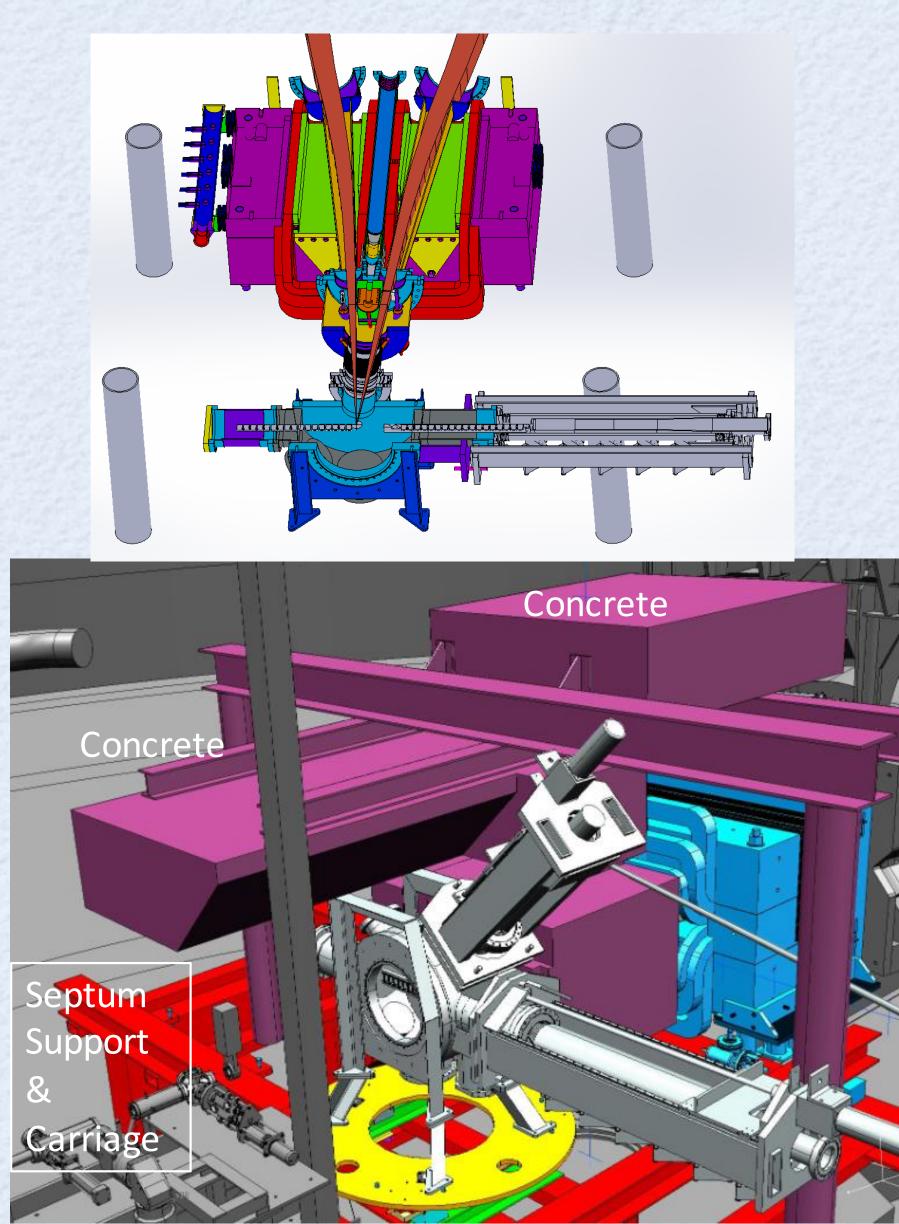
PREX/CREX Parameters

	DDEVI	DDEV II	CDEV
	PREX-I	PREX-II	CREX
Ebeam	1.0 GeV	1.0 GeV	2.1 GeV
A _{PV}	0.65 ppm	0.65 ppm	2.5 ppm
Rate	1 GHz	1.5 GHz	40 MHz
$\delta(A_{PV})_{stat}$	9%	3.5%	4%
$\delta(\mathbf{R}_n)$	0.18 fm	0.07 fm	0.02 fm
Charge	0.1%	0.1%	0.1%
Beam	1.1%	0.5%	0.3%
Non-linearity	1.0%	0.3%	0.3%
Transverse	0.2%	0.2%	0.1%
Beam Polarization	1.1%	0.8%	0.8%
Inelastics	0.1%	0.1%	0.2%
Effective Q ²	0.4%	0.4%	0.4%
Total Systematic	2%	1.1%	1%

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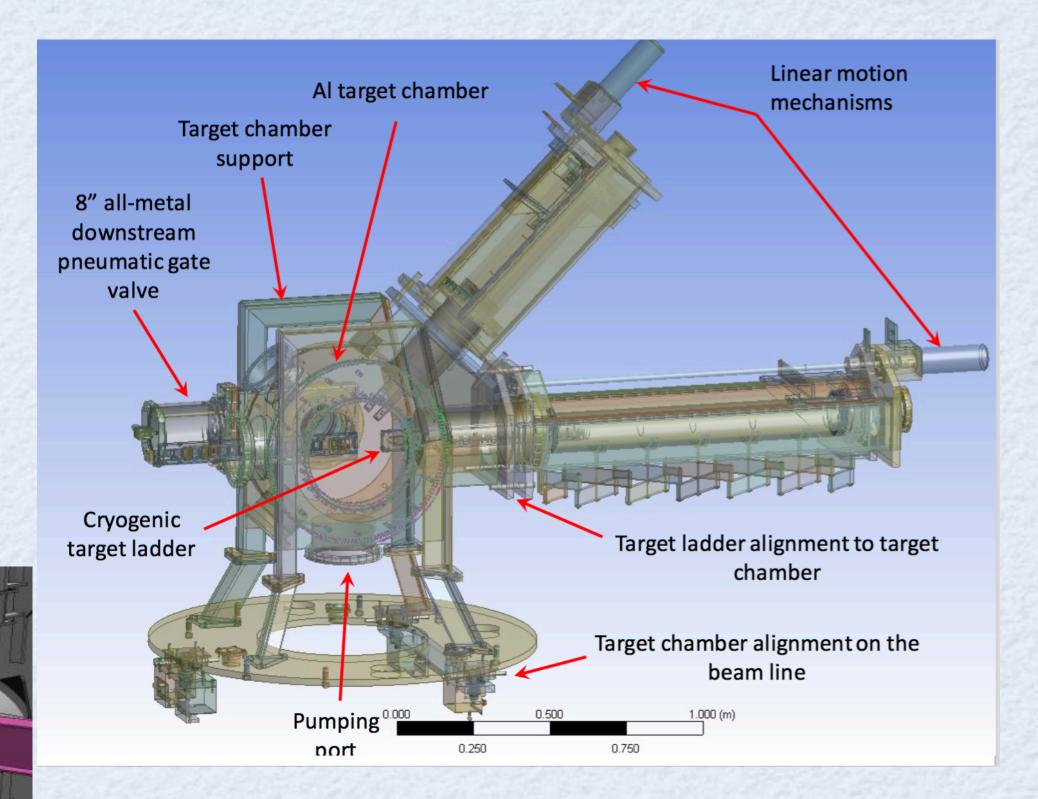


Target Region Redesign



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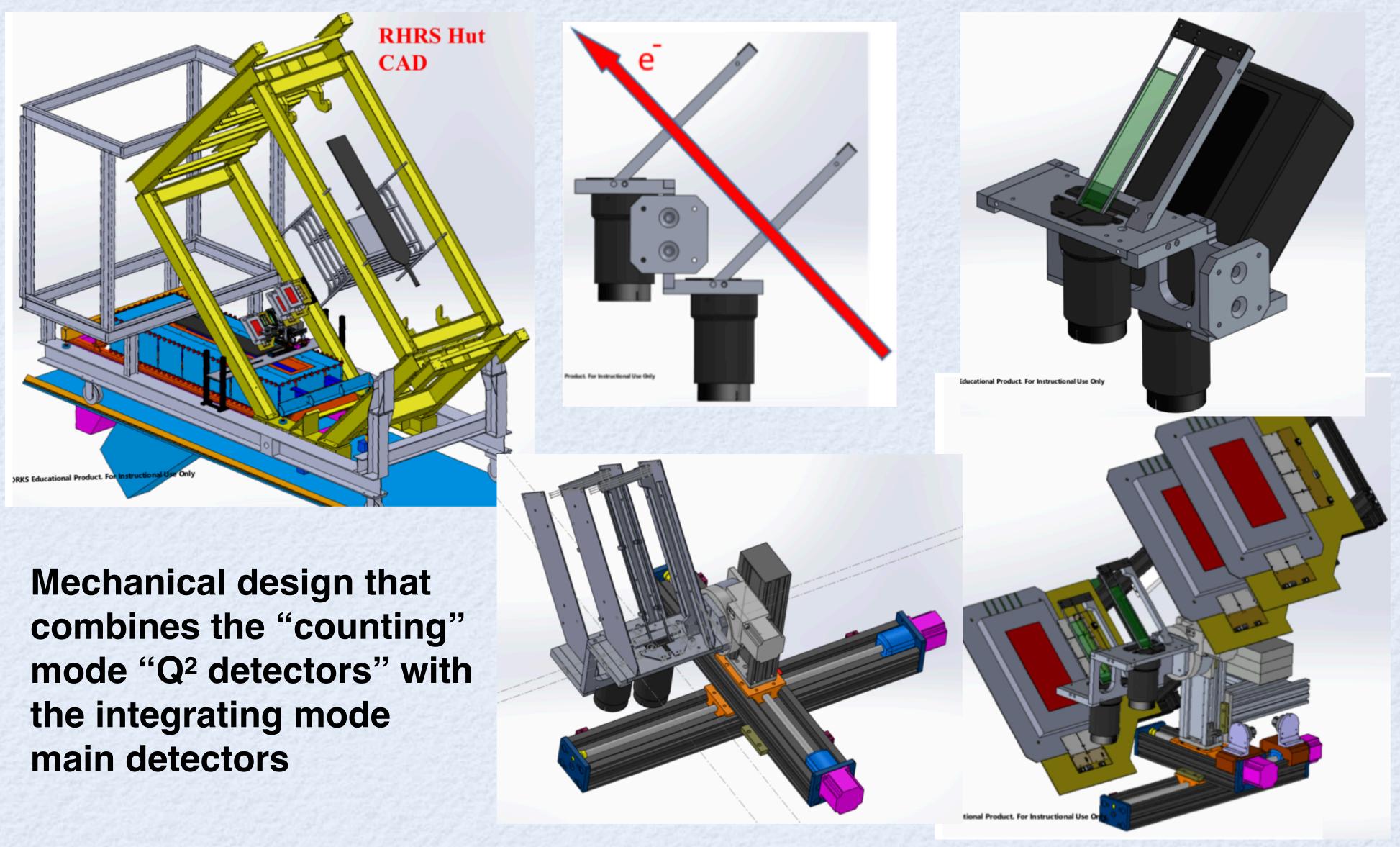
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Extensive simulation, design and engineering effort ongoing for robust, efficient and safe operation of these high luminosity experiments

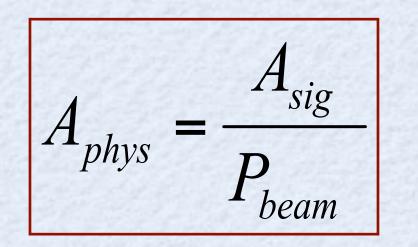


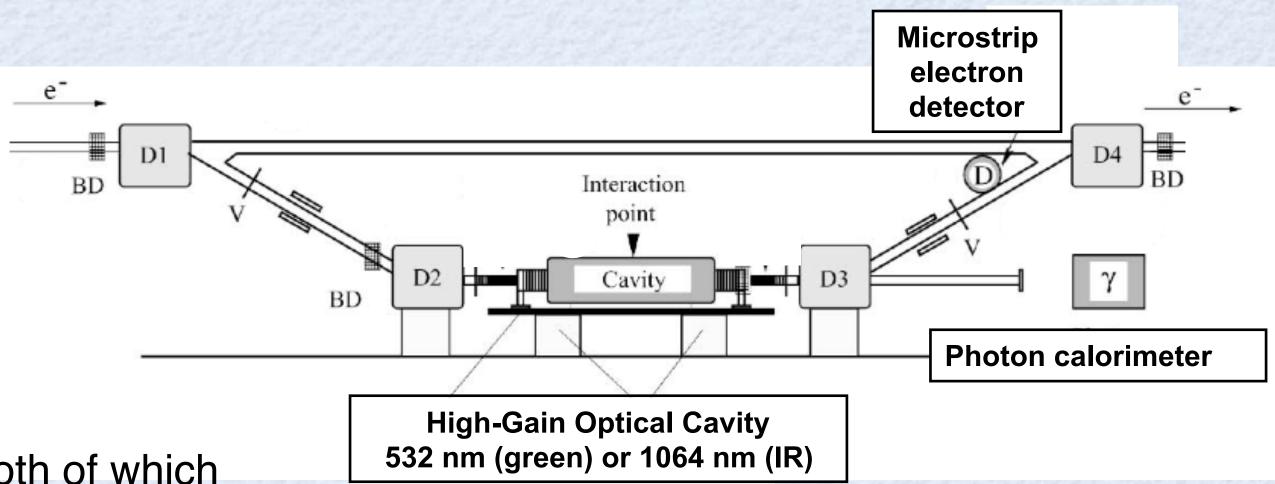
Focal Plane Detectors





Beam Polarimetry

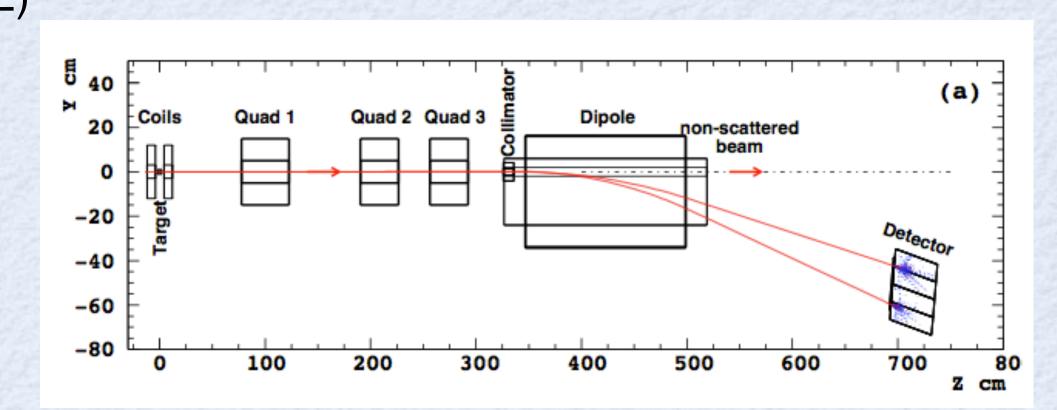




Two independent methods, both of which received recent upgrades

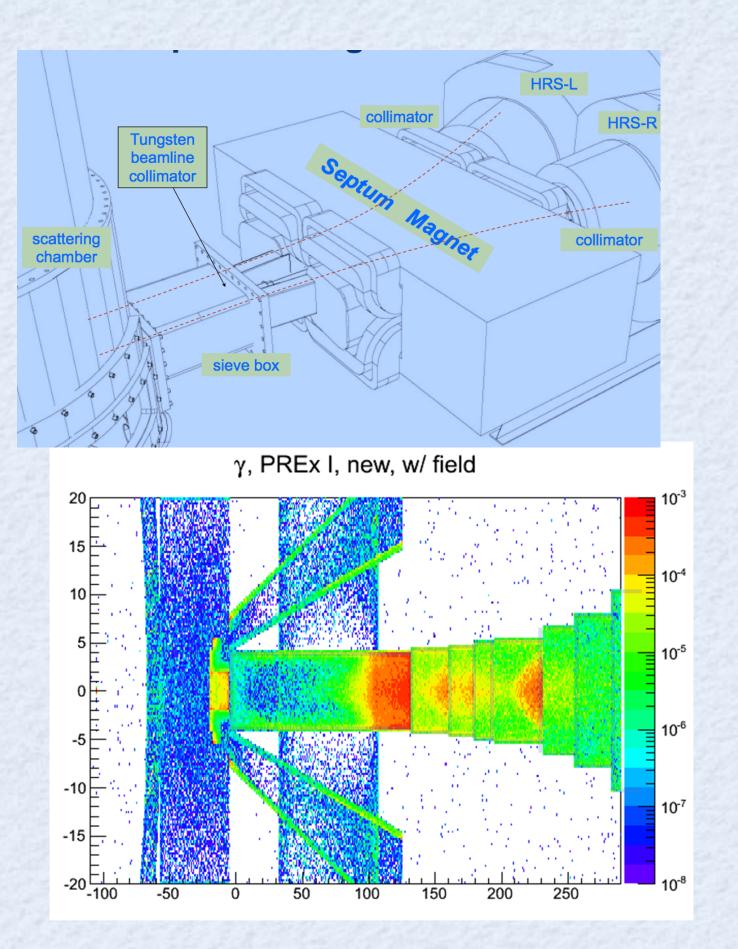
- Compton Polarimeter
 - green laser (increased sensitivity at low E)
 - integrating method (analyzing power)
 - new photon & electron detectors
- Møller Polarimeter
 - electronics and DAQ
 - High field magnet for foil saturation: improved calibration of foil polarization

Both methods expected to reach sub-1% for future measurements: ultimate goal is sub-0.5%



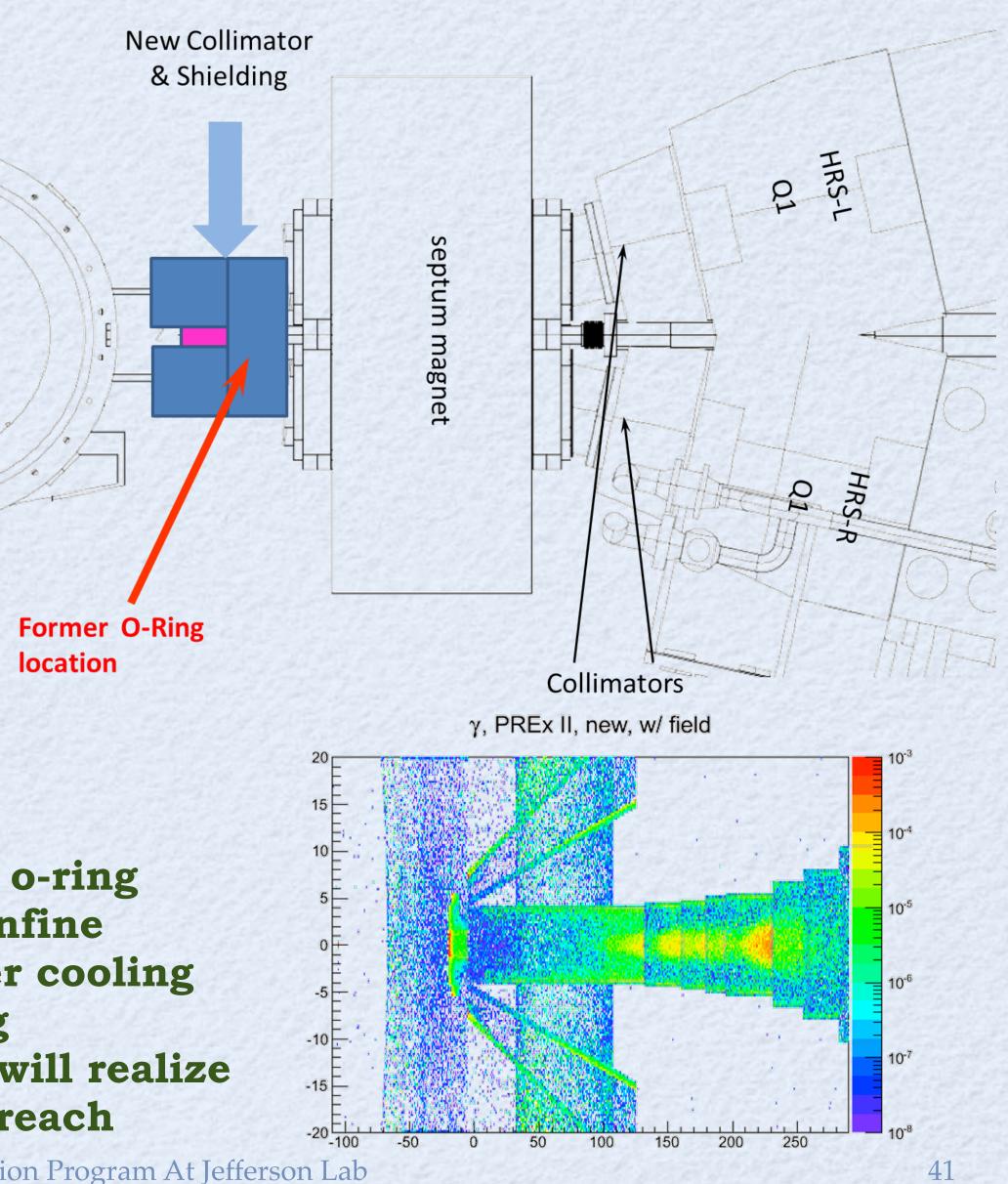
New Beamline Design

target

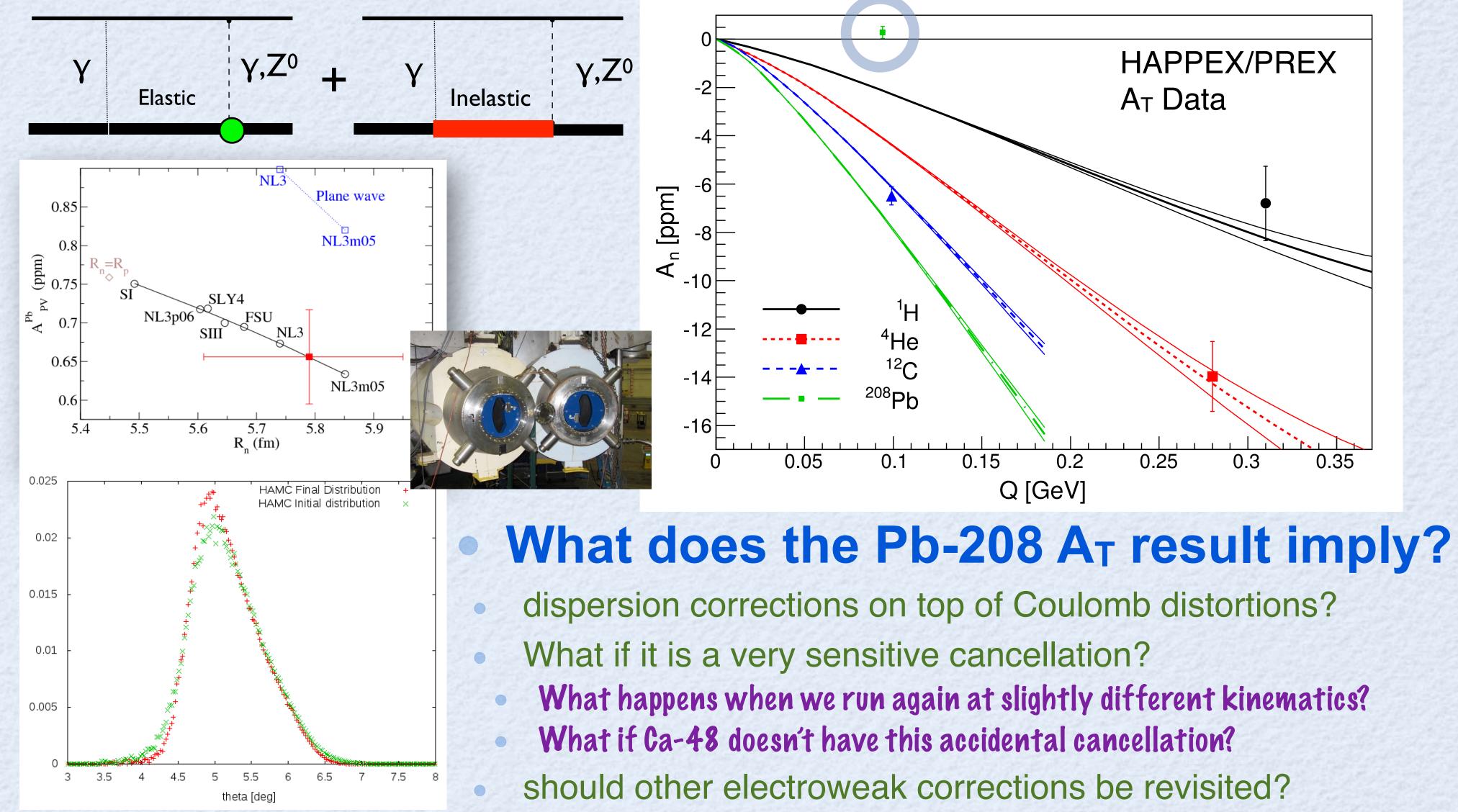


•Redesign beamline seals to eliminate o-ring Neck down tungsten collimator to confine neutrons to one location and add water cooling •Neutrons moderated by new shielding •Small adjustment to septum current will realize an additional ~25% gain in statistical reach

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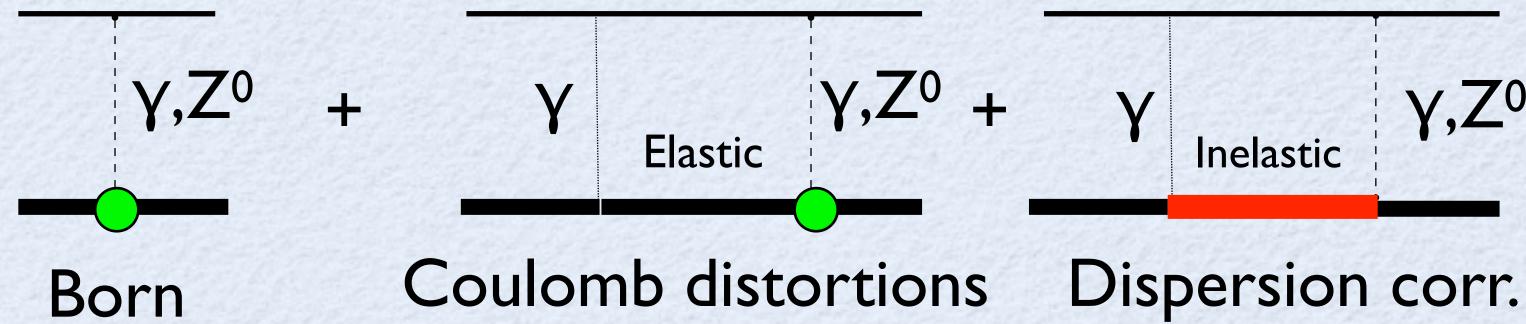


Input from Vector Analyzing Power



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- Sum elastic intermediate states to all orders in $Z\alpha$ by solving Dirac equation for electron moving in coulomb (V) + weak potential (A) of nucleus.
- Coulomb distortions reduce A_{pv} by ~30%, but accurately calculated (uncertainty estimated to be sub-1% of correction)
- Dispersion corrections are of order α (not $Z\alpha$).
- Note: Both Coulomb distortion and dispersion corrections can be important for Transverse Beam Asymmetry An for ²⁰⁸Pb

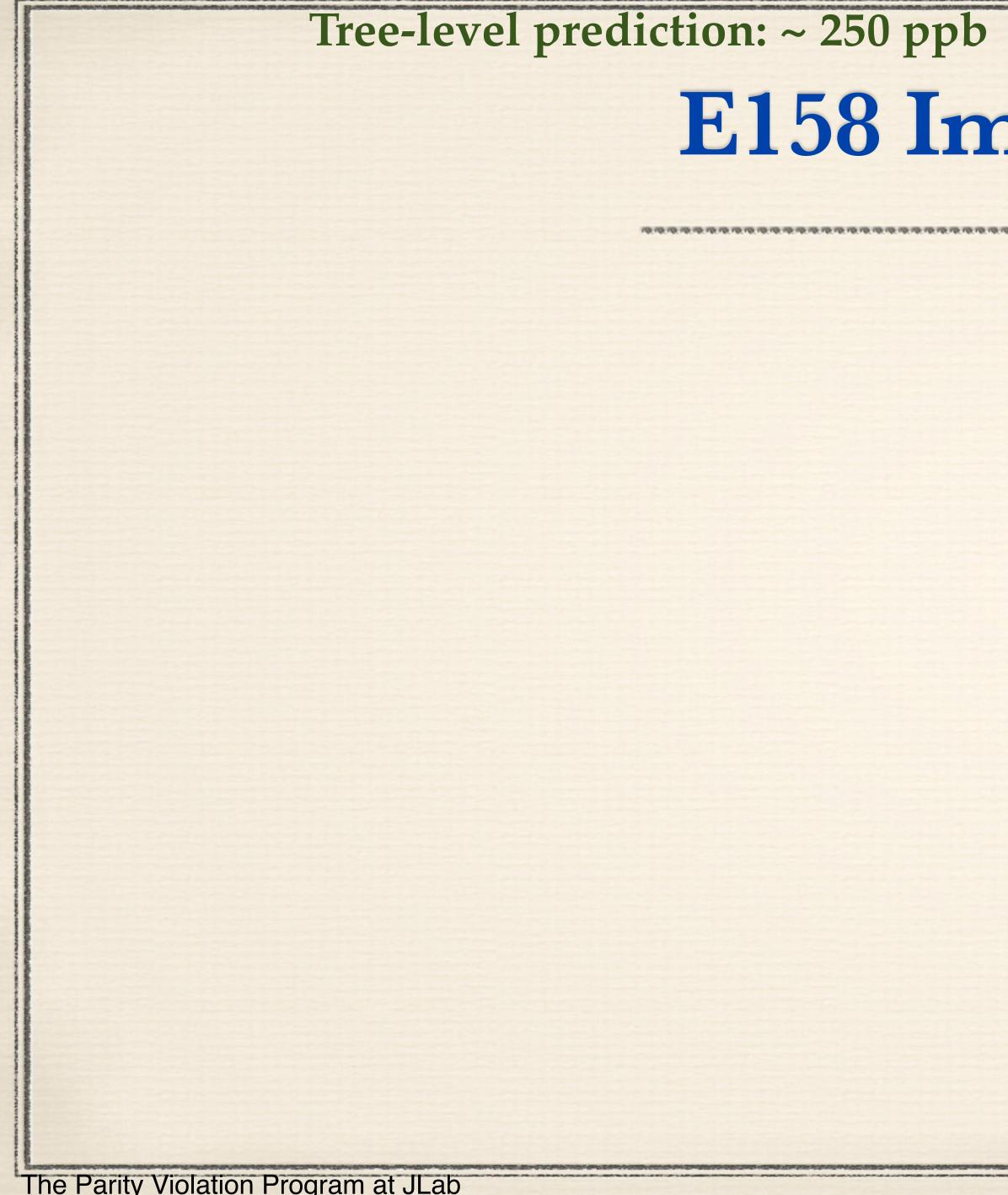
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Radiative Corrections

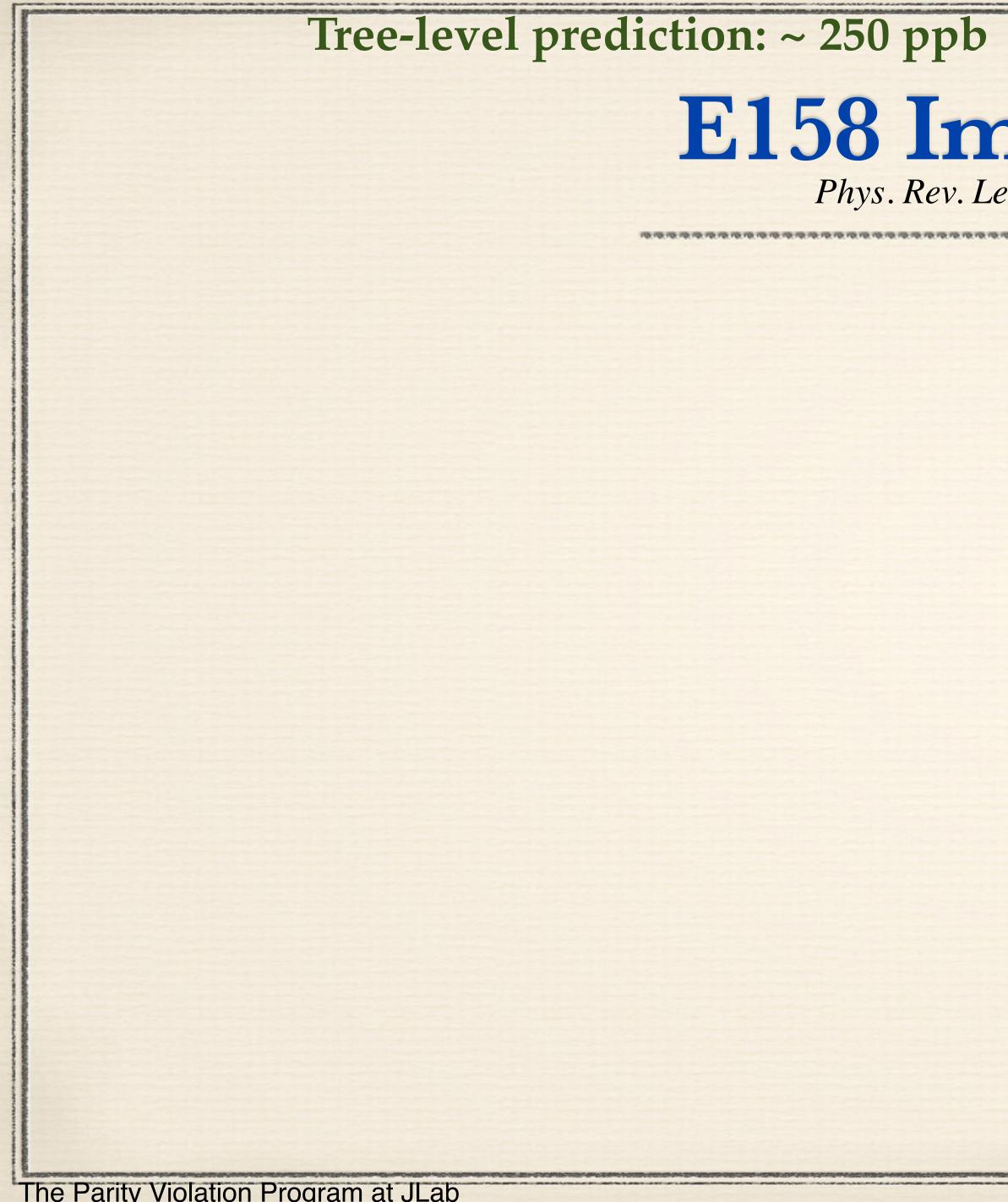
- Coulomb distortions are coherent, order Z α . Important for PREX (Z=82)





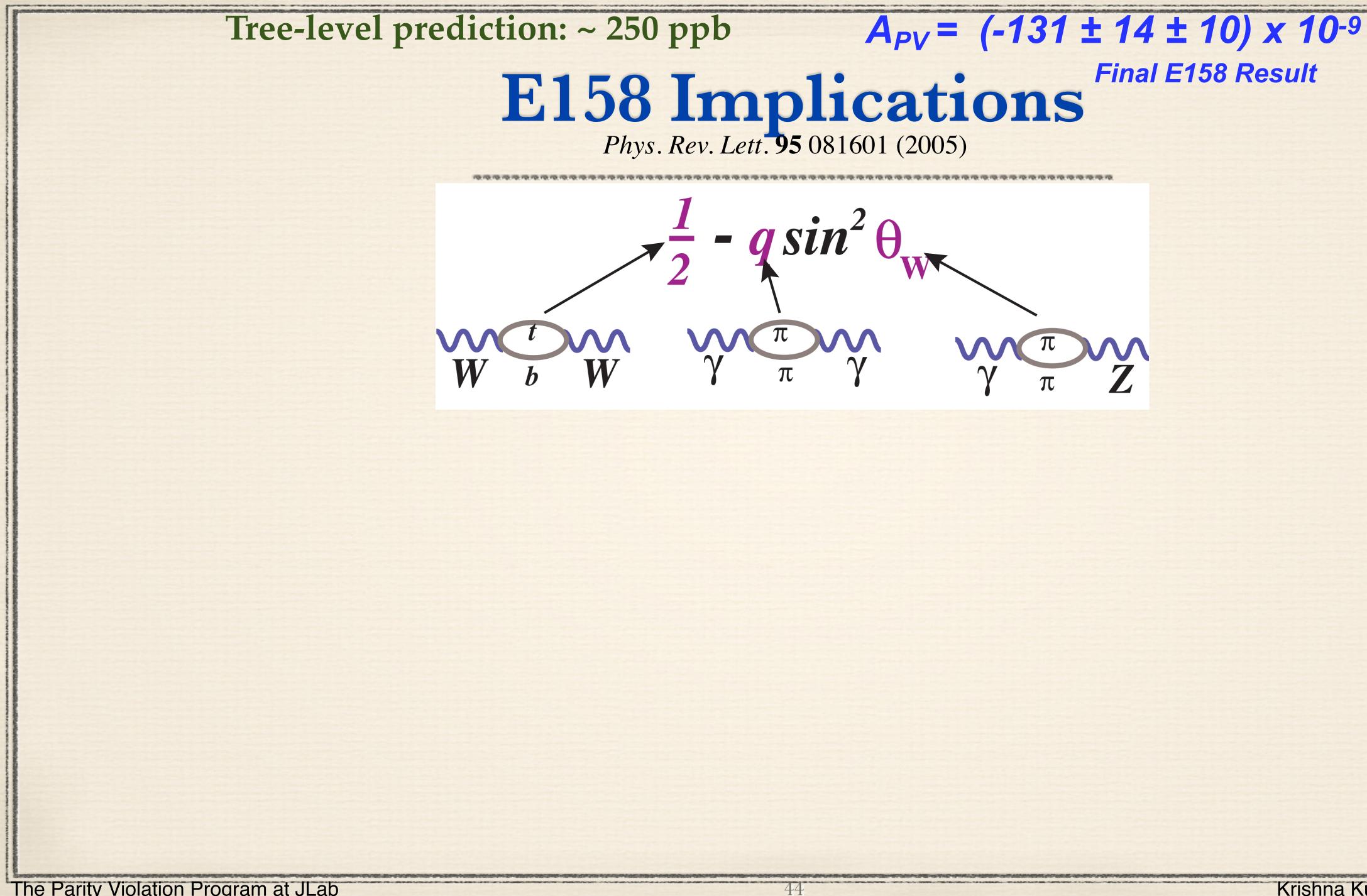
E158 Implications



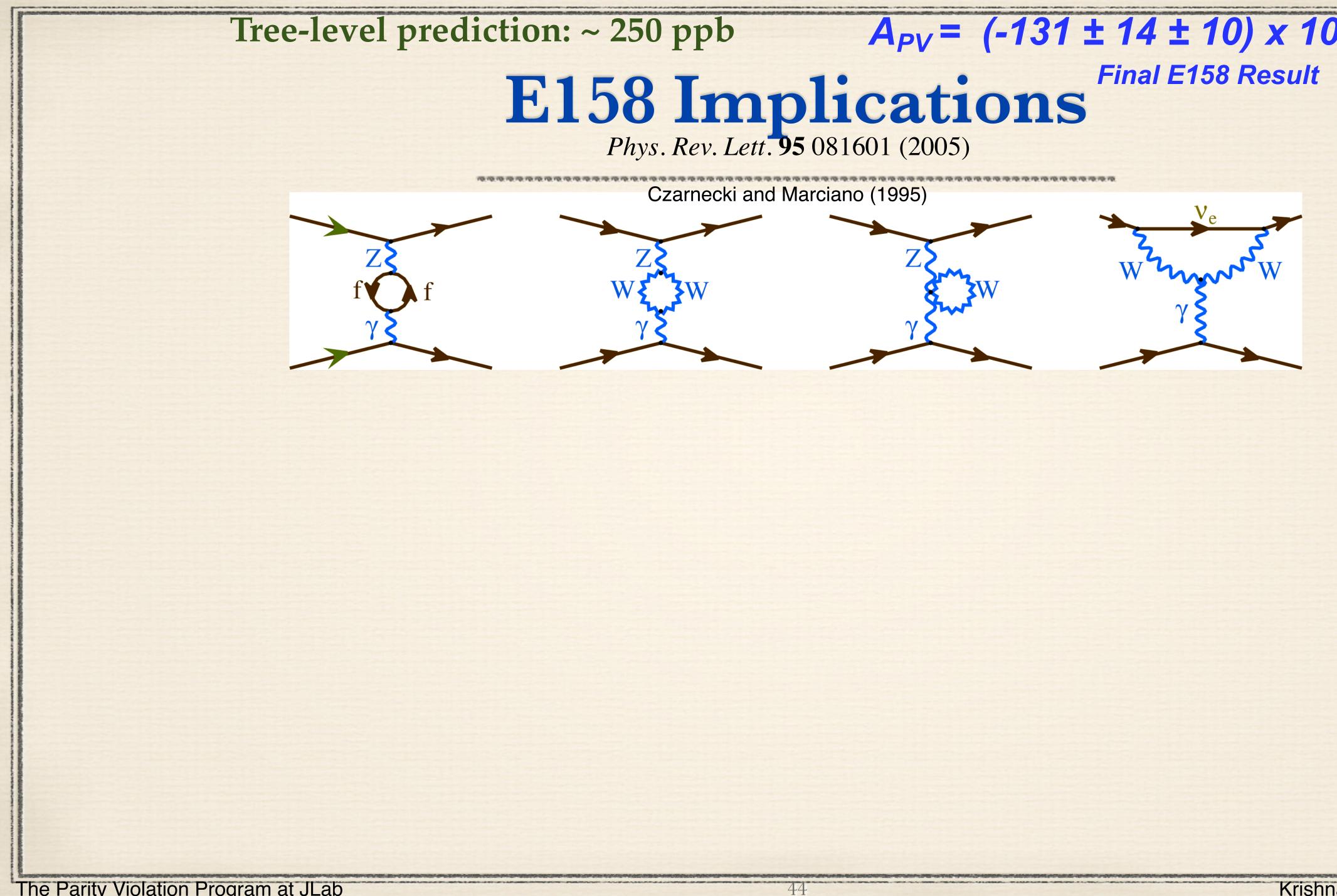


Tree-level prediction: ~ 250 ppb $A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$ E158 Implications Final E158 Result Phys. Rev. Lett. 95 081601 (2005) Final E158 Result







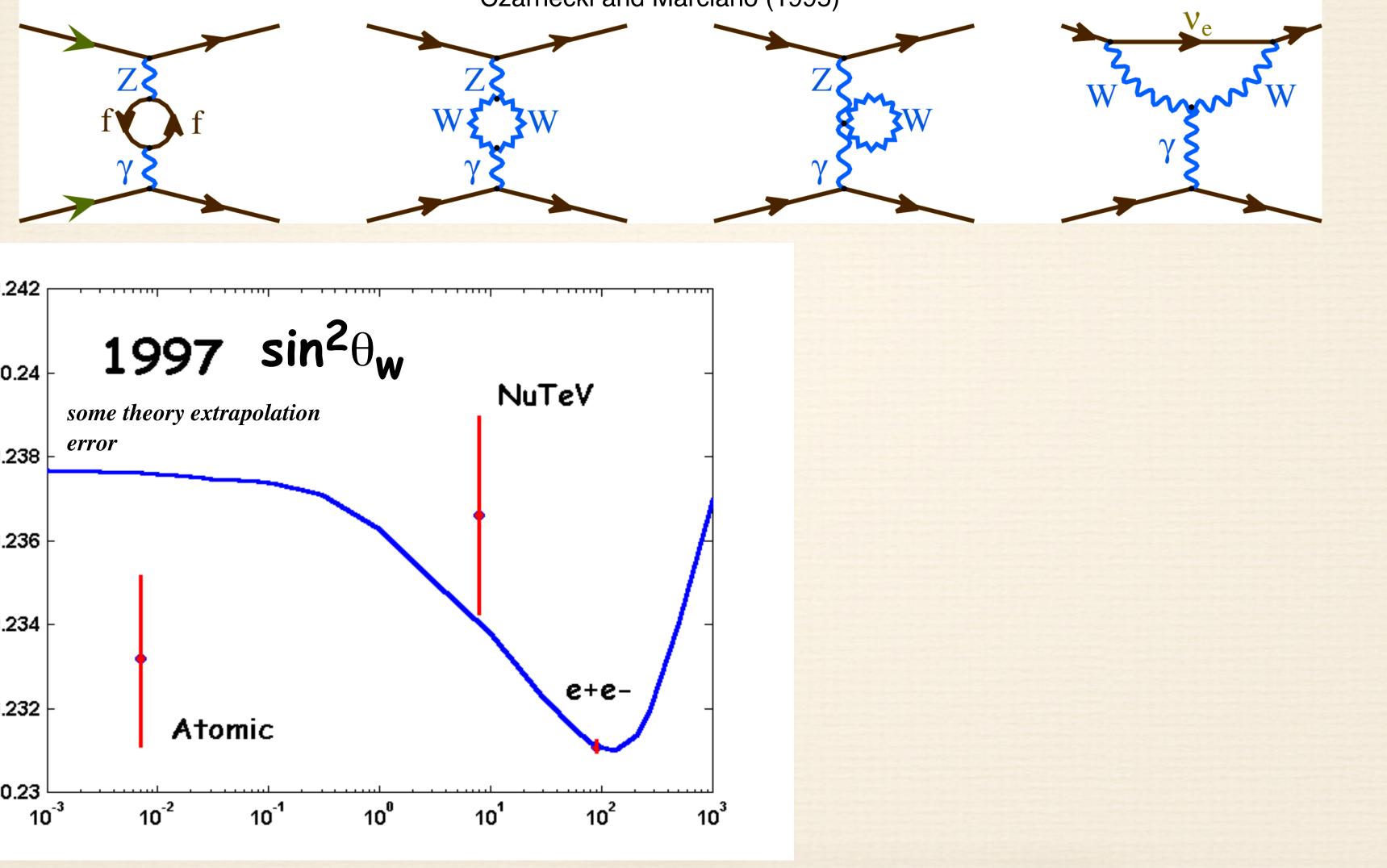


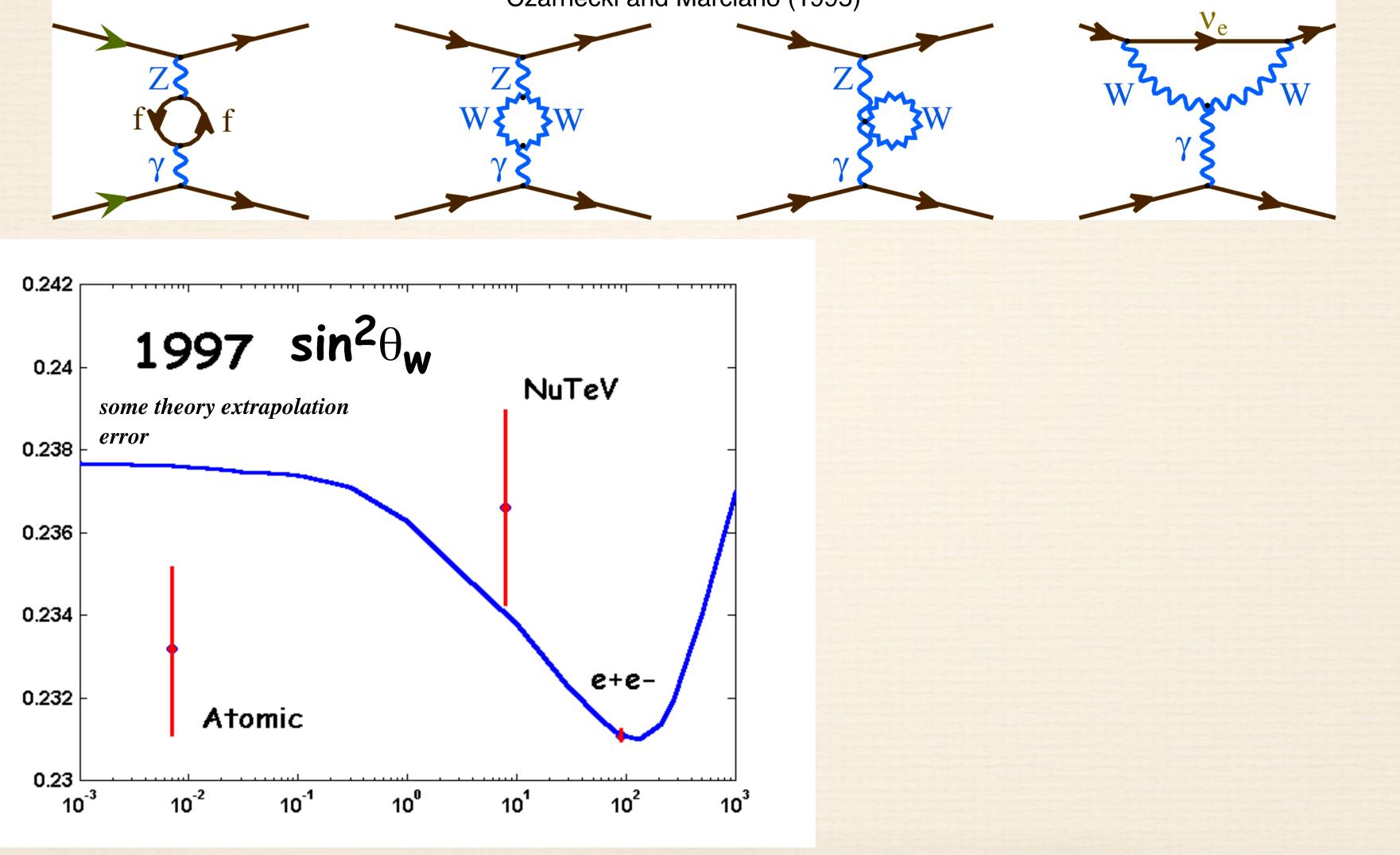
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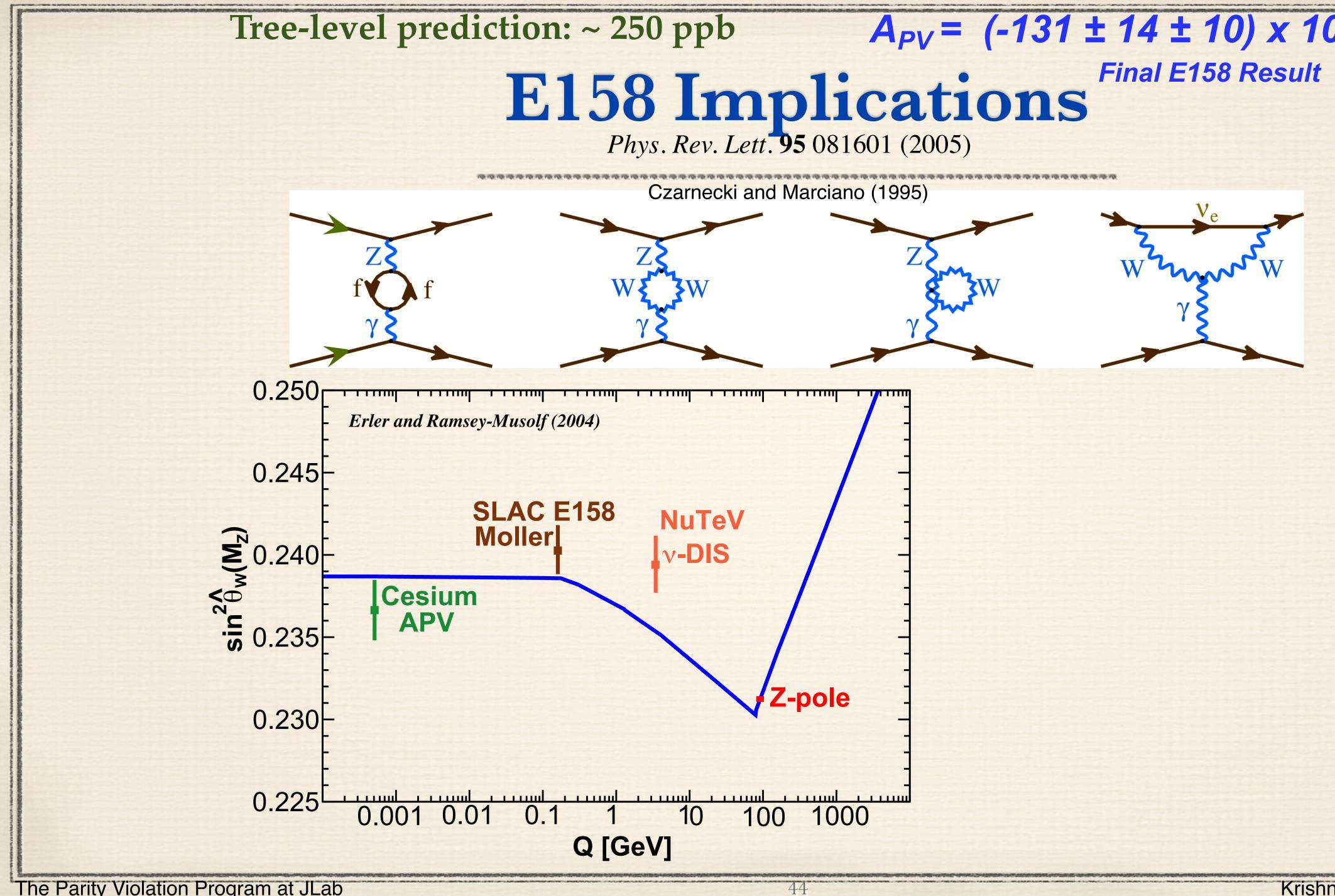
Czarnecki and Marciano (1995)





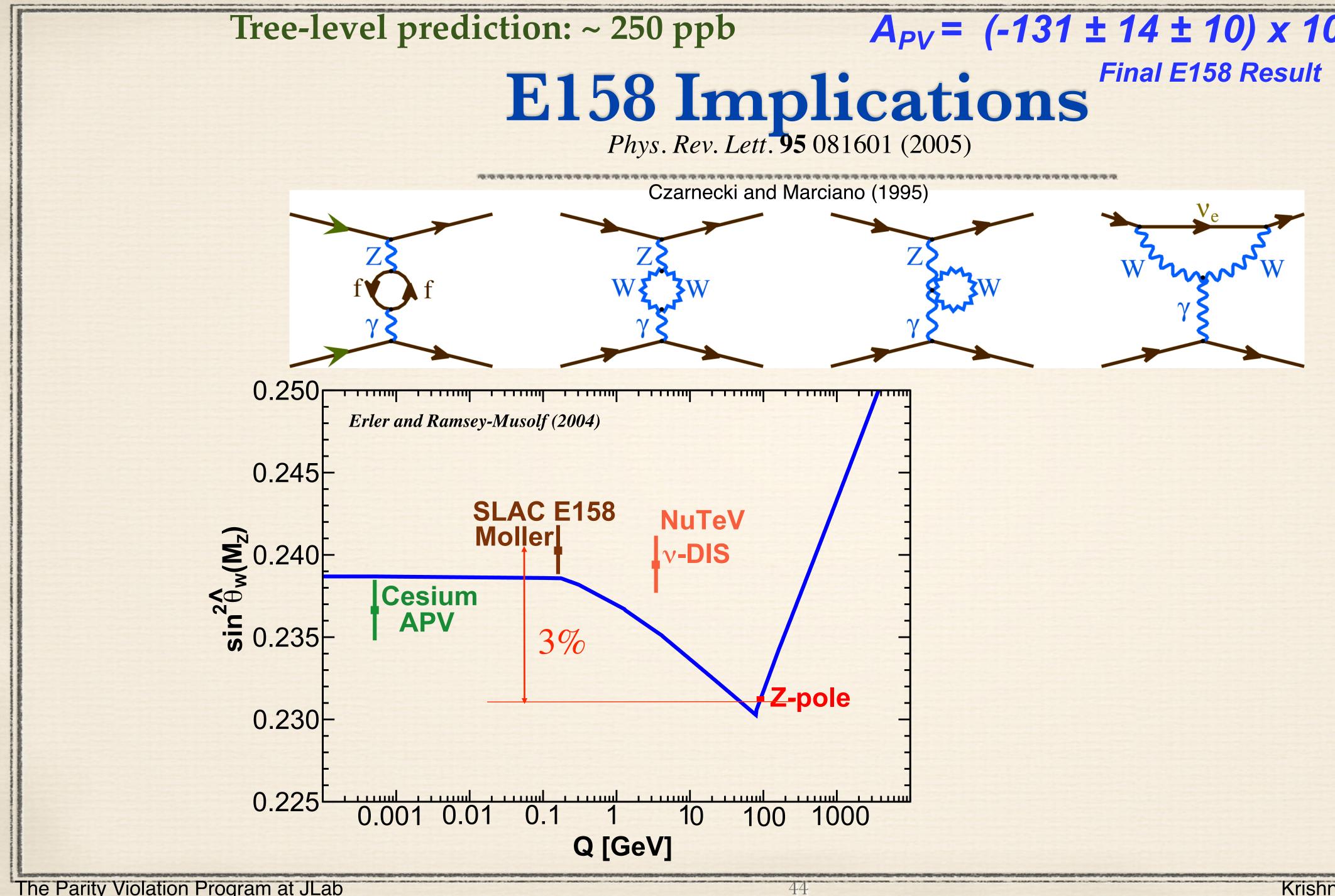
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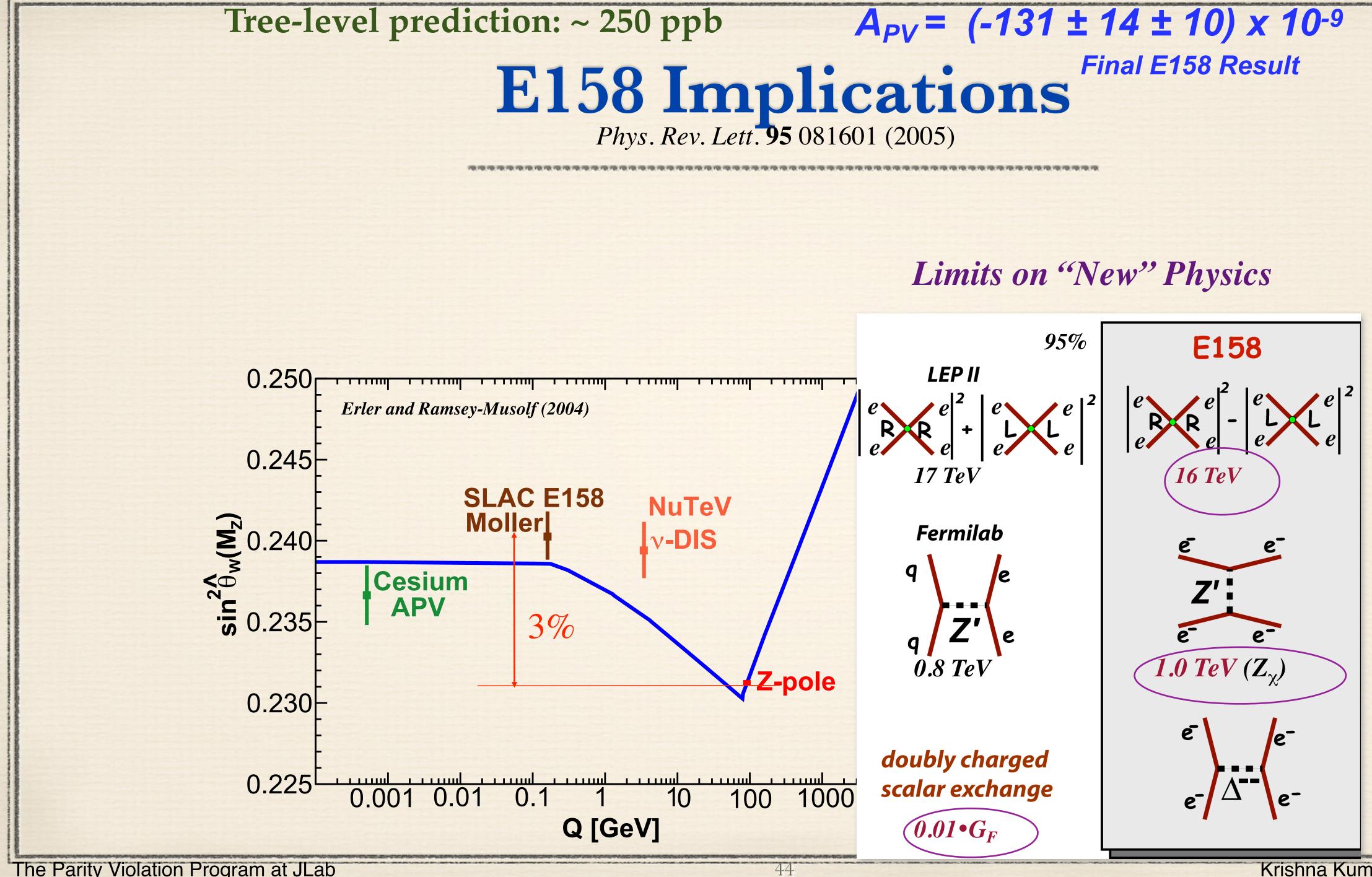
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$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$







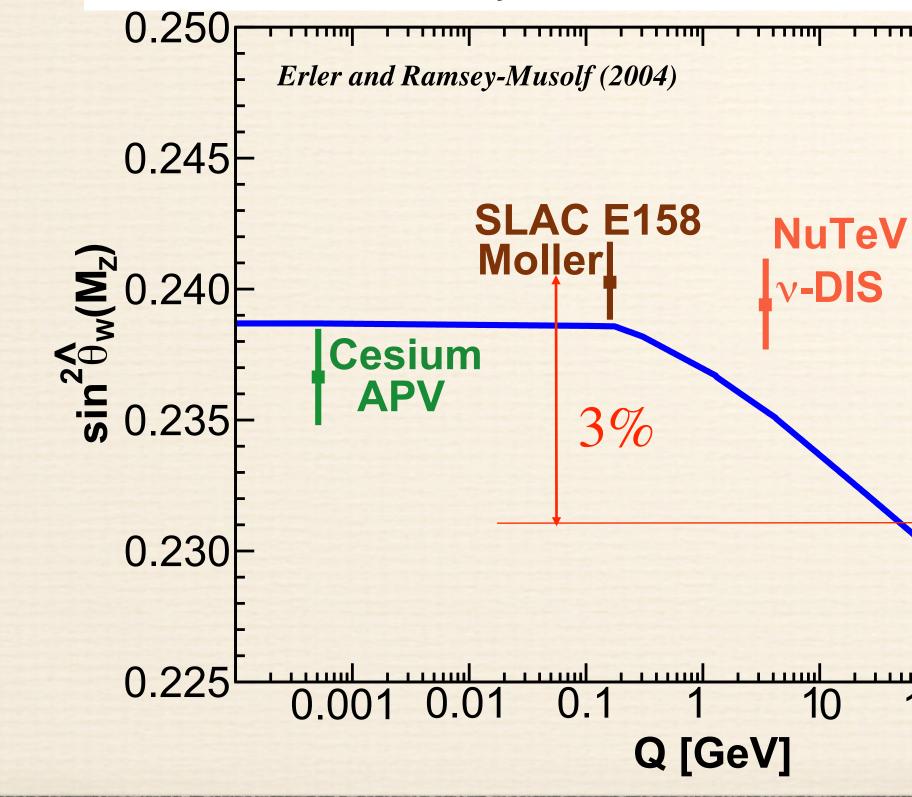
NEWS AND VIEWS

PARTICLE PHYSICS

Electrons are not ambidextrous

Andrzej Czarnecki and William J. Marciano

The best low-energy measurement yet obtained of the electroweak mixing angle — a central parameter of the standard model of particle physics — is 95% E158 the last hurrah for Stanford's powerful two-mile linear accelerator. LEP II 0.250 Erler and Ramsey-Musolf (2004) 0.245 17 TeV **16 TeV SLAC E158 NuTeV** Moller Fermilab v-DIS Cesium **APV** 3% $0 TeV(Z_{\gamma})$ 0.8 TeV -pole 0.230 ē doubly charged 0.225 scalar exchange 0.001 0.01 10 1000 100 Q [GeV] $0.01 \cdot G_F$



The Parity Violation Program at JLab

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Limits on "New" Physics



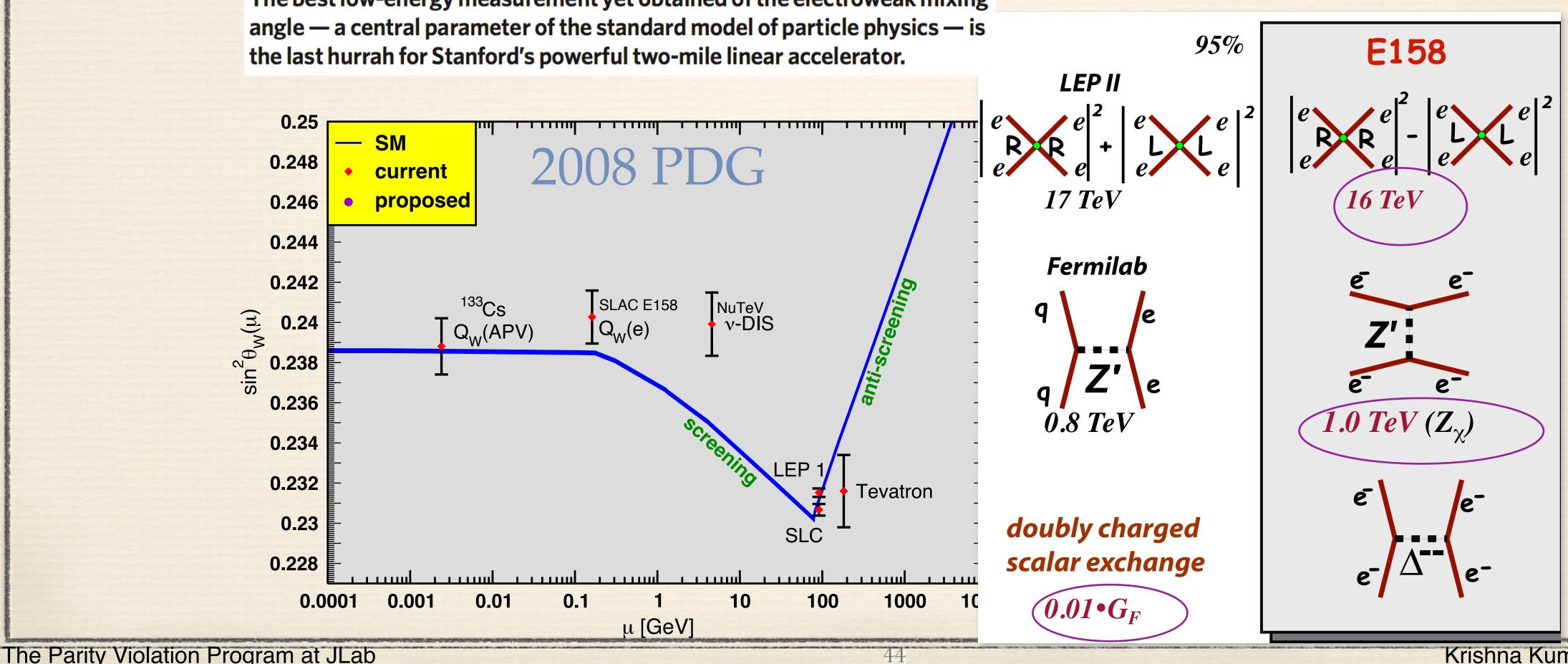
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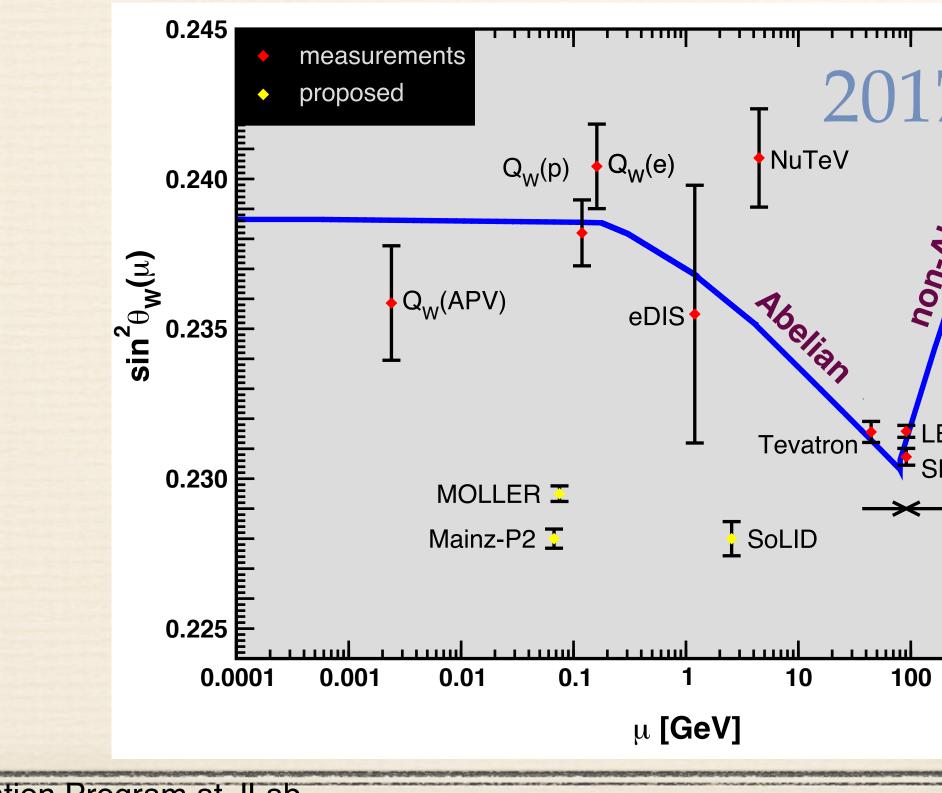
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The Parity Violation Program at JLab

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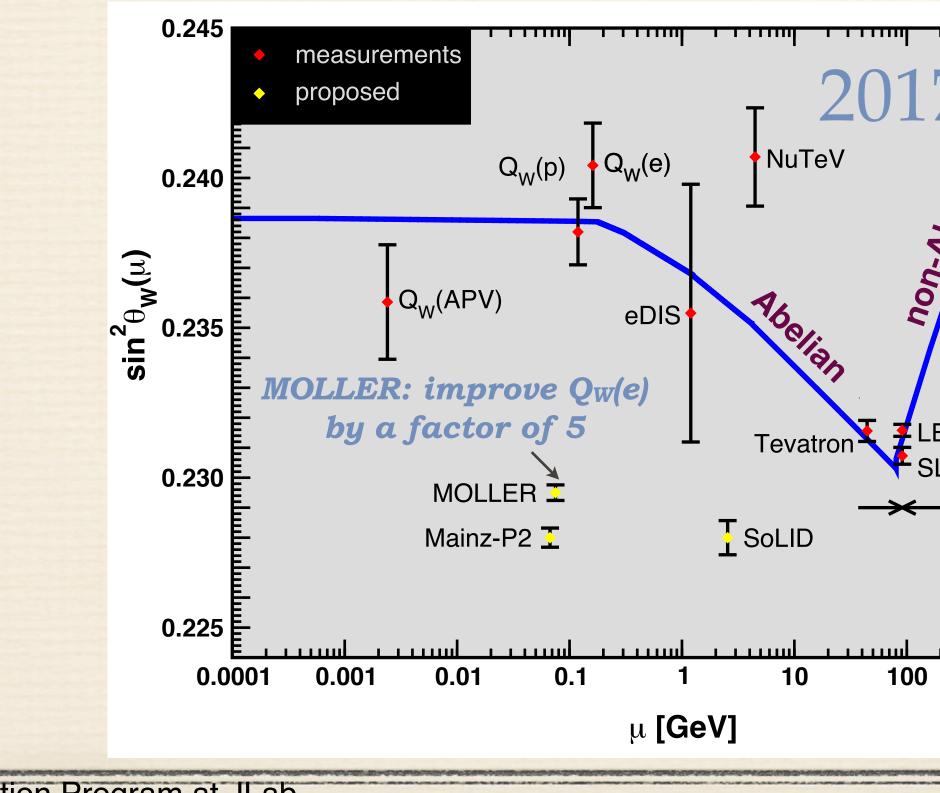
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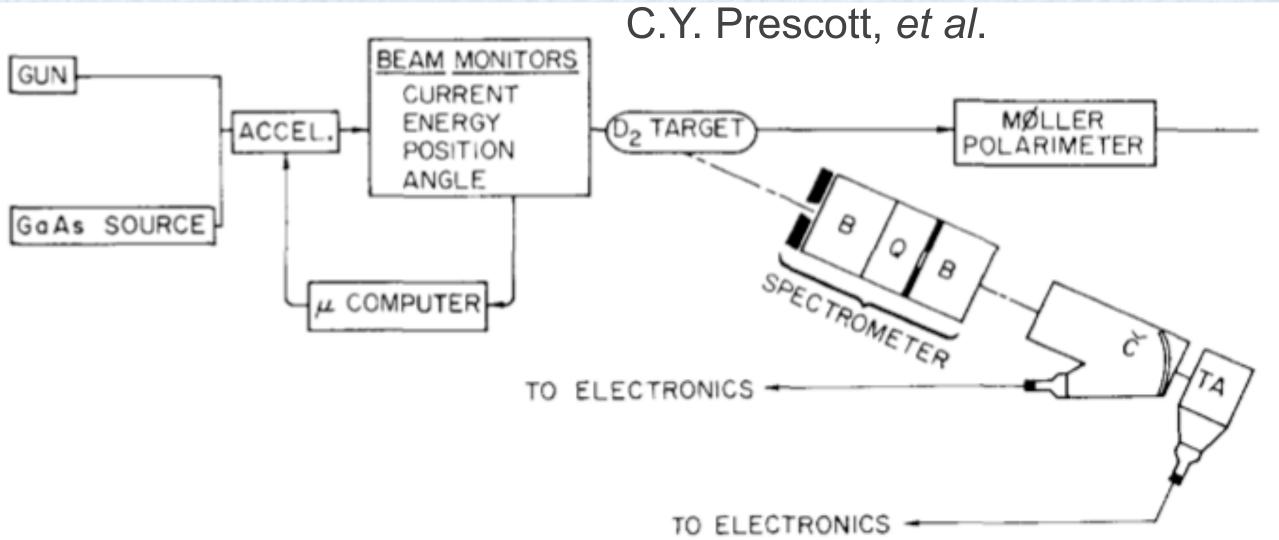
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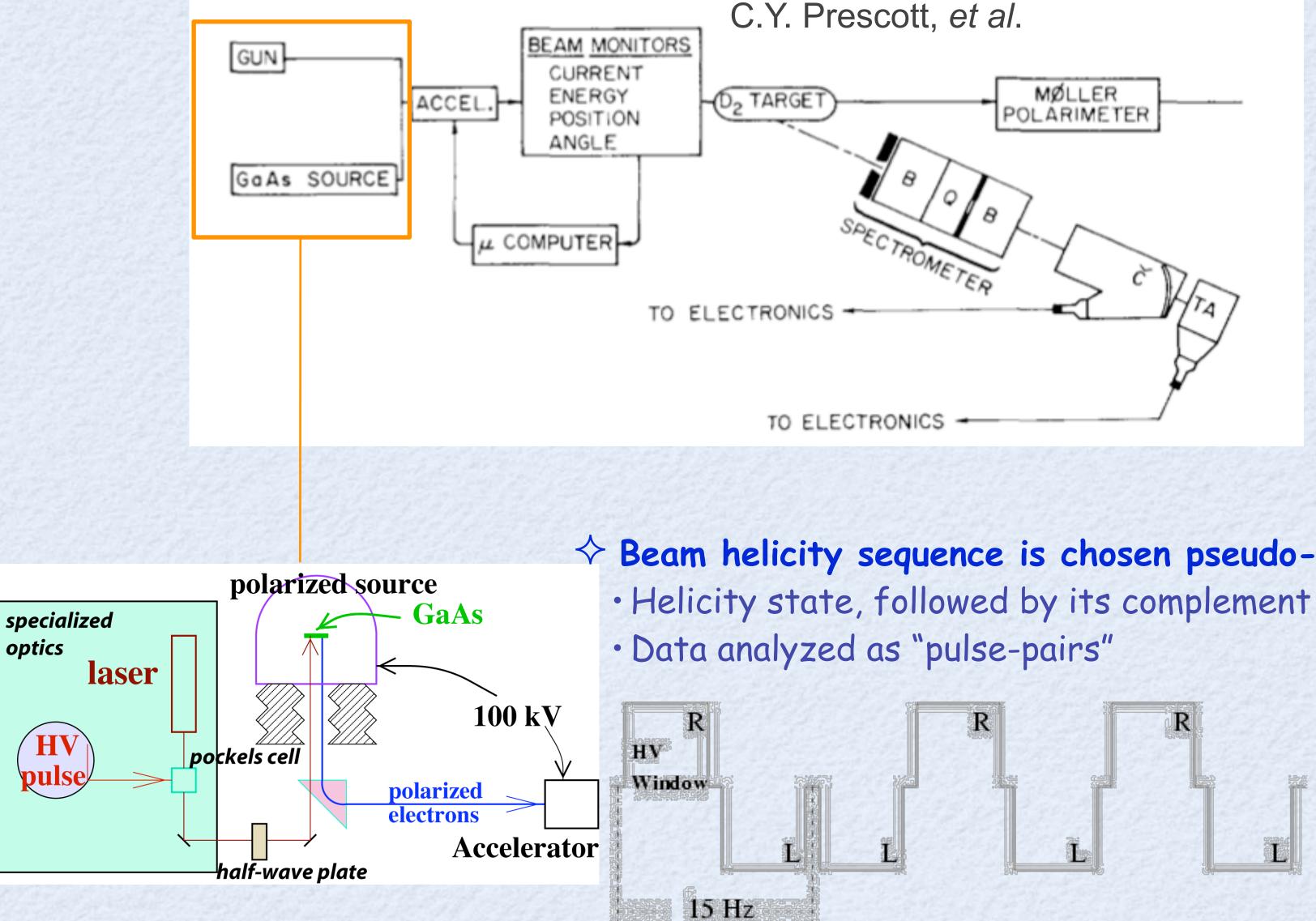


The E122 Experiment at the Stanford Linear Accelerator Center



Krishna S. Kumar

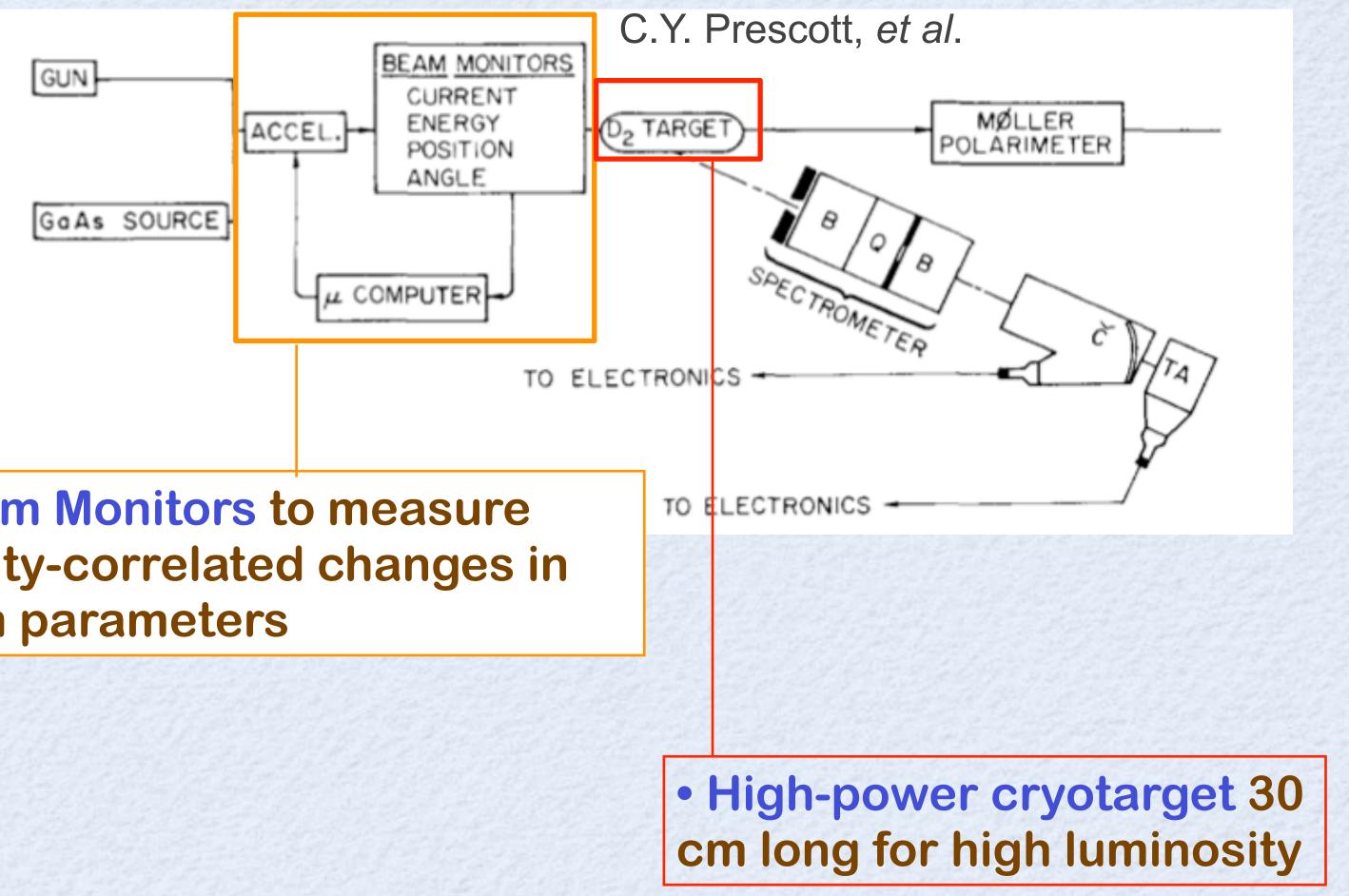
The E122 Experiment at the Stanford Linear Accelerator Center



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 \diamond Beam helicity sequence is chosen pseudo-randomly

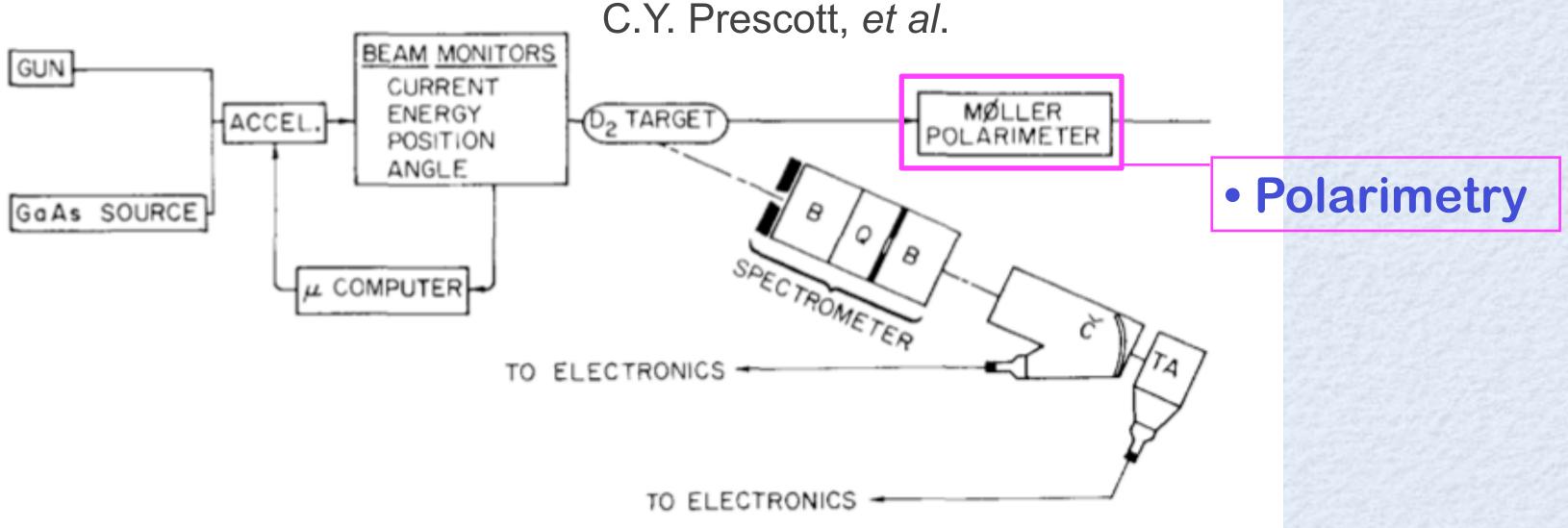
The E122 Experiment at the Stanford Linear Accelerator Center



• Beam Monitors to measure helicity-correlated changes in **beam parameters**

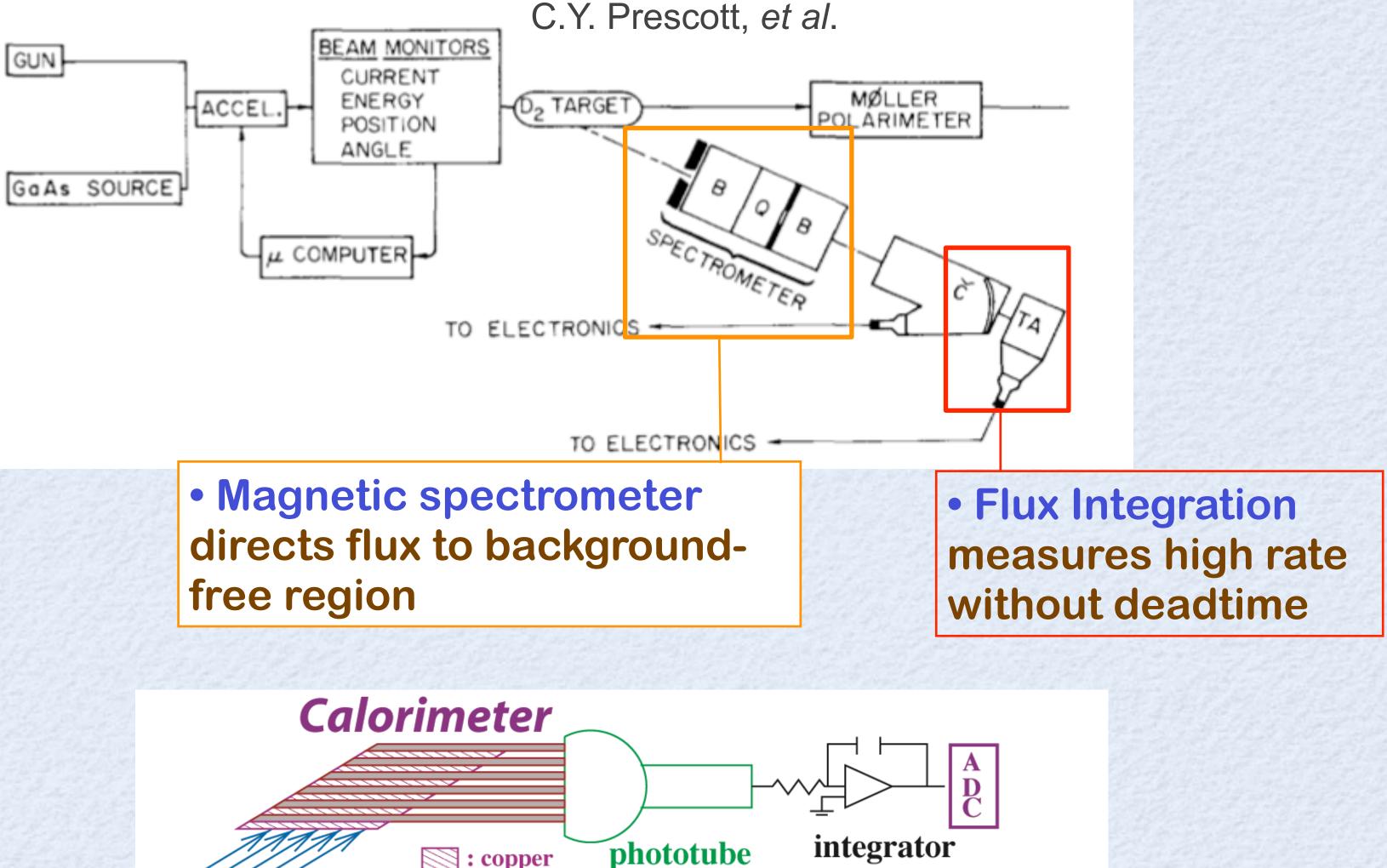
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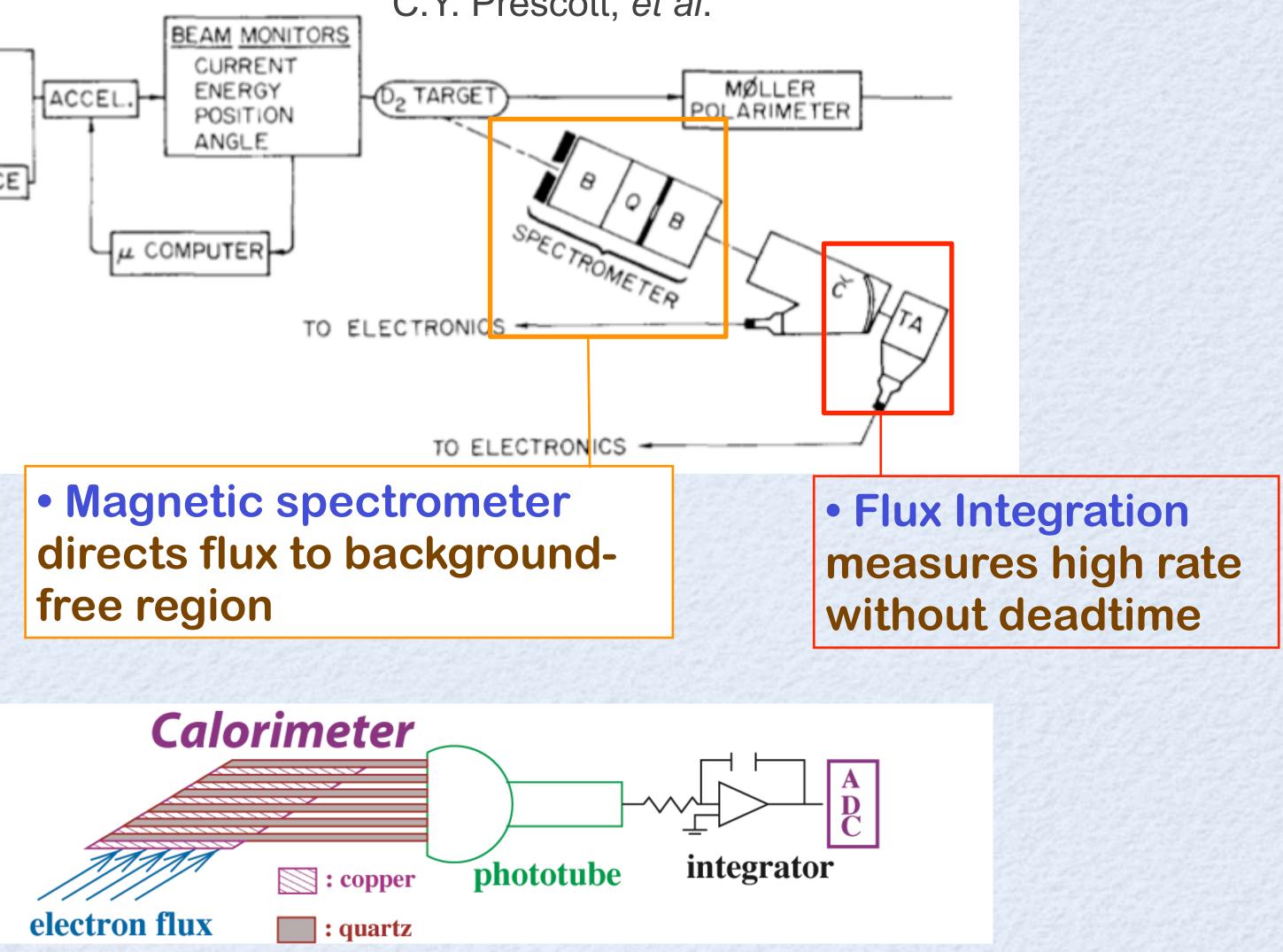
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